

# Exotic and conventional mesons from lattice QCD

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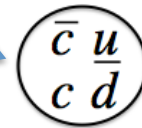


**KRAKÓW, POLAND**  
**29<sup>th</sup> May - 3<sup>rd</sup> June 2014**

# Outline

## Spectrum from lattice QCD

- States well below strong decay threshold:  $\underline{c}\bar{c}$
- Excited states: single-hadron approximation  $\underline{c}\bar{c}$ ,  $D$ ,  $D_s$ , isoscalar mesons  $\underline{u}\bar{u}$  &  $\underline{d}\bar{d}$  &  $\underline{s}\bar{s}$
- Near-threshold states: rigorous treatment
  - ★ first evidence for  $Z_c^+$  from lattice QCD
  - ★  $D_{s0}^*$ (2317) near DK,  $D_{s1}$ (2460) near  $D^*K$
  - ★ first evidence for X(3872) from lattice QCD
- Resonance above threshold: rigorous treatment  $\rho$ ,  $K^*$ ,  $a_1$ ,  $b_1$ ,  $D_0^*$ (2400),  $D_1$ (2430)



Lattice status:

Straightforward procedure,  
many precision results

extensive results completed recently  
highly excited states, many  $J^{PC}$

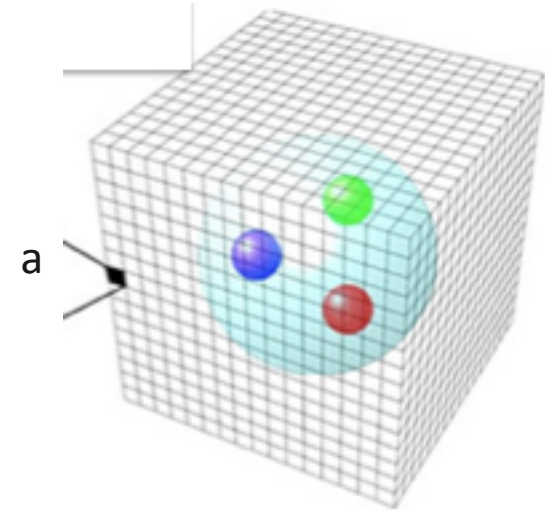
Only few pioneering results available

Only few resonances have been  
studied rigorously

# QCD on lattice: ab initio non-perturbative method

$$L_{QCD} = -\frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu} + \sum_{q=u,d,s,c,b,t} \bar{q} i \gamma_\mu (\partial^\mu + i g_s G_a^\mu T^a) q - m_q \bar{q} q$$

input :  $g_s$  ,  $m_q$



## Evaluation of Feynman path integrals in discretized space-time

$$\langle C \rangle = \int DG Dq D\bar{q} C e^{-S_{QCD}}$$

$$S_{QCD} = \int d^4x L_{QCD}[G(x), q(x), \bar{q}(x)]$$

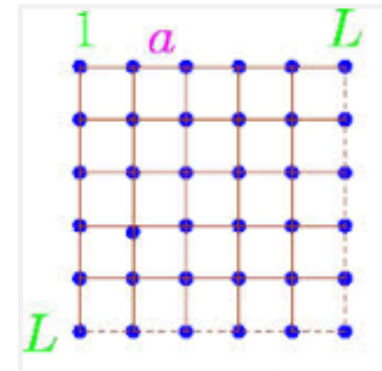
$$S = \int dt L[x(t)]$$

**"Precision" spectrum:  
States well below strong decay threshold**

# States well below threshold

Lattice QCD already determined masses of these states very reliably and precisely  $O(10 \text{ MeV})$ :

