

Pseudoscalar Transition Form Factors from Rational Approximants

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Work in collaboration with R. Escribano and P. Masjuan
[Phys.Rev. D86 (2012) 094021, Phys.Rev. D89 (2014) 034014, ...]

MESON 2014, Cracow, 2 June 2014

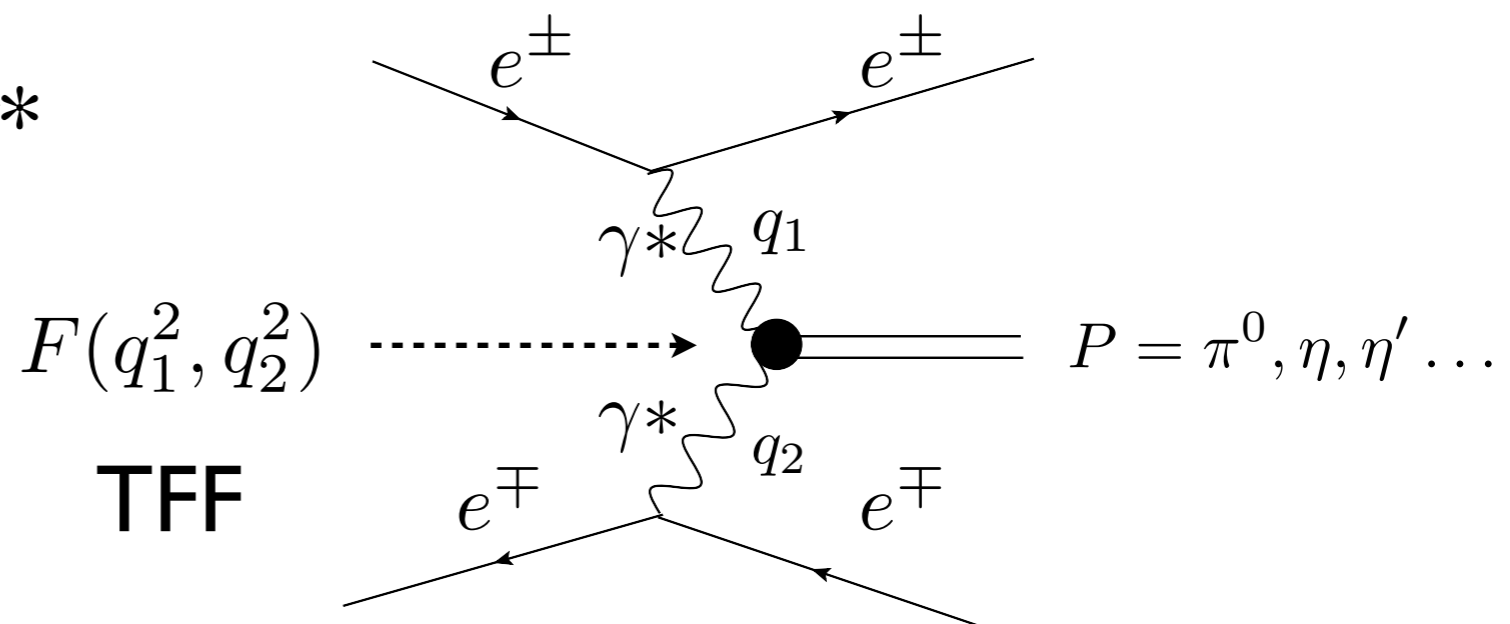
Outline

Pseudoscalar Transition Form Factors

- Pseudoscalar Transition Form Factors
- Parameterization using Rational Approximants
- Results from space-like data alone
 - ▶ η - η' mixing
- Implications from space-like + time-like data
- Conclusions and outlook

Pseudoscalar Transition Form Factors

- Study of $ee \rightarrow ee\gamma^*\gamma^*$ with $\gamma^*\gamma^* \rightarrow \pi, \eta, \eta'$



- Meson Structure
 - Transition Form Factors (TFF) give access to Meson Distribution Amplitudes
- Precision Tests of the Standard Model
 - Relation to mixing parameters and muon anomaly $(g-2)_\mu$

How do we do that?

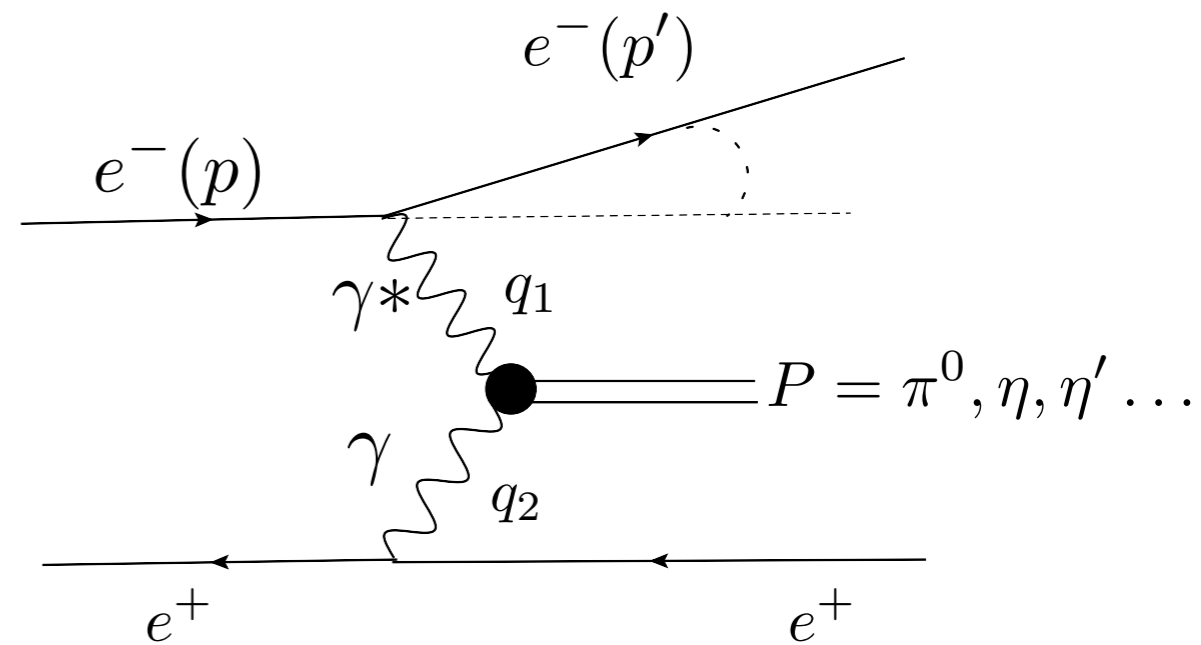
- Single Tag Method can access the Meson Transition Form Factor

Selection criteria

- 1 e^- detected
- 1 e^+ along beam axis
- Meson full reconstructed

Momentum transfer

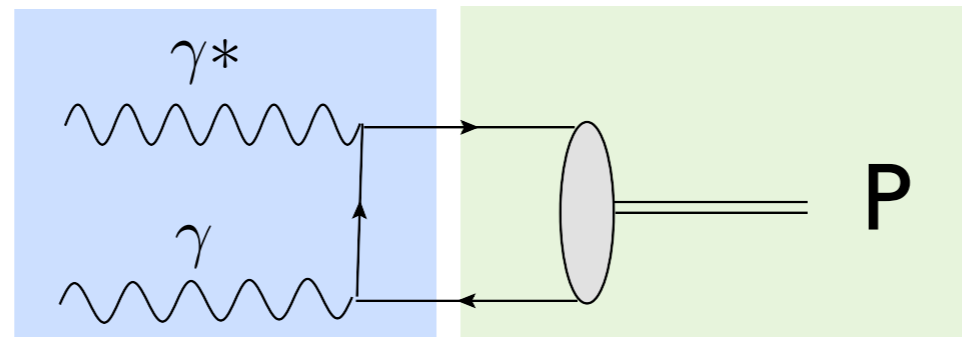
- tagged: $Q^2 = -q_1^2 = -(p - p')^2$
⇒ highly virtual space-like (SL) photon
- untagged: $Q^2 = -q_2^2 \sim 0 \text{ GeV}^2$
⇒ quasi-real photon



How do we do that?

Cross section for P production depends only on $F(q_1^2, q_2^2)$

With the Single Tag Method: $F(q_1^2, q_2^2) \rightarrow F(Q^2)$



$$F(Q^2) = \int T_H(x, Q^2) \Phi_P(x, \mu_F) dx$$

$$T_H(\gamma^* \gamma \rightarrow q\bar{q}) \quad \Phi_P(q\bar{q} \rightarrow P)$$

- μ_F is scale between soft and hard
- x -dependence of $\Phi_P(x, Q^2)$
not known from first principles \rightarrow model
- Experimental data on $F(Q^2)$ is needed
- $F(0)$ from χ PT
- $F(Q^2) \sim Q^{-2}$ at high Q^2 from pQCD

convolution of perturbative and non-perturbative regimes

Our proposal use Padé Approximants

$$F_{P\gamma^*\gamma}(Q^2, 0) = a_0^P \left(1 + b_P \frac{Q^2}{m_P^2} + c_P \frac{Q^4}{m_P^4} + \dots \right)$$

$\Gamma_{P \rightarrow \gamma\gamma}$ slope curvature

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$\Gamma_{P \rightarrow \gamma\gamma}$ slope curvature

We have published space-like data for $Q^2 F_{P\gamma^*\gamma}(Q^2, 0)$

$$Q^2 F_{P\gamma^*\gamma}(Q^2, 0) = a_0 Q^2 + a_1 Q^4 + a_2 Q^6 + \dots$$

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Use Padé Approximants instead (better convergence properties)

$$P_M^N(Q^2) = \frac{T_N(Q^2)}{R_M(Q^2)} = a_0 Q^2 + a_1 Q^4 + a_2 Q^6 + \dots + \mathcal{O}((Q^2)^{N+M+1})$$

Our proposal use Padé Approximants

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$$P_1^1(Q^2) = \frac{a_0 Q^2}{1 - \frac{a_1}{a_0} Q^2}$$

Our proposal use Padé Approximants

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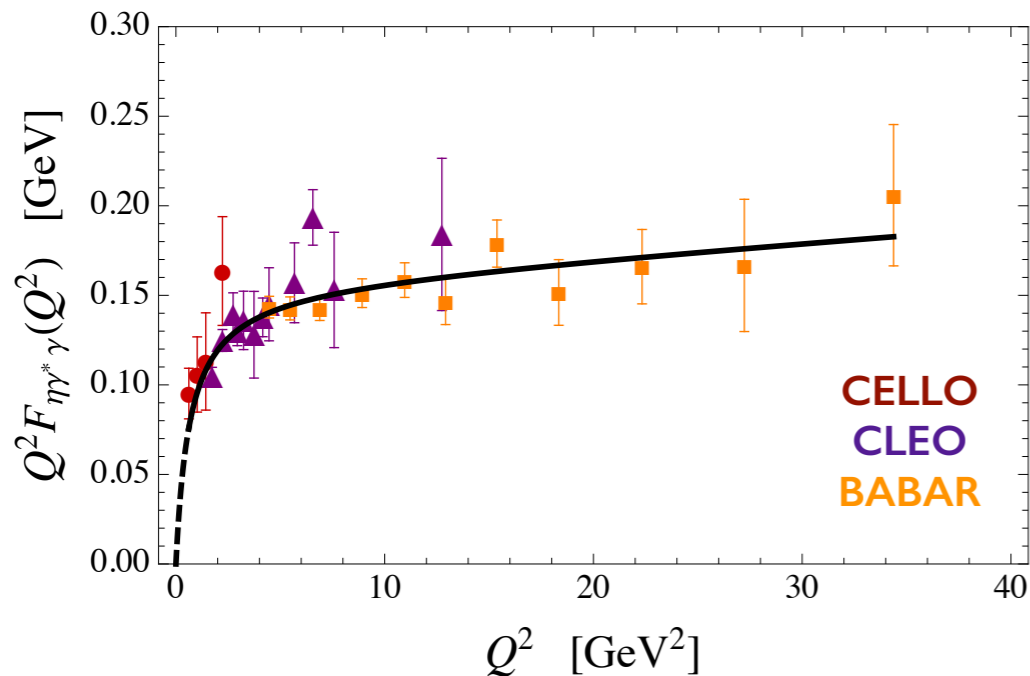
Use Padé Approximants instead (better convergence properties)

$$P_1^1(Q^2) = \frac{a_0 Q^2}{1 - \frac{a_1}{a_0} Q^2} \longrightarrow \begin{aligned} P_1^N(Q^2) &= P_1^1(Q^2), P_1^2(Q^2), P_1^3(Q^2), \dots \\ P_N^N(Q^2) &= P_1^1(Q^2), P_2^2(Q^2), P_3^3(Q^2), \dots \end{aligned}$$

η -TFF

Fit to Space-like data: CELLO'91, CLEO'98, BABAR'11

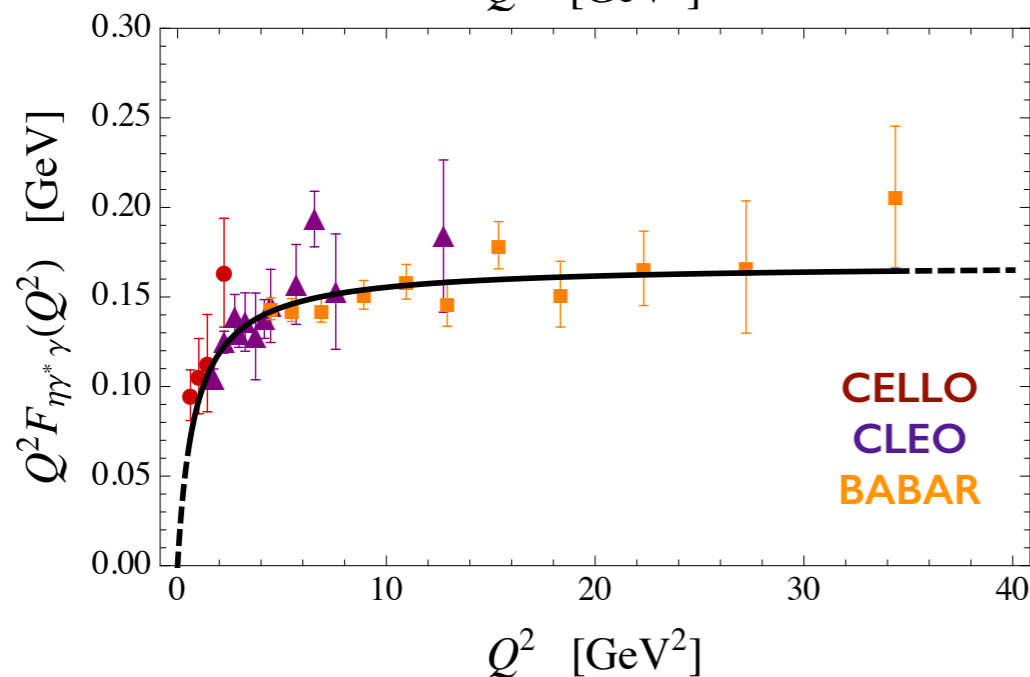
[Phys.Rev. D89 (2014) 034014]



$$P_1^N(Q^2) \quad \text{up to } N=2$$

$$\Gamma_{\eta \rightarrow \gamma\gamma}^{pred} = (0.38 \pm 0.17) \text{ keV}$$

$$\Gamma_{\eta \rightarrow \gamma\gamma}^{PDG} = (0.516 \pm 0.018) \text{ keV}$$



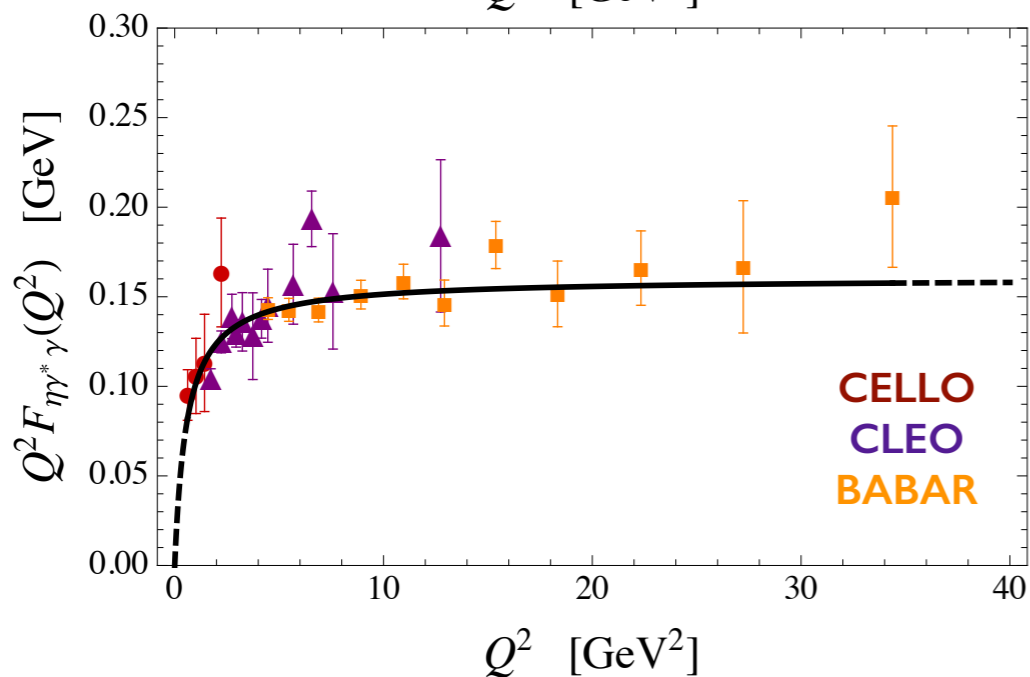
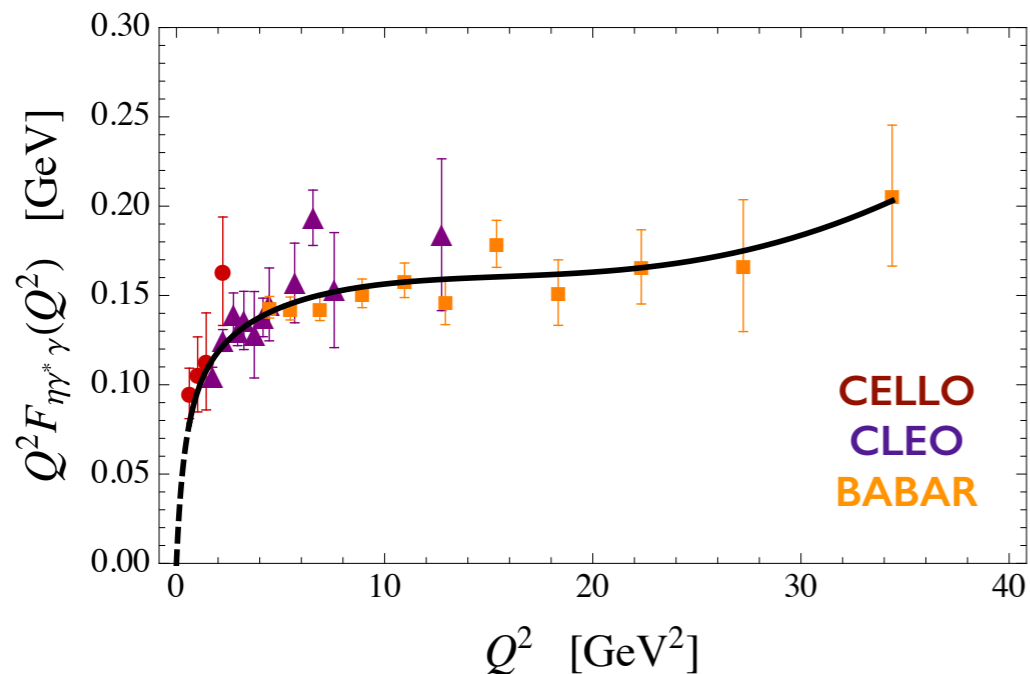
$$P_N^N(Q^2) \quad \text{up to } N=1$$

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma^*\gamma}(Q^2, 0) = 0.17(6) \text{ GeV}$$

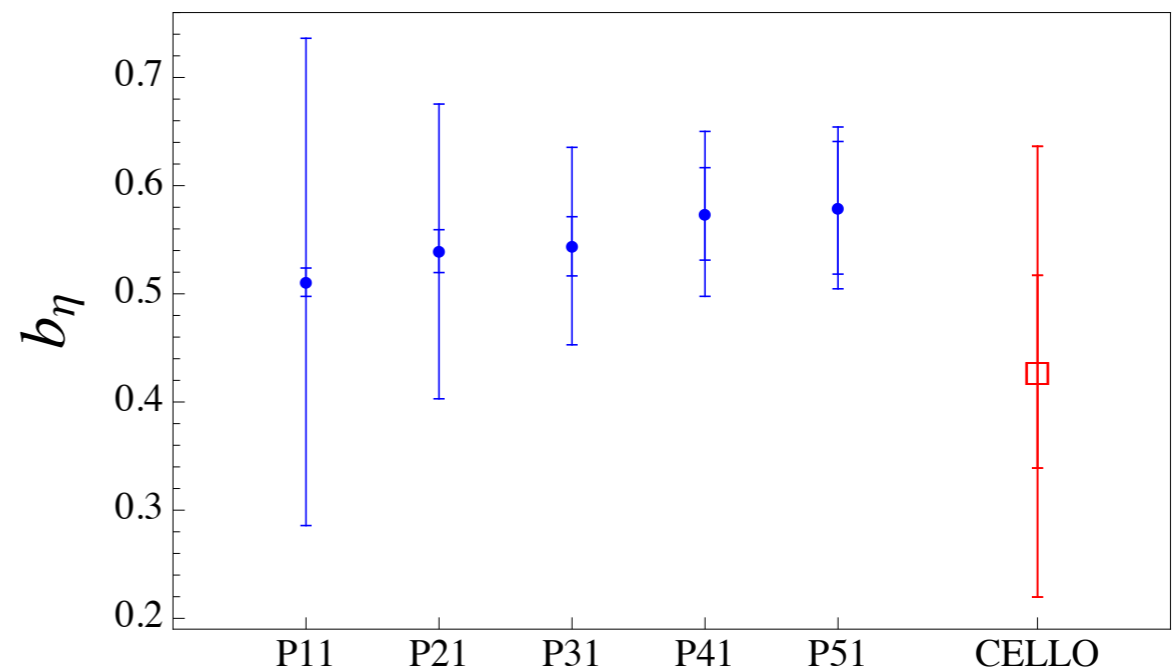
η -TFF

Fit to Space-like data: CELLO'91, CLEO'98, BABAR'11 + $\Gamma_{\eta \rightarrow \gamma\gamma}$

[Phys.Rev. D89 (2014) 034014]



$P_1^N(Q^2)$ up to N=5



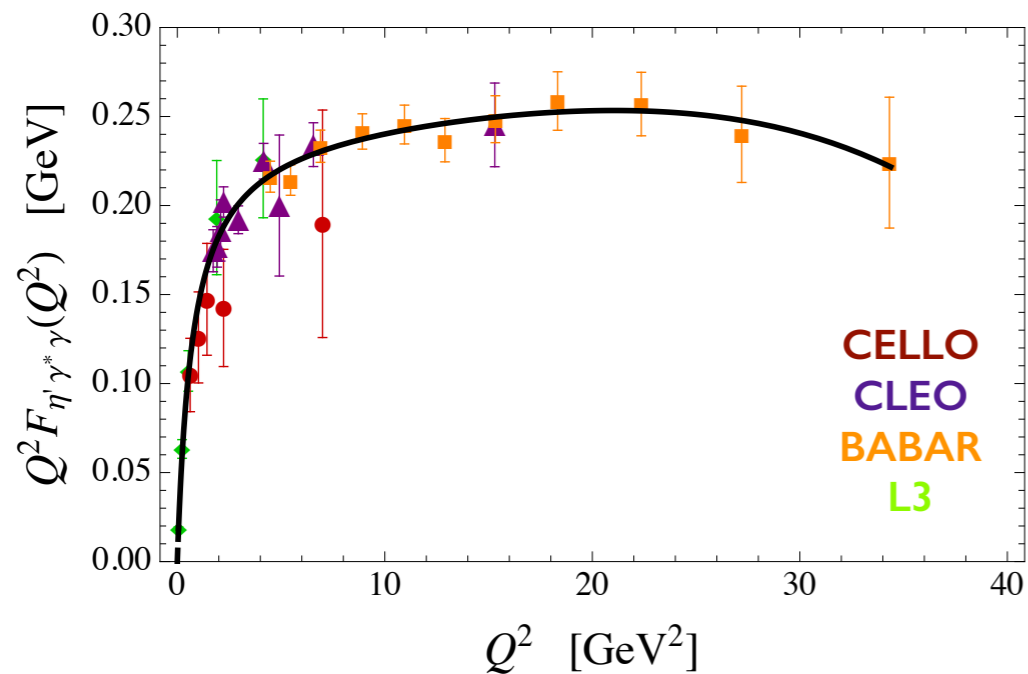
$P_N^N(Q^2)$ up to N=2

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma^*\gamma}(Q^2, 0) = 0.160(24) \text{ GeV}$$