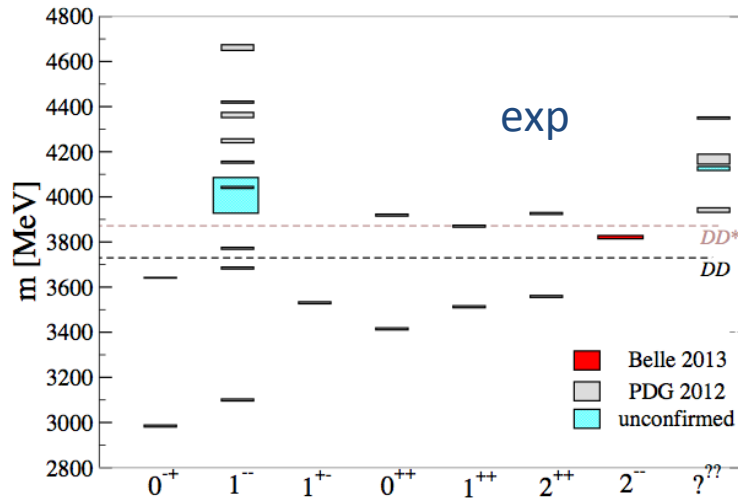
- $m=E$  (for  $P=0$ )
- extrapolation :  $a \rightarrow 0, L \rightarrow \infty$
- extrapolation or interpolation :  $m_q \rightarrow m_q^{\text{phy}}$
- particular care needed for  $am_c$  and  $am_b$  discretization errors: several complementary methods give compatible results



Many precision lattice results available for a number of years!

# States well below threshold

## example: charmonium

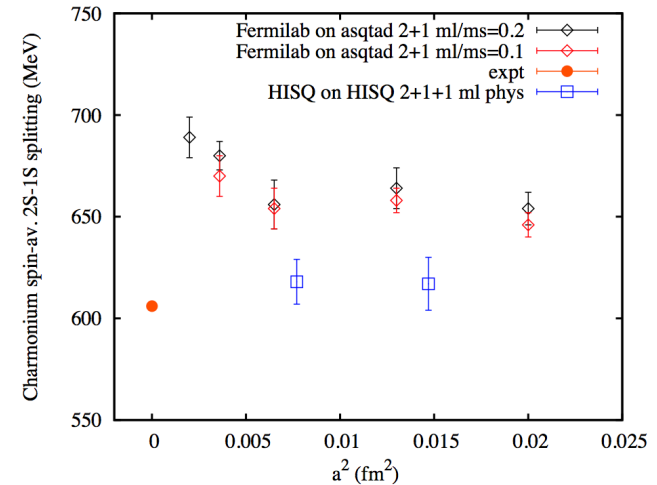


Many results available,  
only few examples shown

[HPQCD/MILC, 2014: Knecht, Galloway,  
Koponen, Davies and DeTar]

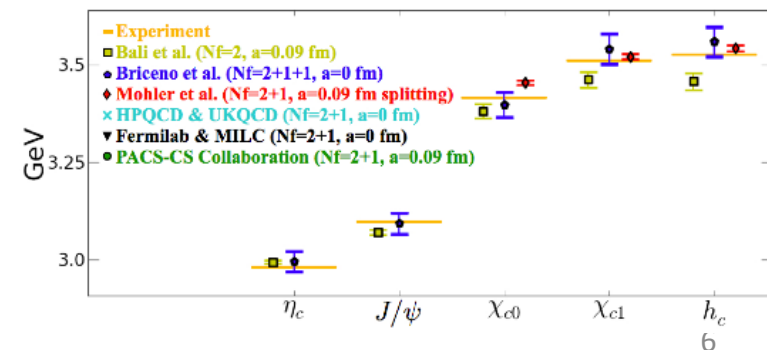
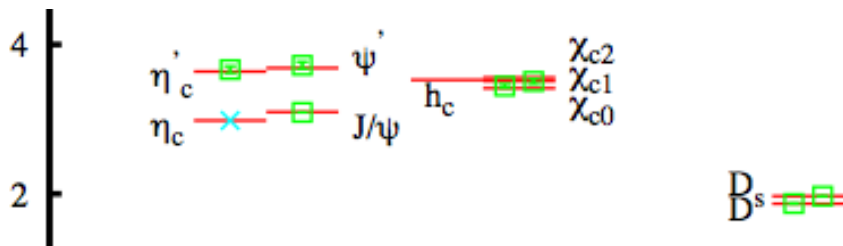
$$\frac{1}{4} M[\eta_c(2S)] + \frac{3}{4} M[\psi(2S)]$$

$$- \frac{1}{4} M[\eta_c(1S)] - \frac{3}{4} M[\psi(1S)]$$



[Briceno, Lin, Bolton, 1207.3536, PRD]

m [GeV] [HPQCD: 1207.5149, PRD]