η' -TFF

Fit to Space-like data: CELLO'91, CLEO'98, L3'98, BABAR'11

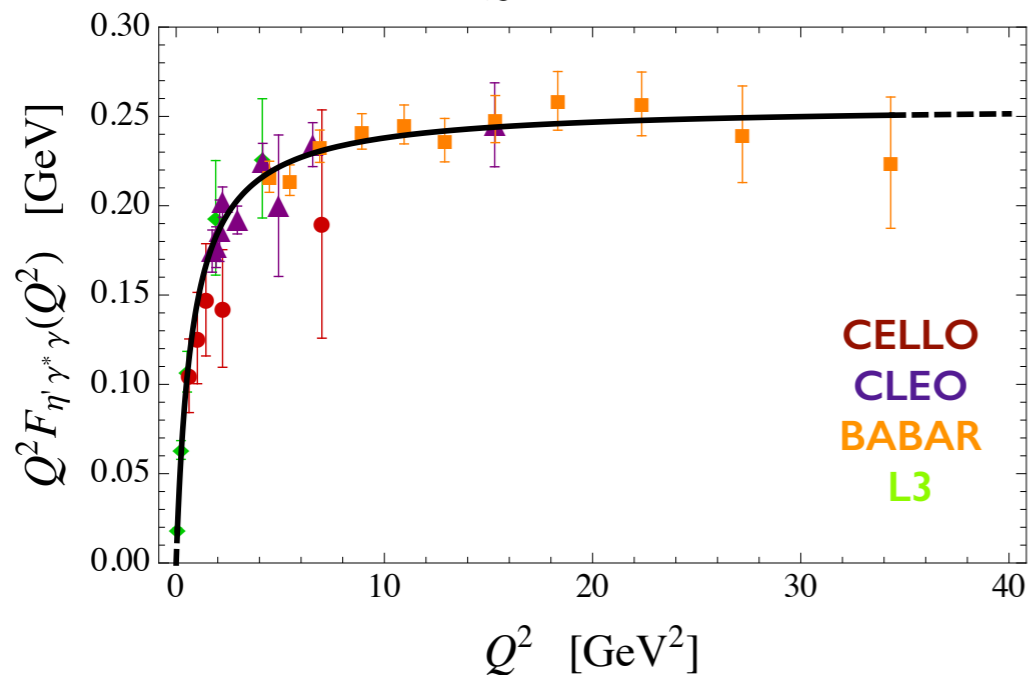
[Phys.Rev. D89 (2014) 034014]



$$P_1^N(Q^2) \quad \text{up to } N=5$$

$$\Gamma_{\eta' \rightarrow \gamma\gamma}^{pred} = (4.21 \pm 0.43) \text{ keV}$$

$$\Gamma_{\eta' \rightarrow \gamma\gamma}^{PDG} = (4.34 \pm 0.14) \text{ keV}$$



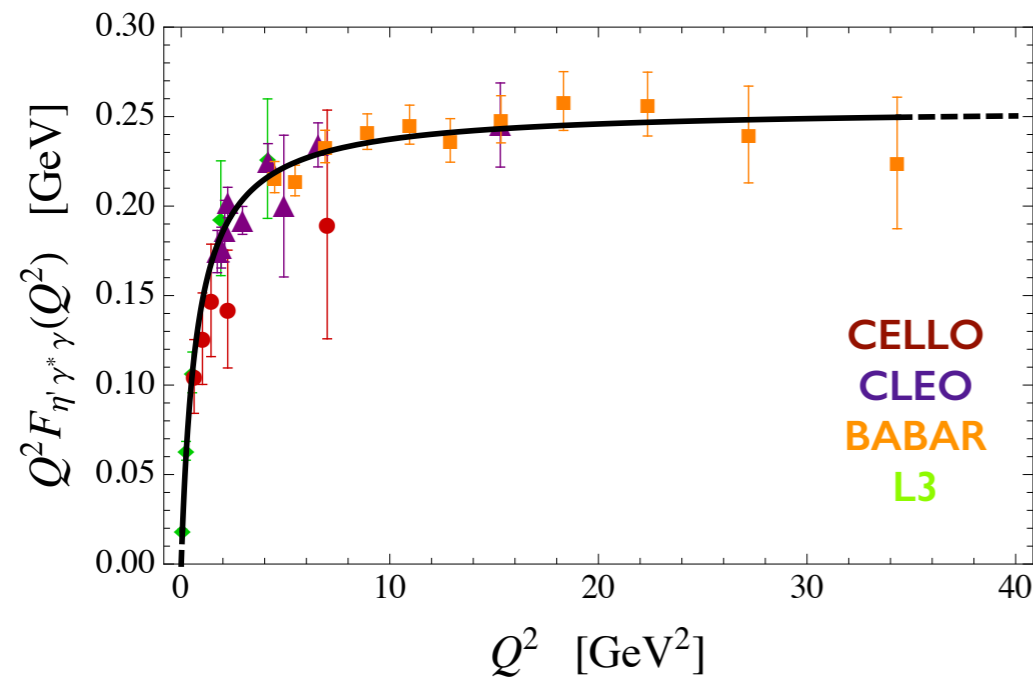
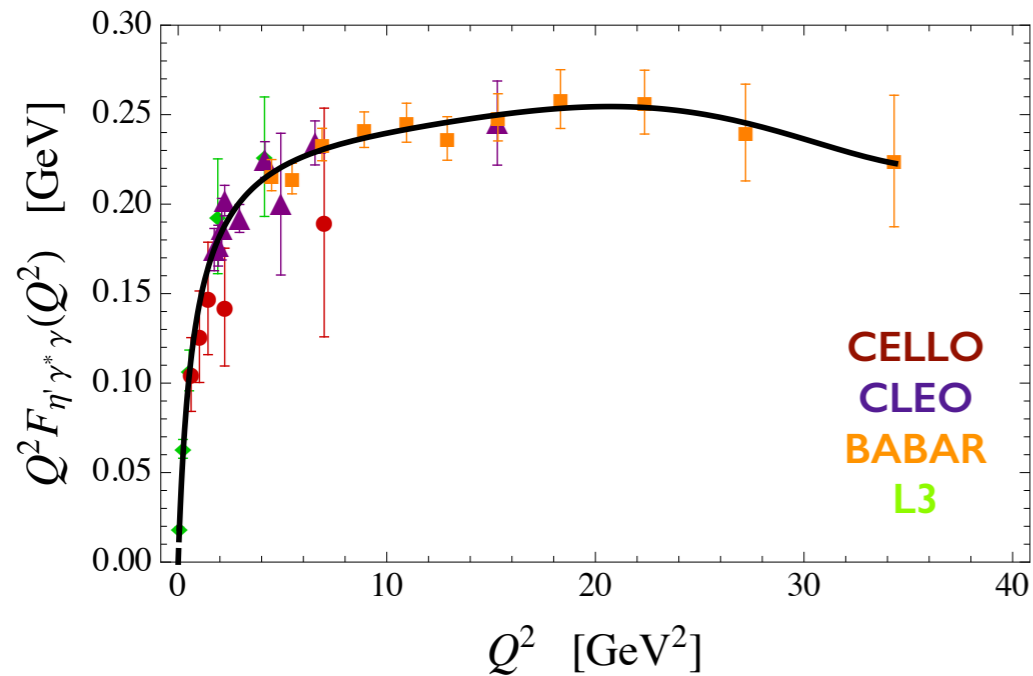
$$P_N^N(Q^2) \quad \text{up to } N=1$$

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta' \gamma^* \gamma}(Q^2, 0) = 0.256(4) \text{ GeV}$$

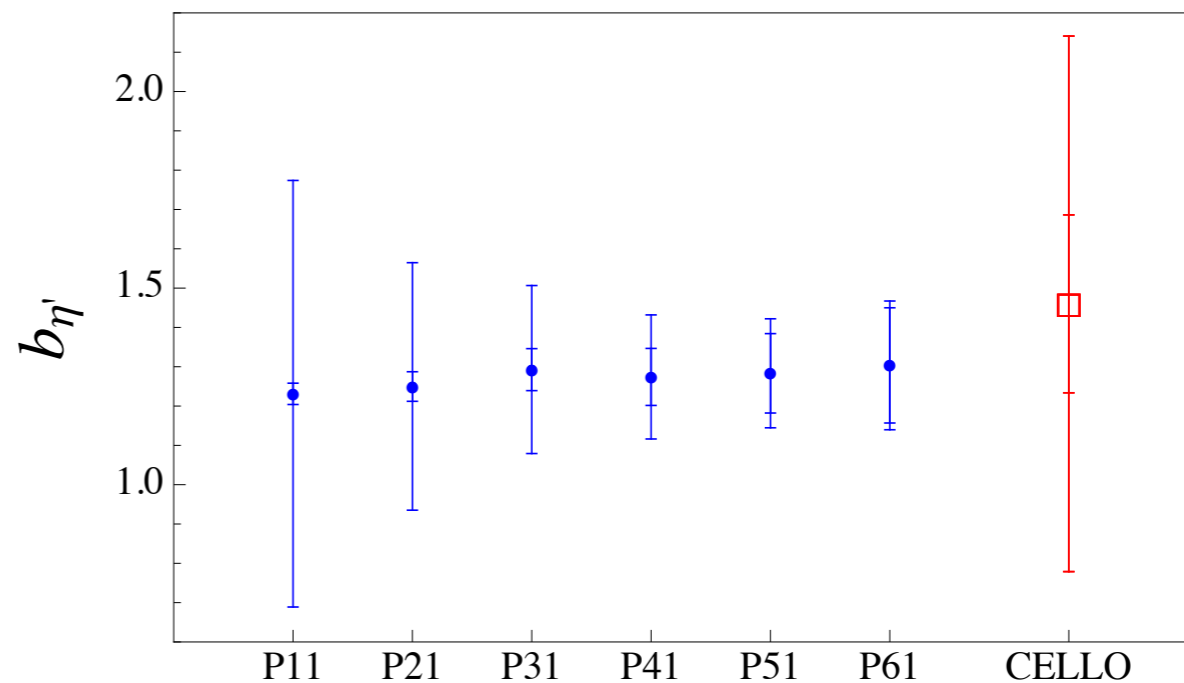
η' -TFF

Fit to Space-like data: CELLO'91, CLEO'98, L3'98, BABAR'11 + $\Gamma_{\eta' \rightarrow \gamma\gamma}$

[Phys.Rev. D89 (2014) 034014]



$P_1^N(Q^2)$ up to N=6



$P_N^N(Q^2)$ up to N=1

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta' \gamma^* \gamma}(Q^2, 0) = 0.255(4) \text{ GeV}$$

η - η' mixing

η - η' mixing in the flavor basis

$$\begin{pmatrix} F_{\eta}^q & F_{\eta}^s \\ F_{\eta'}^q & F_{\eta'}^s \end{pmatrix} = \begin{pmatrix} F_q \cos \phi & -F_s \sin \phi \\ F_q \sin \phi & F_s \cos \phi \end{pmatrix}$$

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From the TFFs we can determine F_q, F_s, ϕ

$$F_{\eta\gamma\gamma}(0) = \frac{1}{4\pi^2} \left(\frac{\hat{c}_q}{F_q} \cos \phi - \frac{\hat{c}_s}{F_s} \sin \phi \right) \quad \lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma^*\gamma}(Q^2) = 2(\hat{c}_q F_q \cos \phi - \hat{c}_s F_s \sin \phi)$$
$$F_{\eta'\gamma\gamma}(0) = \frac{1}{4\pi^2} \left(\frac{\hat{c}_q}{F_q} \sin \phi + \frac{\hat{c}_s}{F_s} \cos \phi \right) \quad \lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta'\gamma^*\gamma}(Q^2) = 2(\hat{c}_q F_q \sin \phi + \hat{c}_s F_s \cos \phi)$$

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[Phys.Rev. D89 (2014) 034014]

$$F_q = 1.06(1)F_{\pi}, \quad F_s = 1.56(24)F_{\pi}, \quad \phi = 40.3(1.8)^{\circ}$$

Update of Frere-Escribano '05 with PDG12 using 9 inputs

$$F_q = 1.07(1)F_{\pi}, \quad F_s = 1.63(3)F_{\pi}, \quad \phi = 39.6(0.4)^{\circ}$$

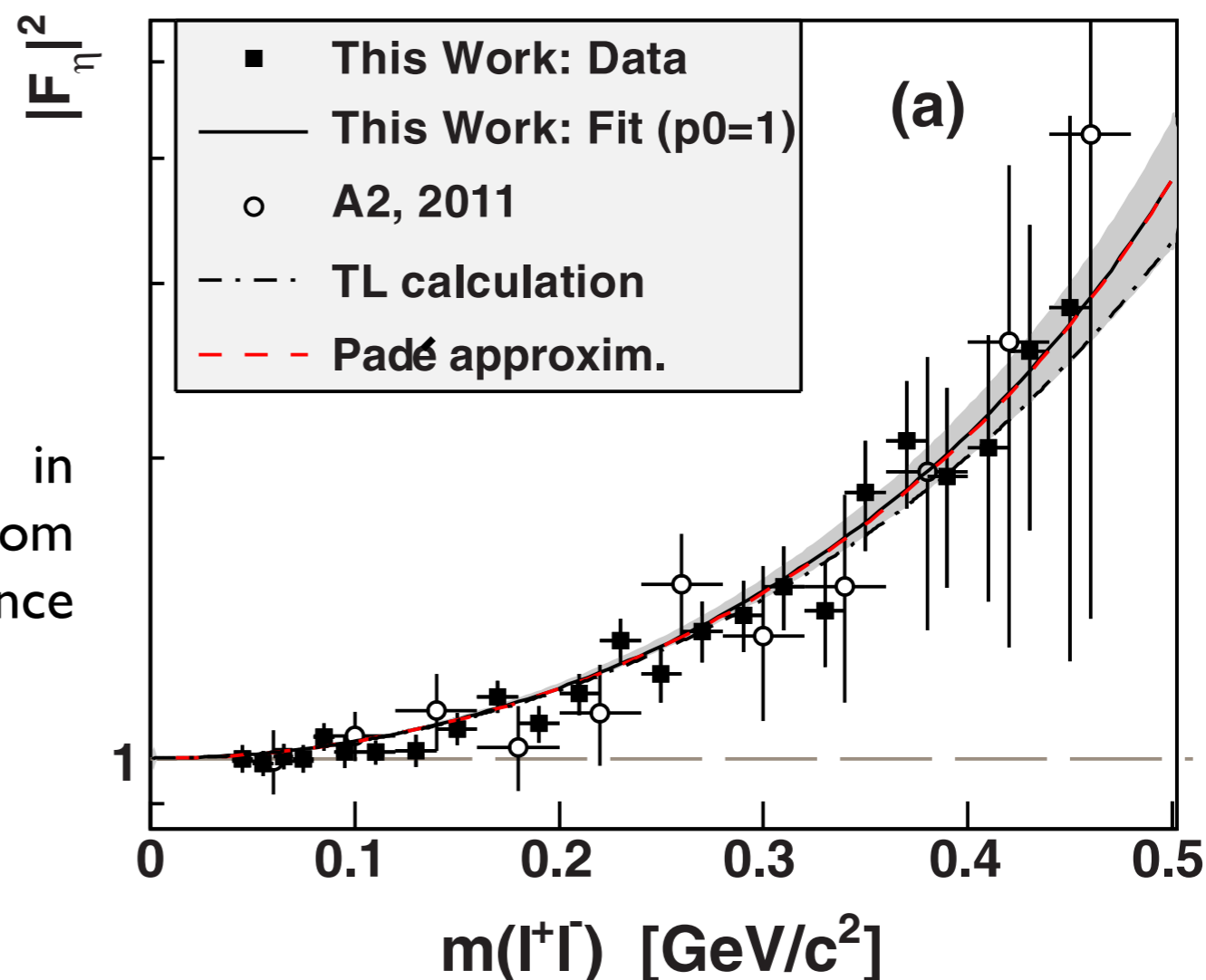
Pseudoscalar Transition Form Factors

- Study Dalitz decays

$$\eta \rightarrow \gamma^* \gamma \rightarrow e^+ e^- \gamma$$

- Our predictions from SL data are in excellent agreement with latest results from A2 Coll. as expected below resonance region.

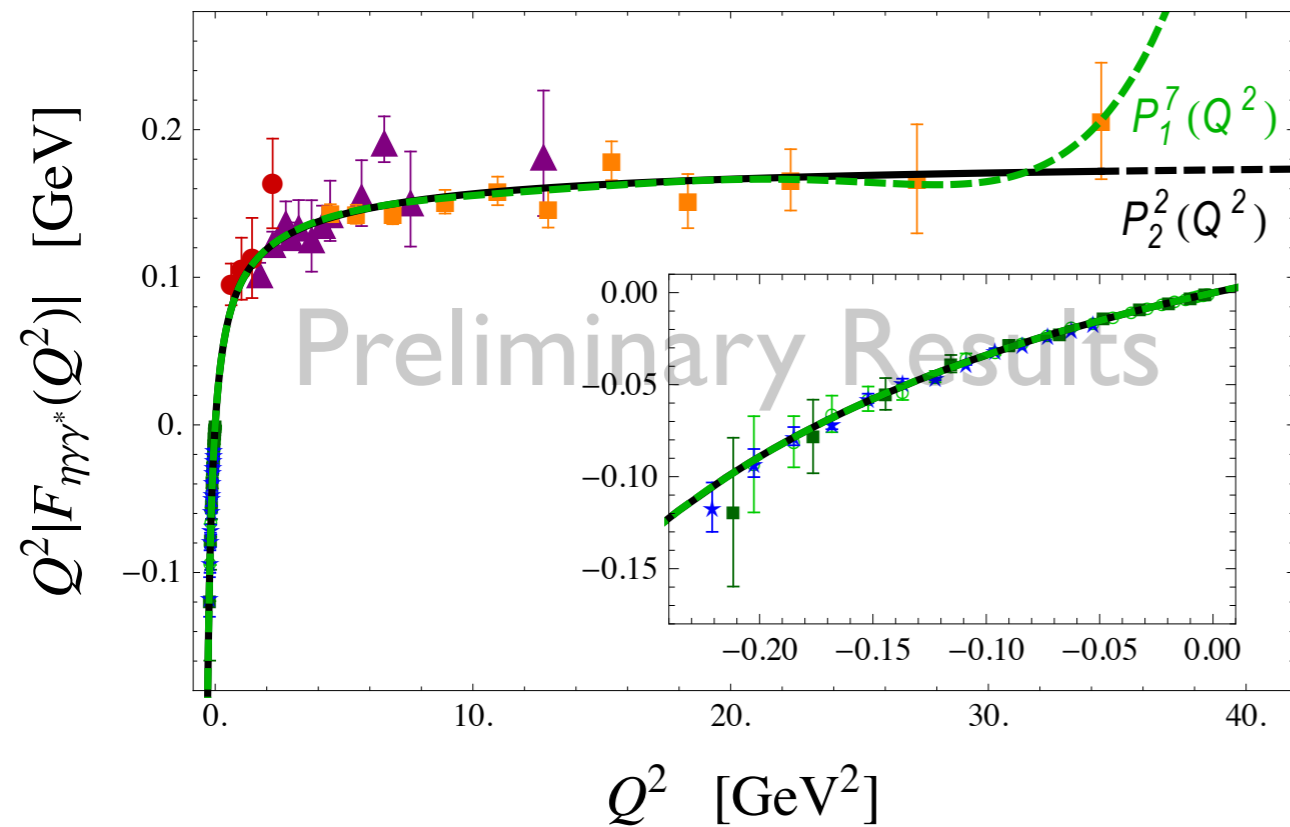
Result from A2 Coll. [Phys.Rev. C89 (2014) 044608]



η -TFF

Fit to Space-like data [CELLO'91, CLEO'98, BABAR'11] + $\Gamma_{\eta \rightarrow \gamma\gamma}$
 + Time-like data [NA60'09, A2'11, A2'13]

[R.Escribano, P.M., P. Sanchez-Puertas, '14]



$P_1^N(Q^2)$ up to $N=7$

$$\Gamma_{\eta \rightarrow \gamma\gamma}^{pred} = (0.42 \pm 0.10) keV$$

$$\Gamma_{\eta \rightarrow \gamma\gamma}^{PDG} = (0.518 \pm 0.018) keV$$

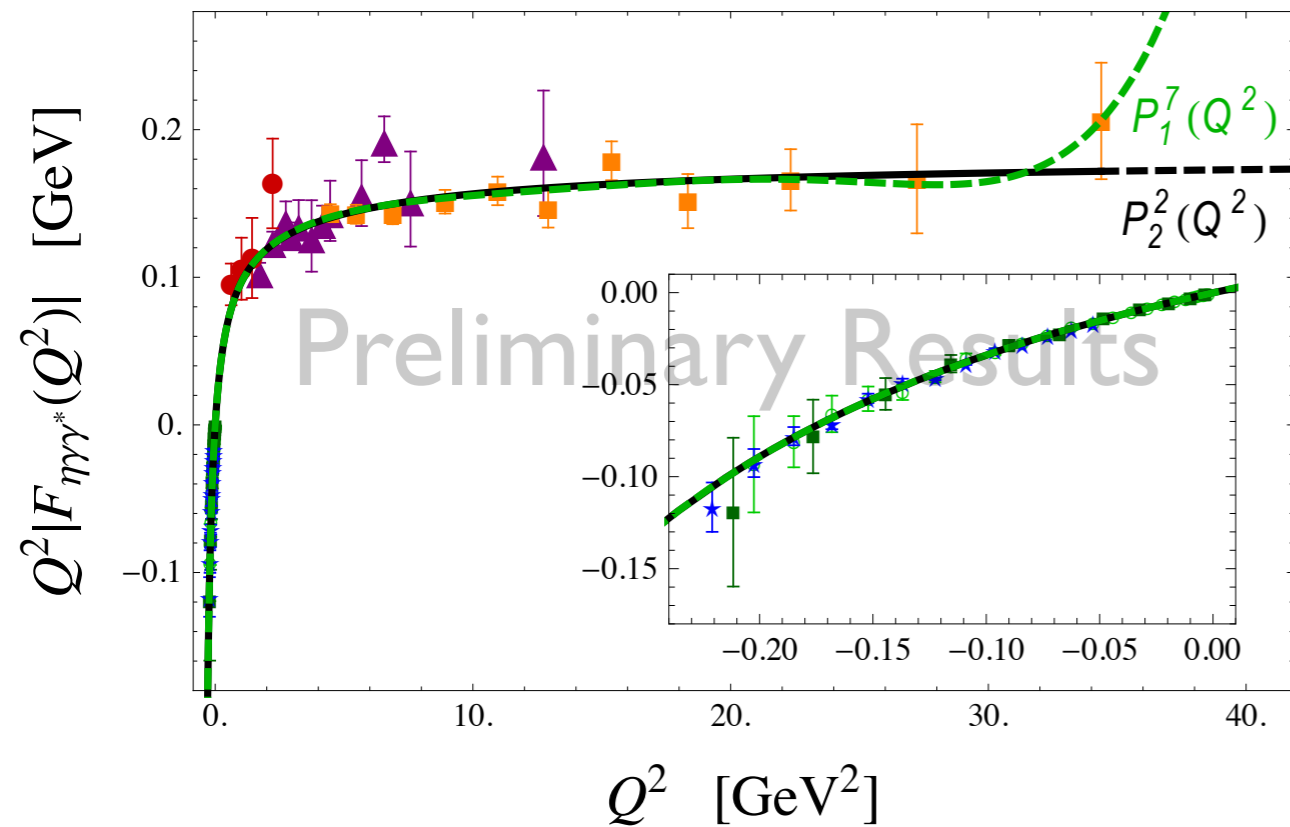
$P_N^N(Q^2)$ up to $N=2$

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma^*}(Q^2, 0) = 0.169(14) GeV$$

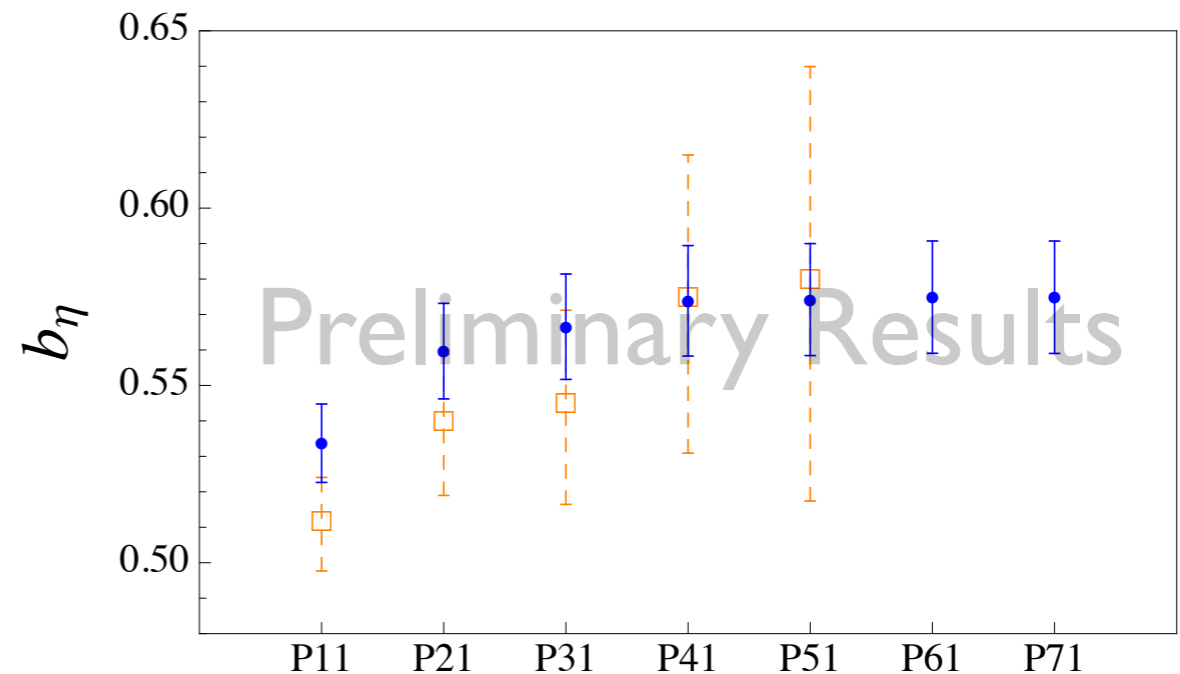
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[R.Escribano, P.M., P. Sanchez-Puertas, '14]