S. Prelovsek, MESON 14, Lattice spectrum

## "Non-precision" spectrum: states near or above threshold

only one or two  $a, L, m_{u/d}$

limits  $a \rightarrow 0, L \rightarrow \infty, m_{u/d} \rightarrow m_{u/d}^{\text{phy}}$  usually not performed

## Excited states: single-hadron approximation

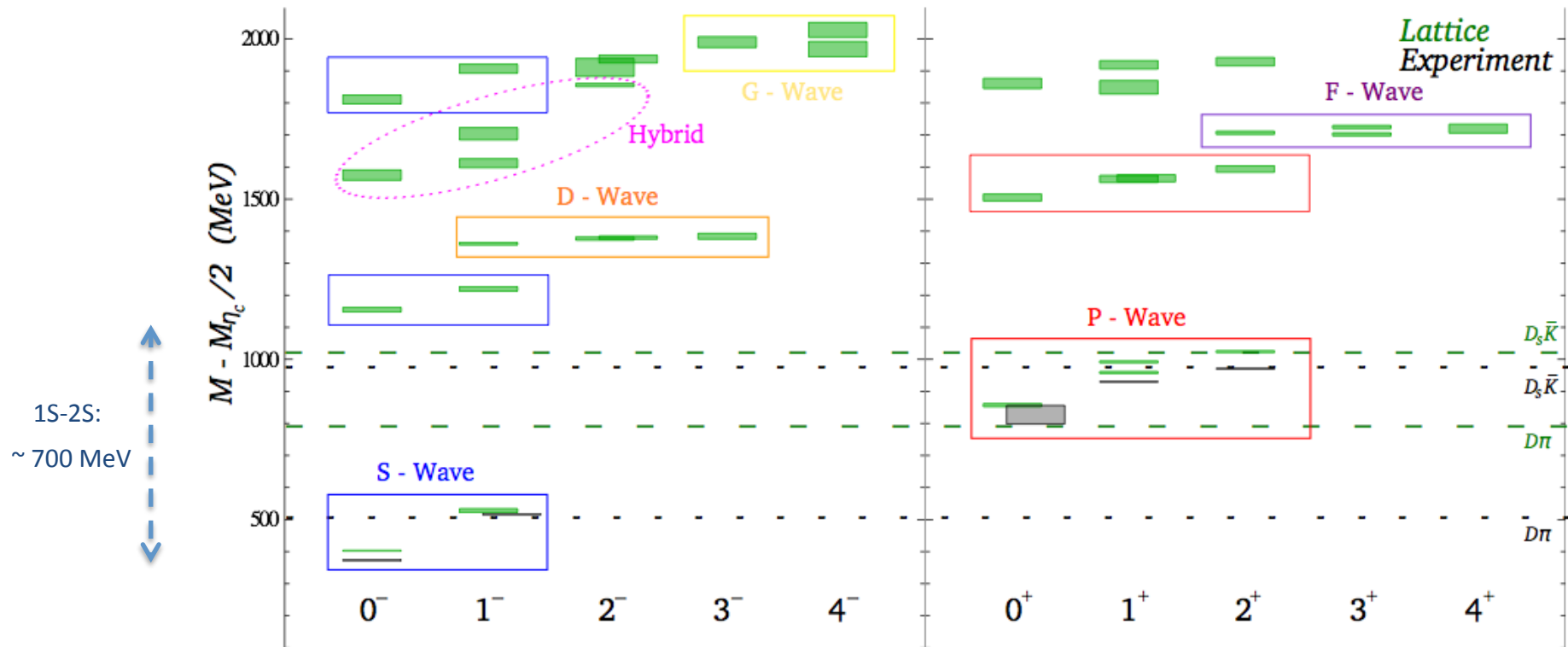
- only interpolating fields  $\mathcal{O} \approx \bar{q} q$
- assumptions: all energy levels correspond to "one-particle" states

$m = E$  (for  $P=0$ )

these are strong assumptions ...