$P_1^N(Q^2)$ up to N=7



$P_N^N(Q^2)$ up to N=2

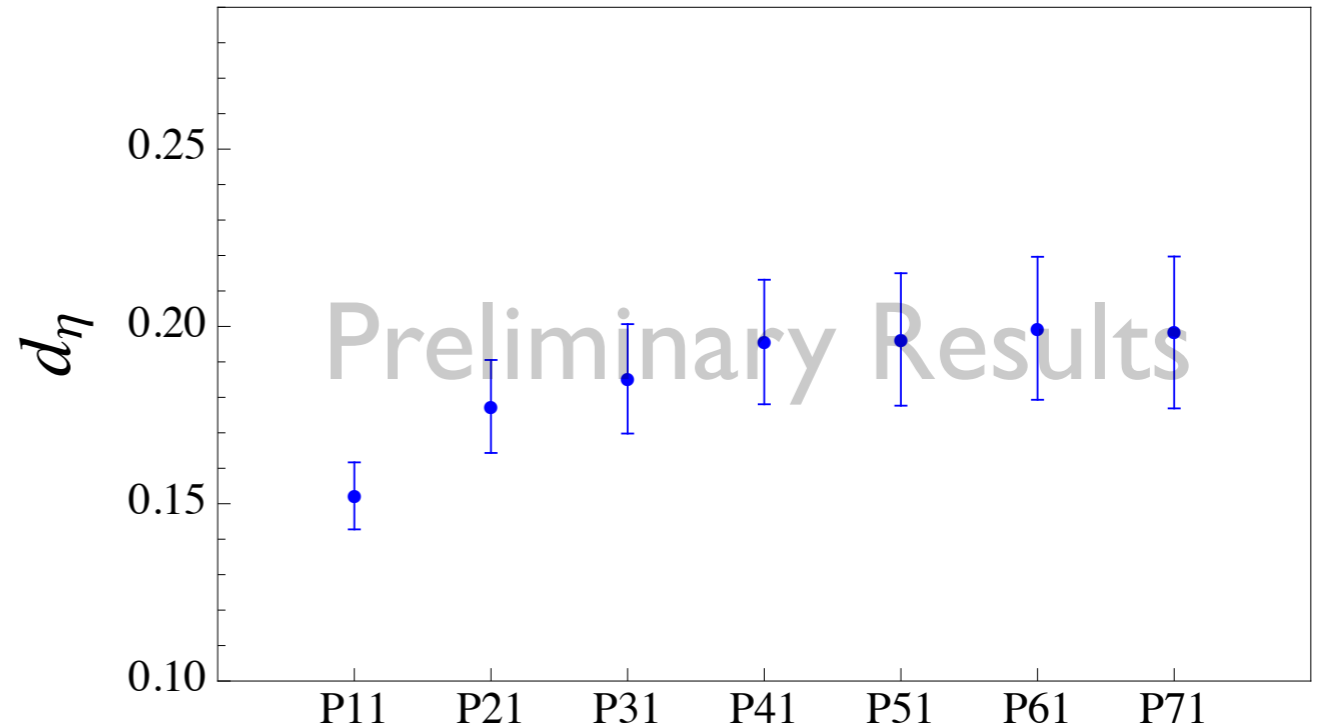
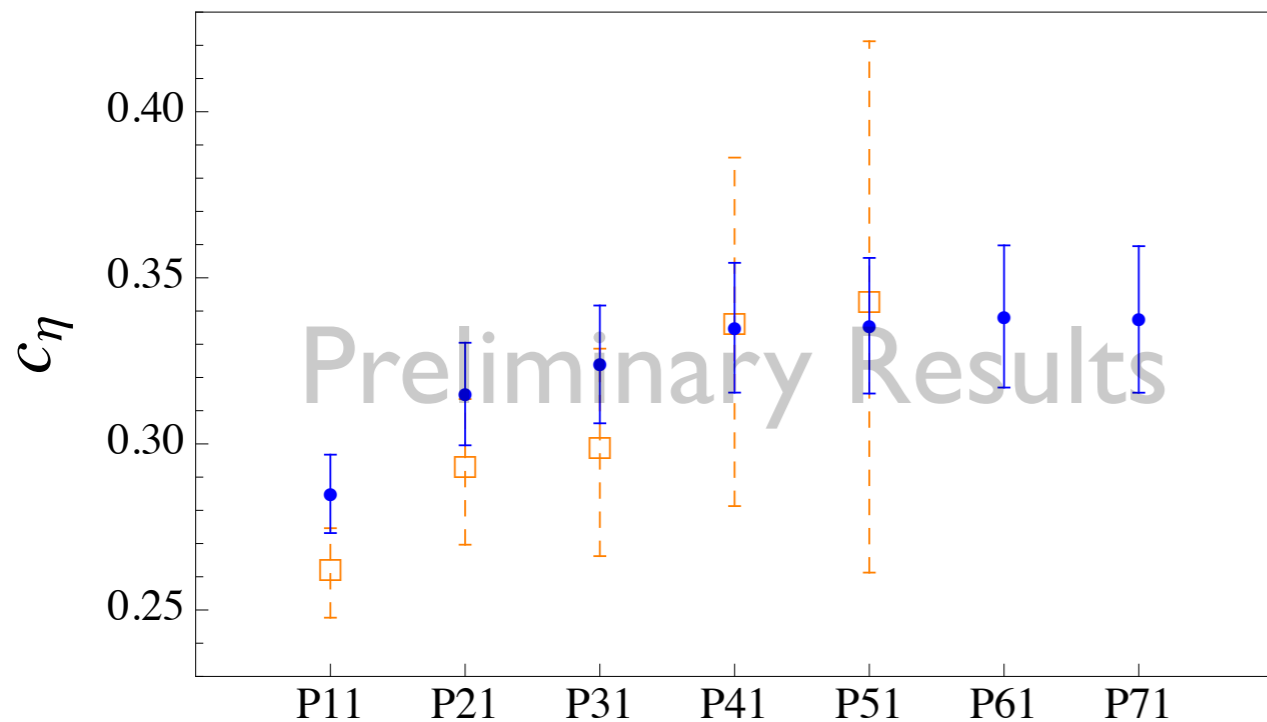
$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma^*\gamma}(Q^2, 0) = 0.177(15) \text{ GeV}$$

η -TFF

Fit to Space-like data [CELLO'91, CLEO'98, BABAR'11] + $\Gamma_{\eta \rightarrow \gamma\gamma}$
+ Time-like data [NA60'09, A2'11, A2'13]

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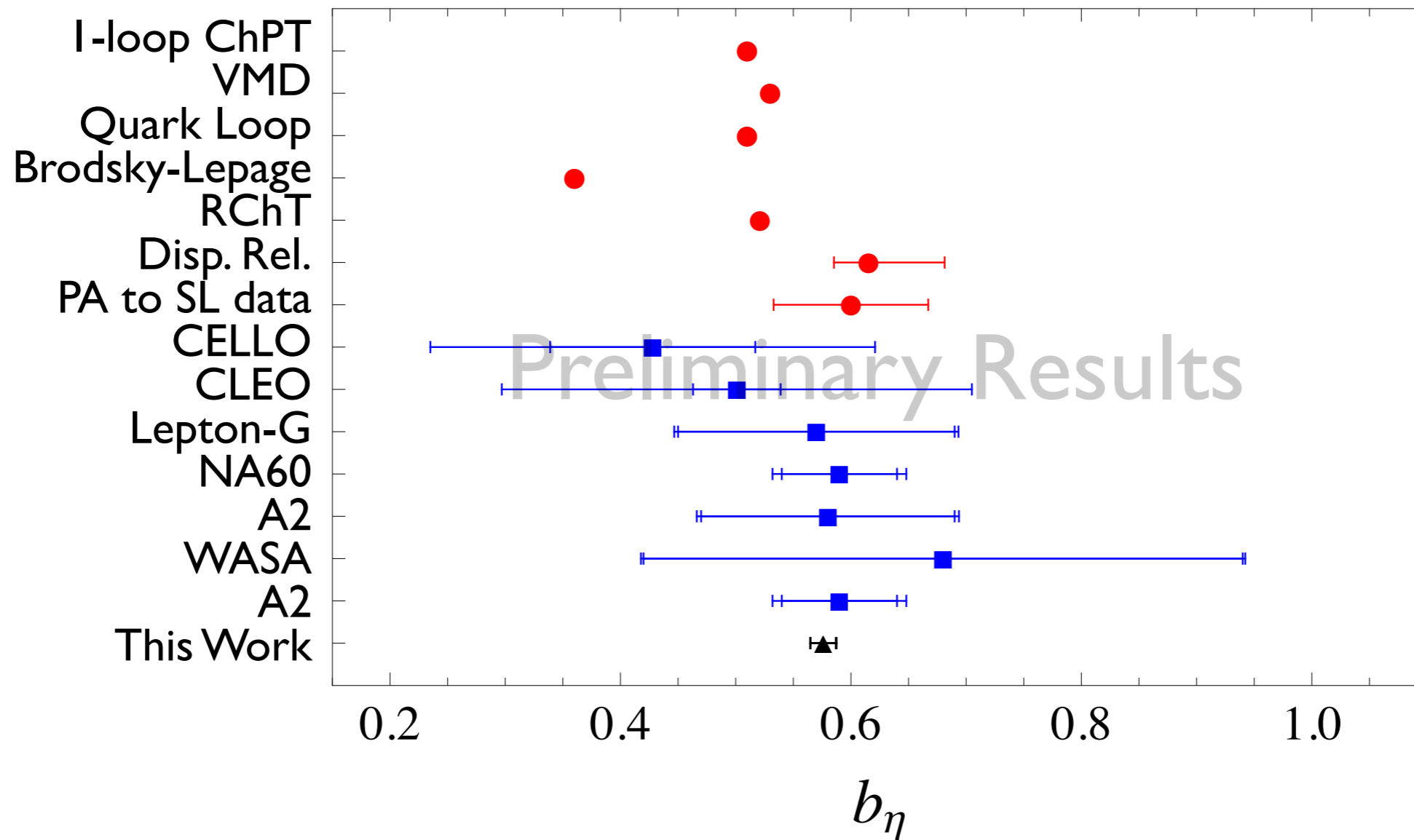
[R. Escribano, P.M., P. Sanchez-Puertas, '14]



slope for the η -TFF

Fit to Space-like data [CELLO'91, CLEO'98, BABAR'11] + $\Gamma_{\eta \rightarrow \gamma\gamma}$
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[R.Escribano, P.M., P. Sanchez-Puertas, '14]



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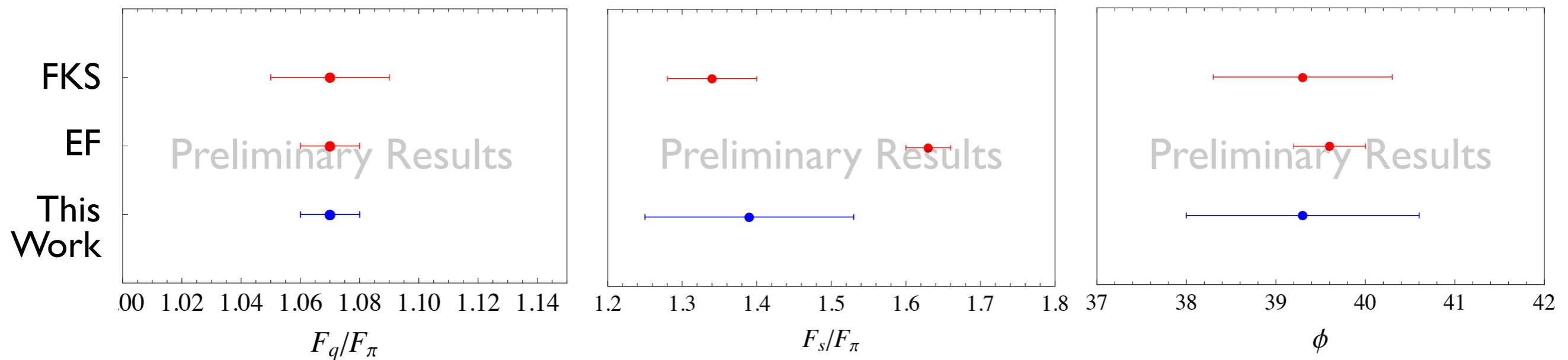
Update of Frere-Escribano '05 with PDG12 using 9 inputs

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η - η' mixing

η - η' mixing in the flavor basis

From the TFFs we can determine F_q, F_s, ϕ



FKS: Feldmann, Kroll, Stech, PRD 58, 114006 (1998)

EF: Escribano, Frere, JHEP 0506, 029 (2005) updated in Escribano, P.M, Sanchez-Puertas, 2013.

Conclusions

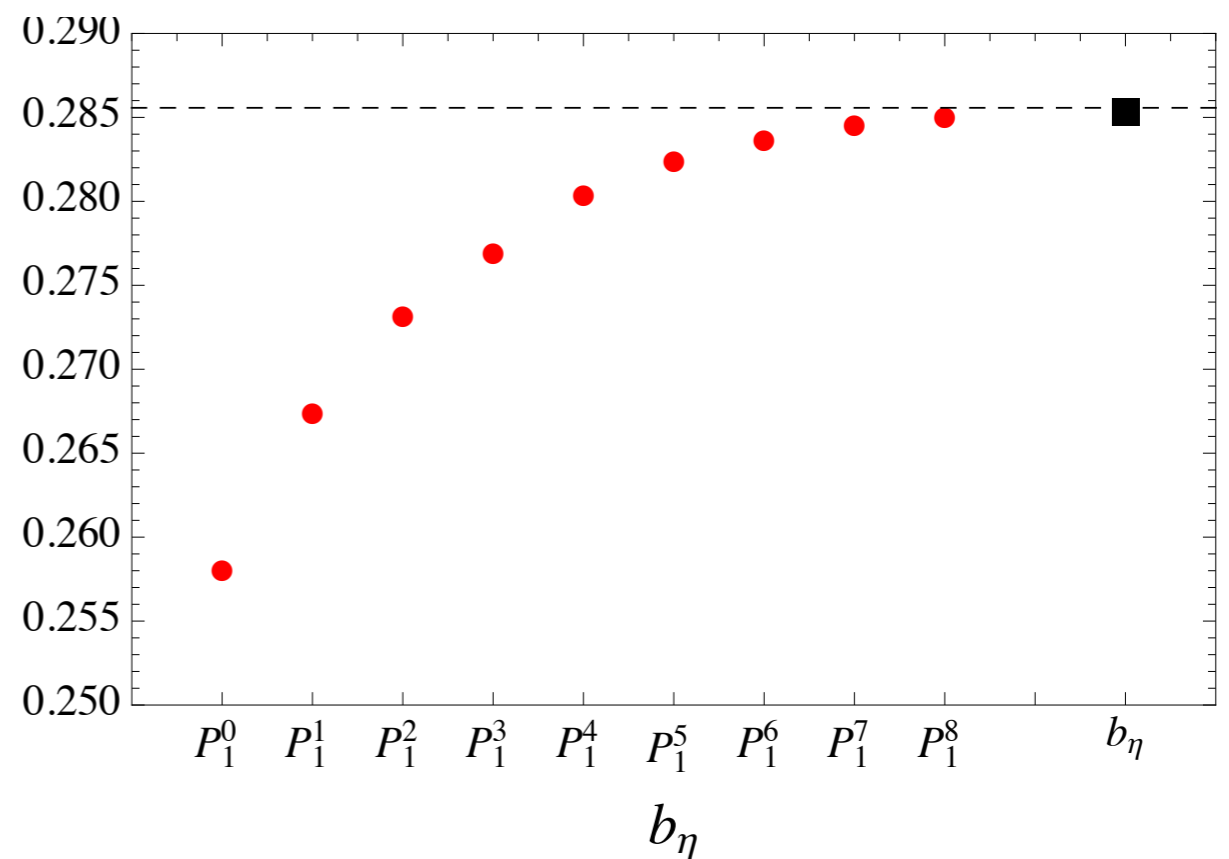
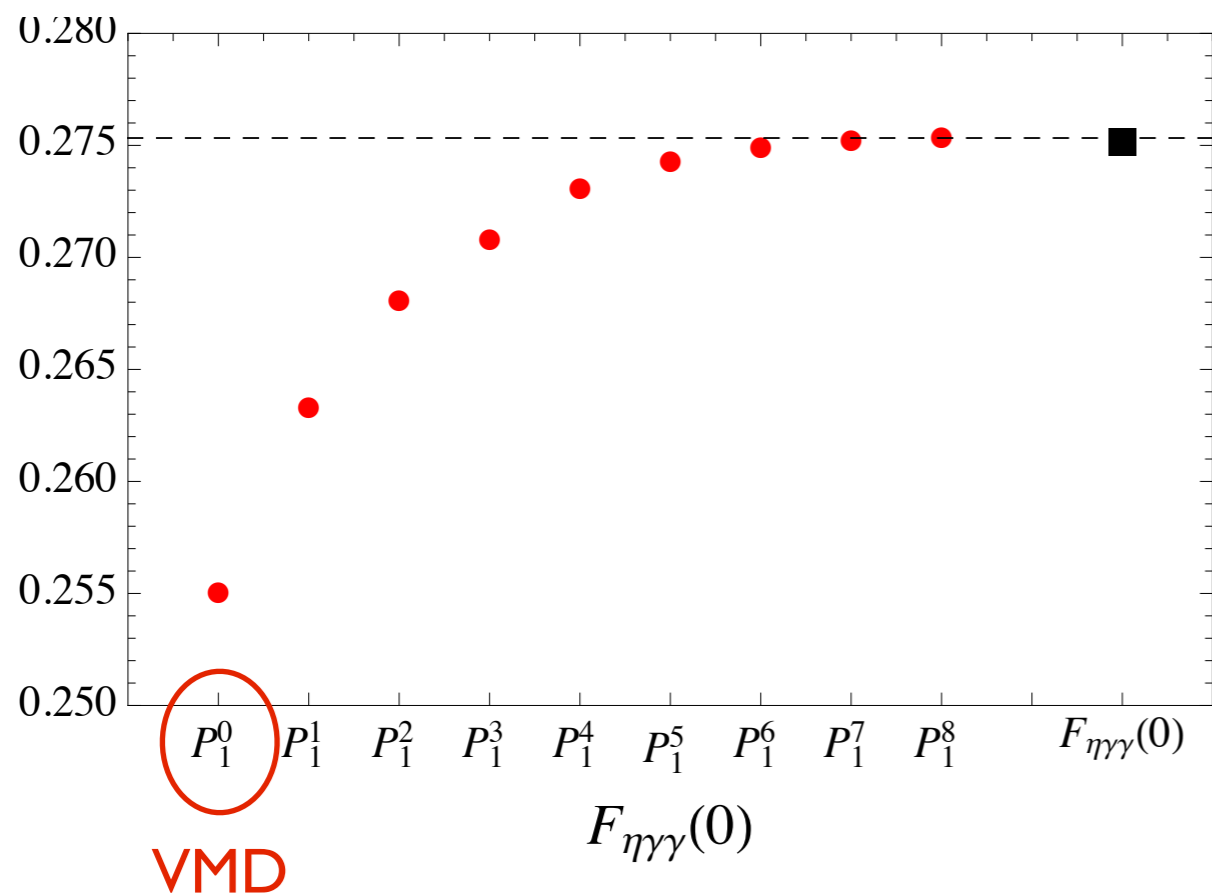
- Transition Form Factors are a good laboratory to study meson properties
- Need for a model independent approach: we use Padé App.
- Considering Space-like and time-like data
 - provides very accurate LECs and asymptotic limits
 - provides insight in mixing scheme and meson structure
- Many more applications
 - Predicts $VP\gamma$, J/Ψ , rare decays ($P \rightarrow e^+e^-(\mu^+\mu^-)$), $(g-2)_\mu$...

Padé Approximants' method is easy, systematic, deals with syst. errors and can be improved upon by including new data

Thank you!

A word on systematics

- Consider a model for η TFF
- Generate a pseudodata set emulating the physical situation (SL+TL)
- Build up your PA sequence
- Fit and compare



η - η' mixing

From the TFFs we can determine F_q, F_s, ϕ

and the $V\rho\gamma$ and J/Ψ decays used in FKS and EF as inputs

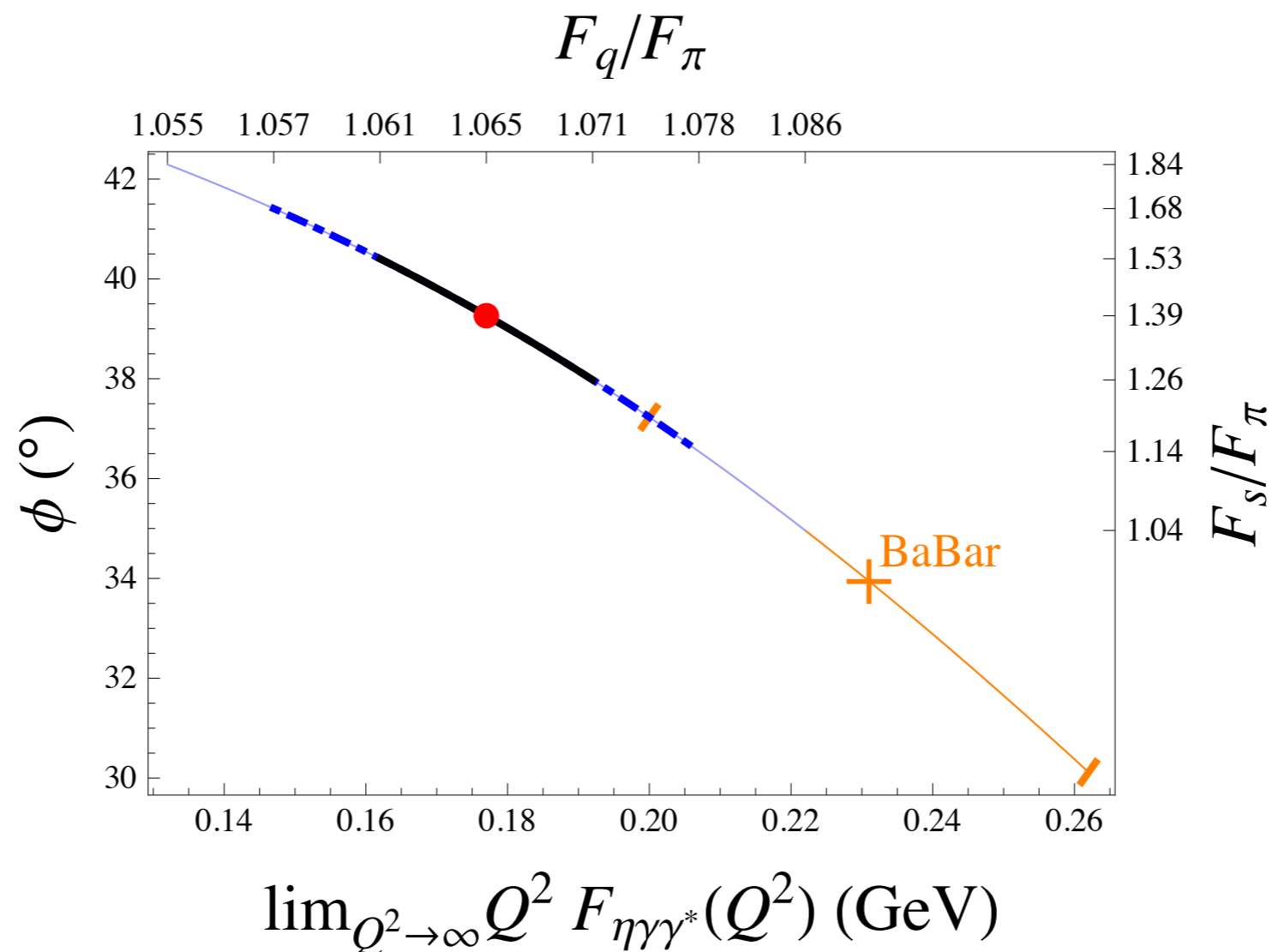
(using $F_{\pi^0} = 131.5 \pm 1.4$ MeV instead of $F_{\pi^-} = 92.21 \pm 0.14$ MeV)

| | Our predictions | Experimental determinations |
|--|-----------------|-----------------------------|
| $\mathcal{G}_{\rho\eta\gamma}$ | 1.46(3) | 1.58(5) |
| $\mathcal{G}_{\rho\eta'\gamma}$ | 1.20(4) | 1.32(3) |
| $\mathcal{G}_{\omega\eta\gamma}$ | 0.56(2) | 0.45(2) |
| $\mathcal{G}_{\omega\eta'\gamma}$ | 0.55(2) | 0.43(2) |
| $\mathcal{G}_{\phi\eta\gamma}$ | -0.78(8) | -0.69(1) |
| $\mathcal{G}_{\phi\eta'\gamma}$ | 0.88(10) | 0.72(1) |
| $\frac{J/\Psi \rightarrow \eta'\gamma}{J/\Psi \rightarrow \eta\gamma}$ | 5.09(47) | 4.67(20) |

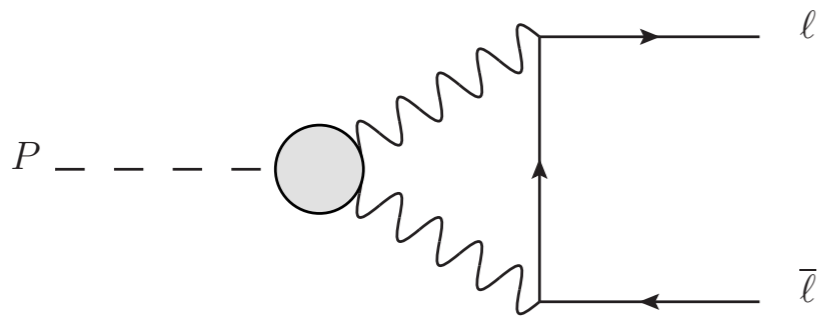
η - η' mixing

η - η' mixing in the flavor basis

From the TFFs we can determine F_q, F_s, ϕ



Dissection of $\eta \rightarrow l^+ l^-$



$$\frac{BR(P \rightarrow \bar{l}l)}{BR(P \rightarrow \gamma\gamma)} = 2 \left(\frac{\alpha m_l}{\pi m_P} \right)^2 \beta_l(m_P^2) |\mathcal{A}(m_P^2)|^2$$

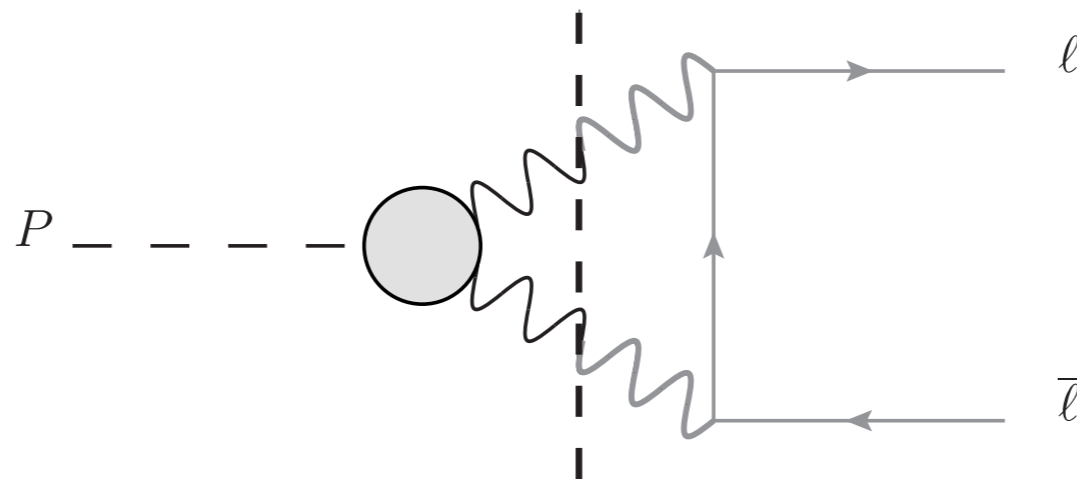
The only unknown $\mathcal{A}(m_P^2)$ from loop calculation where the TFF enters.

$$\mathcal{A}(q^2) = \frac{2i}{\pi^2 q^2} \int d^4 k \frac{(q^2 k^2 - (q \cdot k)^2) F_P(k^2, (q - k)^2)}{(k^2 + i\epsilon)((q - k)^2 + i\epsilon)((p - k)^2 - m^2 + i\epsilon)}$$

Dissection of $\eta \rightarrow l^+ l^-$

As model independent as possible:

Cutkosky rules provides the imaginary part



Dissection of $\eta \rightarrow l^+ l^-$

As model independent as possible:

Cutcosky rules provides the imaginary part

$$\text{Im}\mathcal{A}(q^2) = \frac{\pi}{2\beta_I(q^2)} \ln \left(\frac{1 - \beta_I(q^2)}{1 + \beta_I(q^2)} \right); \quad \beta_I(q^2) = \sqrt{1 - \frac{4m_l^2}{q^2}}$$

$q^2 = m_P^2$

Dissection of $\eta \rightarrow l^+ l^-$

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$q^2 = m_P^2$

Use dispersion relations to get the real part

$$\text{Re}(\mathcal{A}(q^2)) = \mathcal{A}(0) + \frac{1}{\beta_I(q^2)} \left(\frac{\pi^2}{12} + \frac{1}{4} \ln^2 \left(\frac{1 - \beta_I(q^2)}{1 + \beta_I(q^2)} \right) + \text{Li}_2 \left(\frac{1 - \beta_I(q^2)}{1 + \beta_I(q^2)} \right) \right)$$

Dissection of $\eta \rightarrow l^+l^-$

PDG value dominated by the KTeV measurement

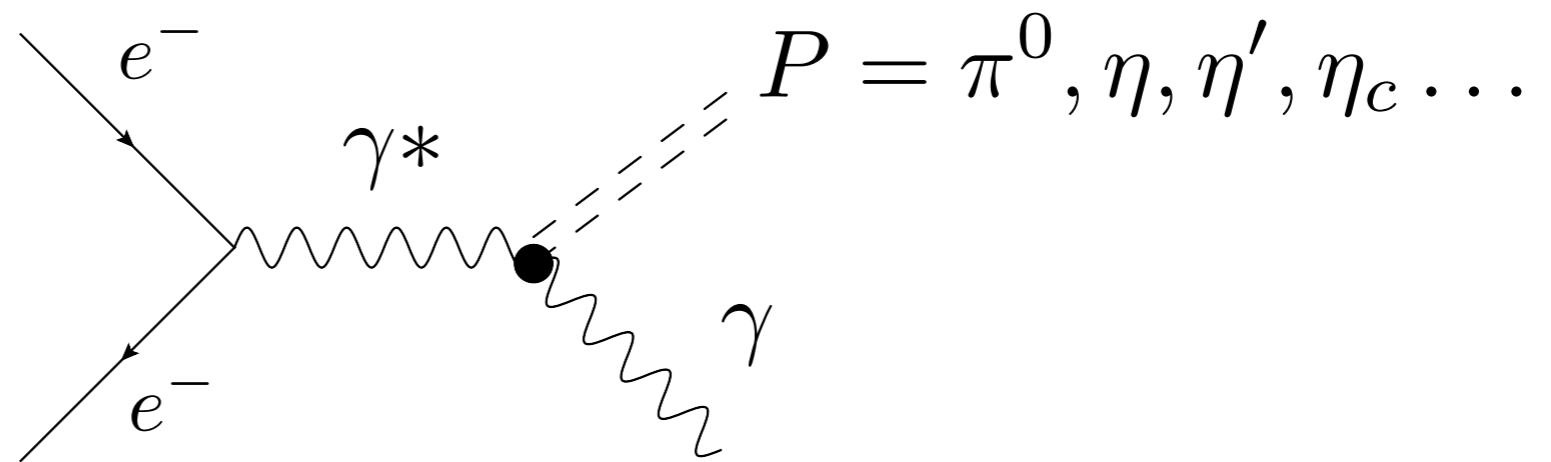
$$\frac{BR(P \rightarrow \bar{l}l)}{BR(P \rightarrow \gamma\gamma)} = 2 \left(\frac{\alpha m_l}{\pi m_P} \right)^2 \beta_l(m_P^2) |\mathcal{A}(m_P^2)|^2 = 5.8(8) \cdot 10^{-6} \quad (\mu^+\mu^-)$$
$$\leq 5.6 \cdot 10^{-6} \quad (e^+e^-)$$

$$\text{Unitary Bound for the } \mu\mu \text{ case} = 4.37 \cdot 10^{-6}$$

$$\text{SM calculations with } m_\eta^2/\Lambda^2 \sim 0 = 4.99 \cdot 10^{-6}$$

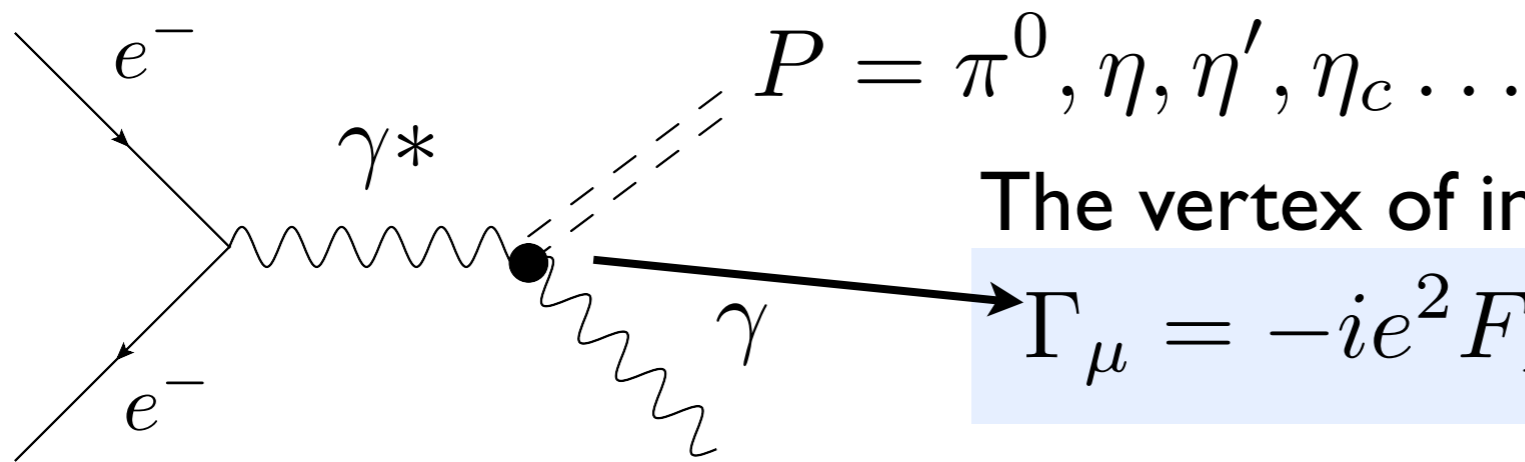
$$\text{Our result from SL+TL (full result)} = 4.51(2) \cdot 10^{-6}$$

Time-like TFF: prediction



- Asymptotic limits in time-like and space-like FFs are expected to be close, is important to measure this time-like FF because:
 - the charmonium region is between the perturbative and non-perturbative regimes of the π -, η -, and η' -TFF
 - background for charmonium decays

Time-like TFF: prediction



The vertex of interest is:

$$\Gamma_\mu = -ie^2 F_P(Q^2) \epsilon_{\mu\nu\rho\sigma} p^\nu \epsilon^\rho q^\sigma$$

Differential cross section:

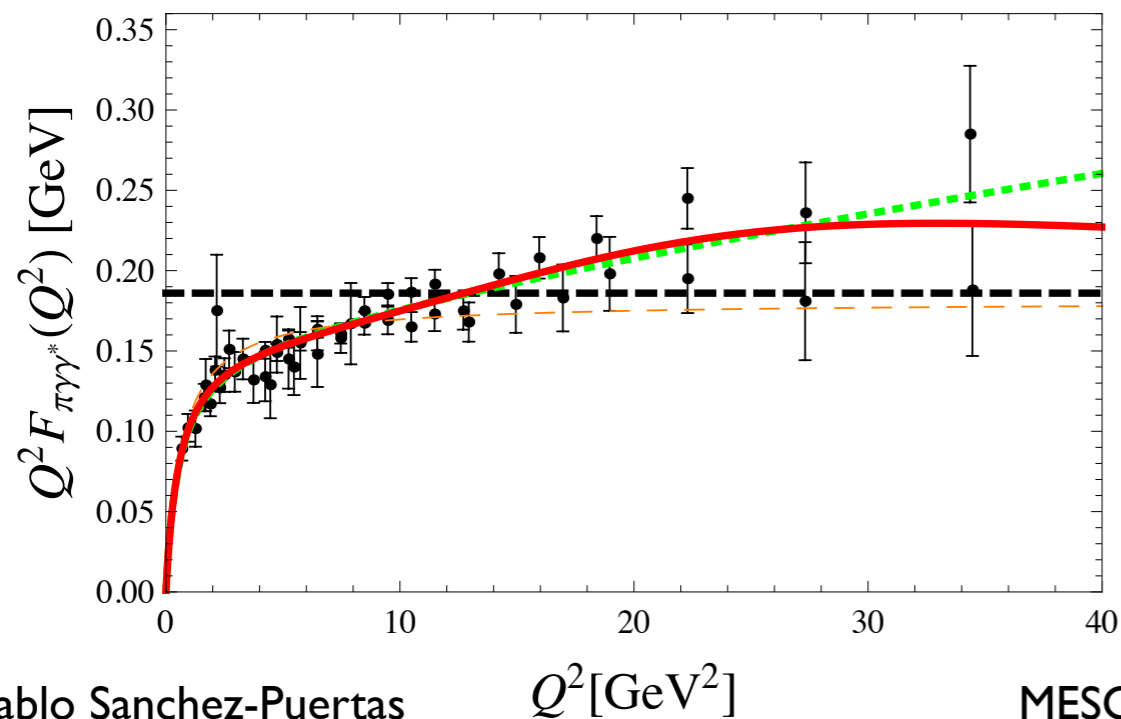
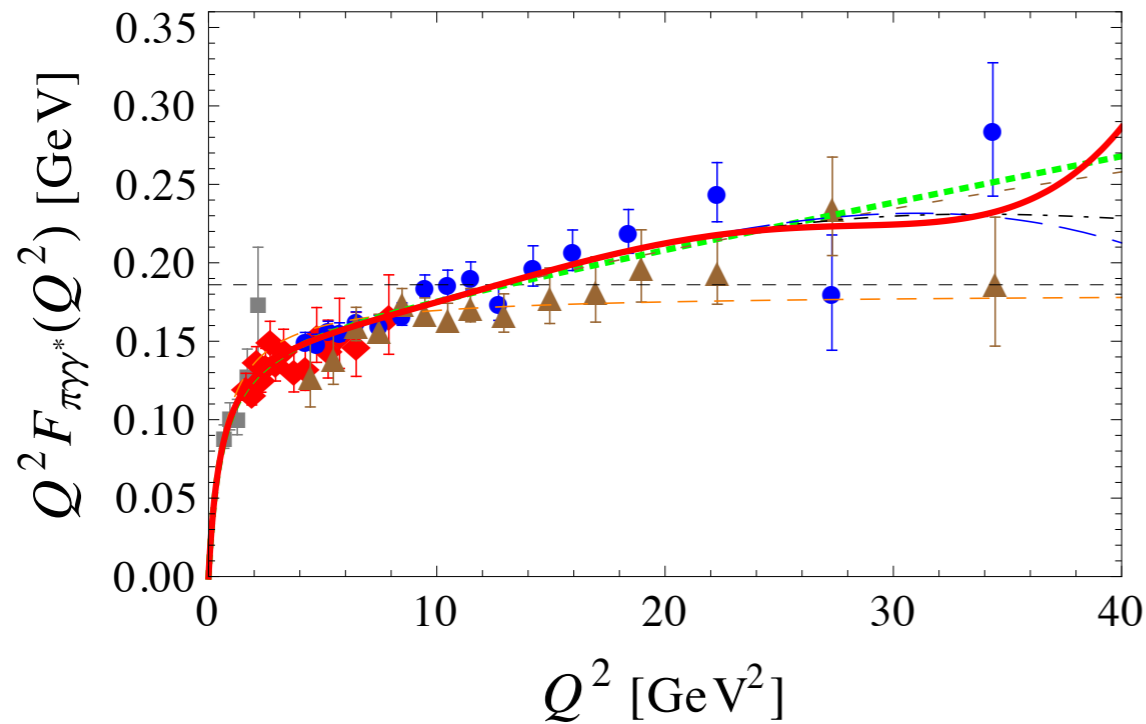
$$\frac{d\sigma(e^+e^- \rightarrow \gamma^* \rightarrow \gamma P)}{d(\cos\theta)} = \frac{\pi^2 \alpha^3}{4} (F_{P\gamma^*\gamma}(s, 0))^2 \left(1 - \frac{M_P^2}{s}\right)^3 (1 + \cos^2\theta)$$

Integrating with respect to $\cos\theta$

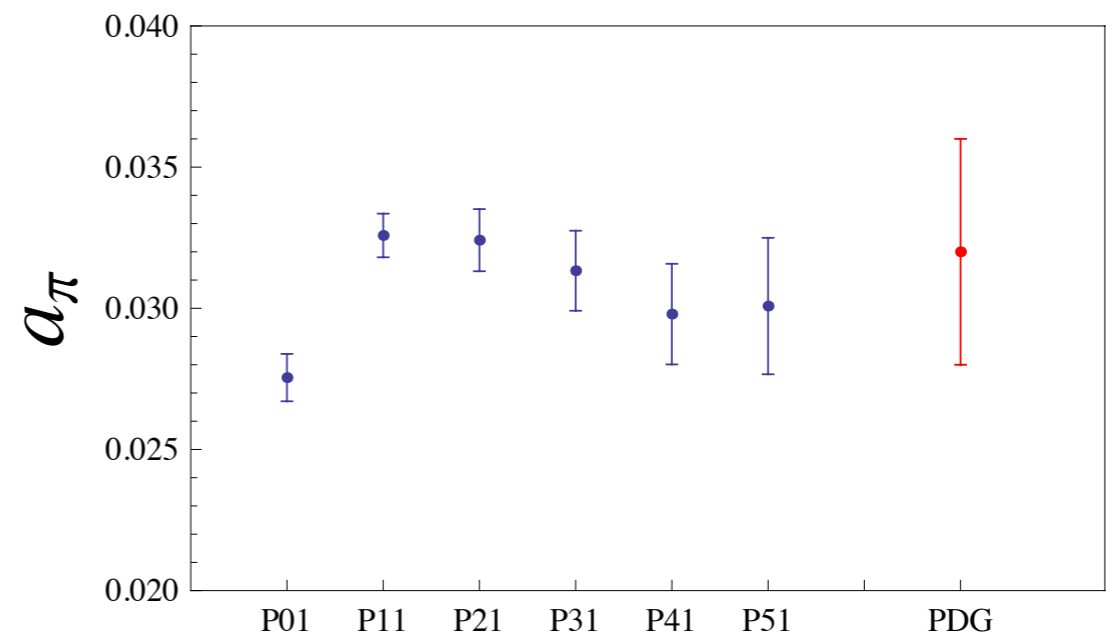
$$\sigma(e^+e^- \rightarrow \gamma^* \rightarrow \gamma P) = \frac{2\pi^2 \alpha^3}{3} (F_{P\gamma^*\gamma}(s, 0))^2 \left(1 - \frac{M_P^2}{s}\right)^3$$

π^0 -TFF

Fit to Space-like data: CELLO'91, CLEO'98, BABAR'09 and Belle'12



$P_1^N(Q^2)$ up to N=5 [P.M, '12]

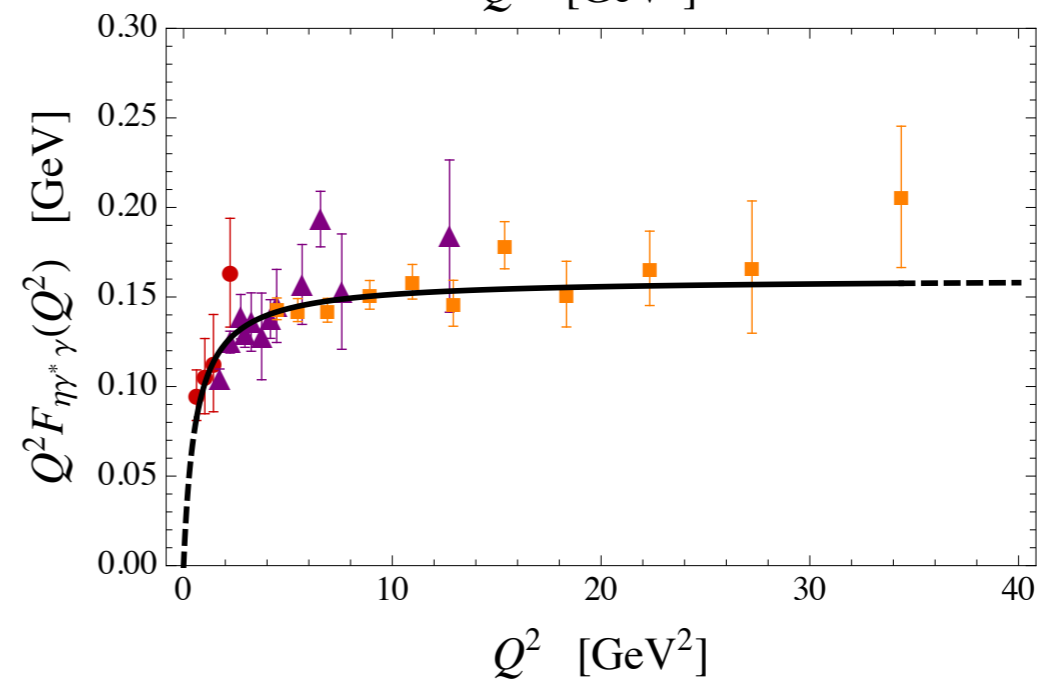
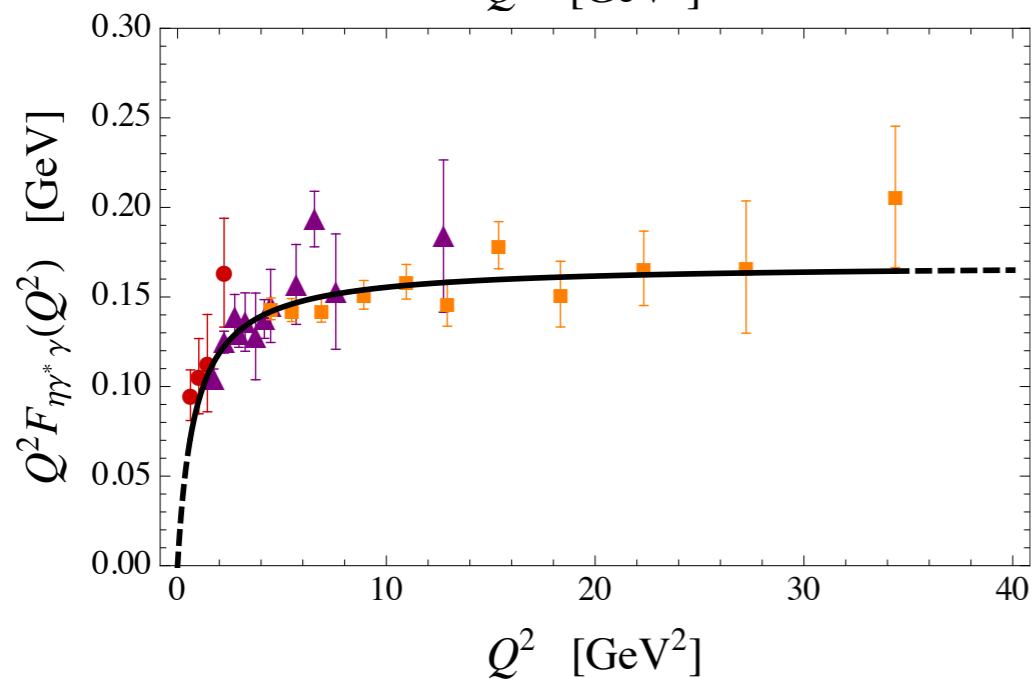
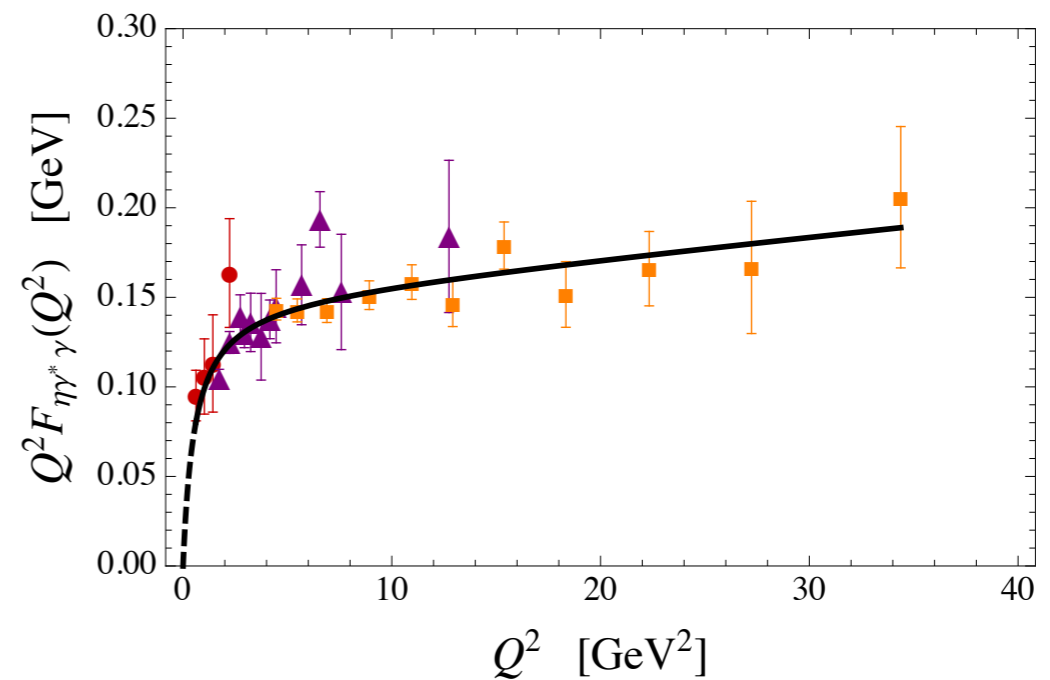
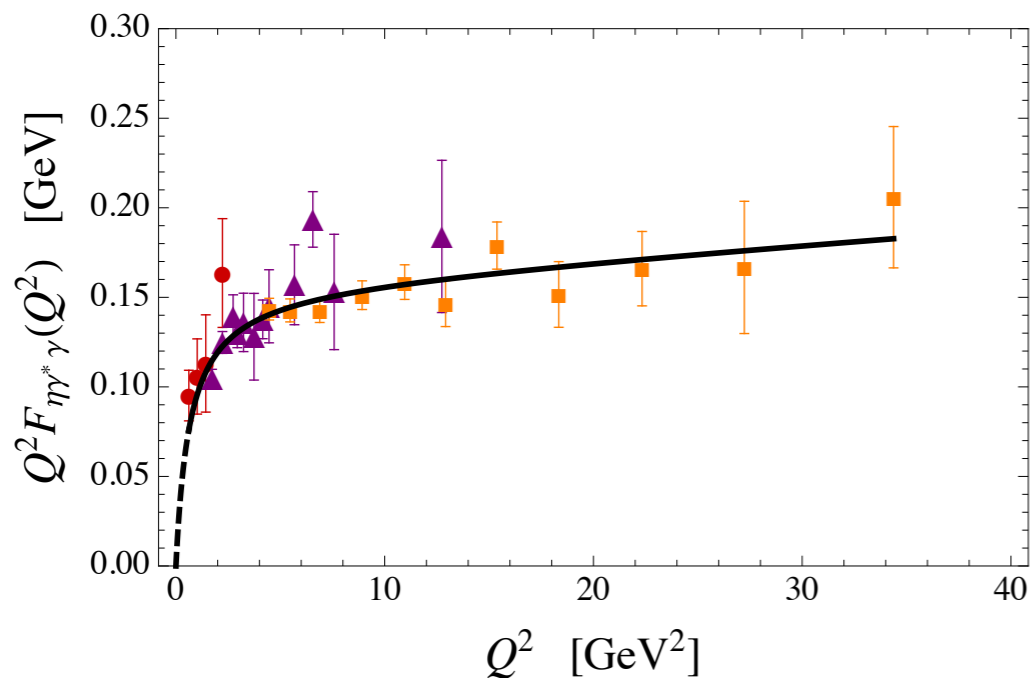


$P_N^N(Q^2)$ up to N=3

η -TFF

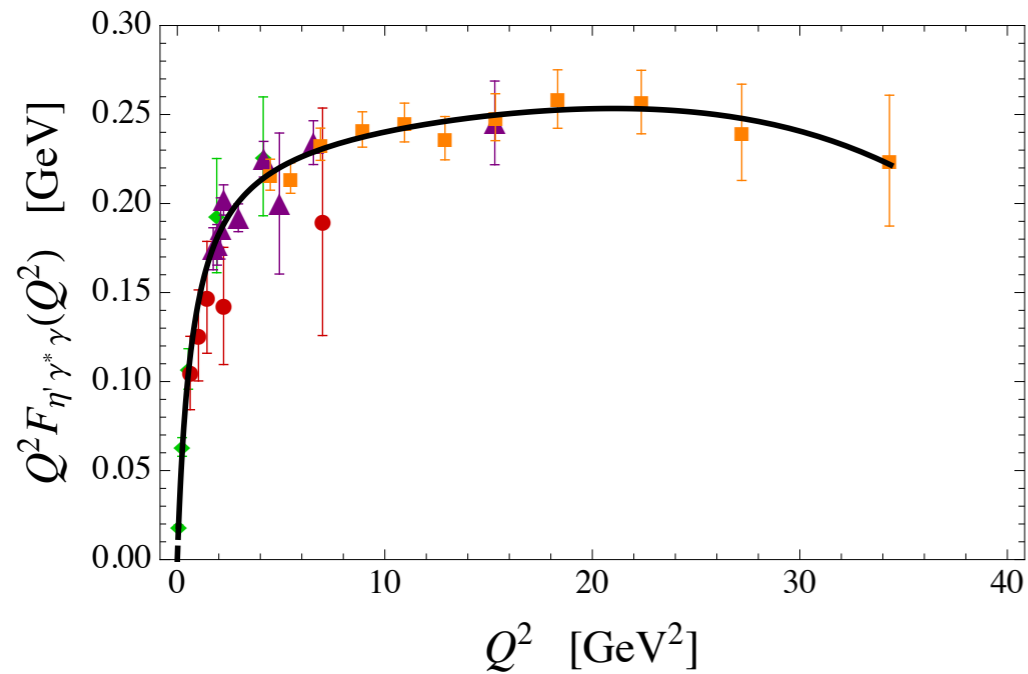
$\Gamma_{\eta \rightarrow \gamma\gamma}$ not included

$\Gamma_{\eta \rightarrow \gamma\gamma}$ included



η' -TFF

$\Gamma_{\eta' \rightarrow \gamma\gamma}$ not included



$\Gamma_{\eta' \rightarrow \gamma\gamma}$ included