# D spectrum: single hadron approximation



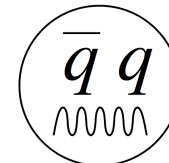
[G. Moir et al, HSC (Hadron Spectrum Coll.): 1301.7670, JHEP]

- $m_{\pi} \approx 400$  MeV,  $L \approx 2.9$  fm,  $N_f = 2+1$
- reliable  $J^P$  determination; many excited states
- identification with  $n^{2S+1}L_J$  multiplets using  $\langle O | n \rangle$
- green: lat, black: exp

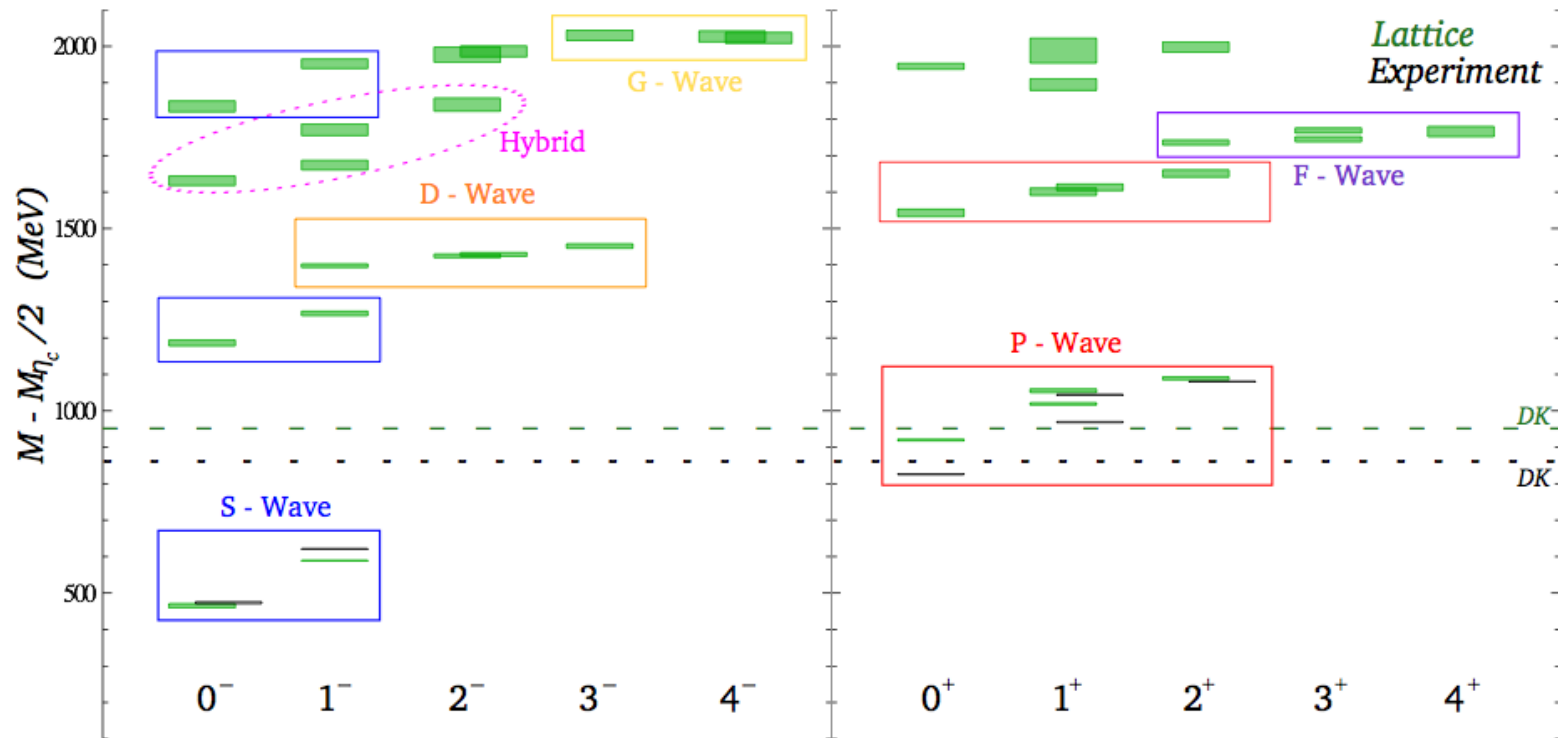
S. Prelovsek, MESON 14, Lattice spectrum

Hybrids:

large overlap with  $O = \underline{q} F_{ij} q$   
gluonic tensor  $F_{ij} = [D_i, D_j]$



# $D_s$ spectrum: single hadron approximation

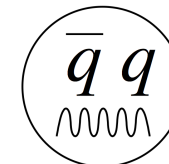


[G. Moir et al., HSC : 1301.7670, JHEP]

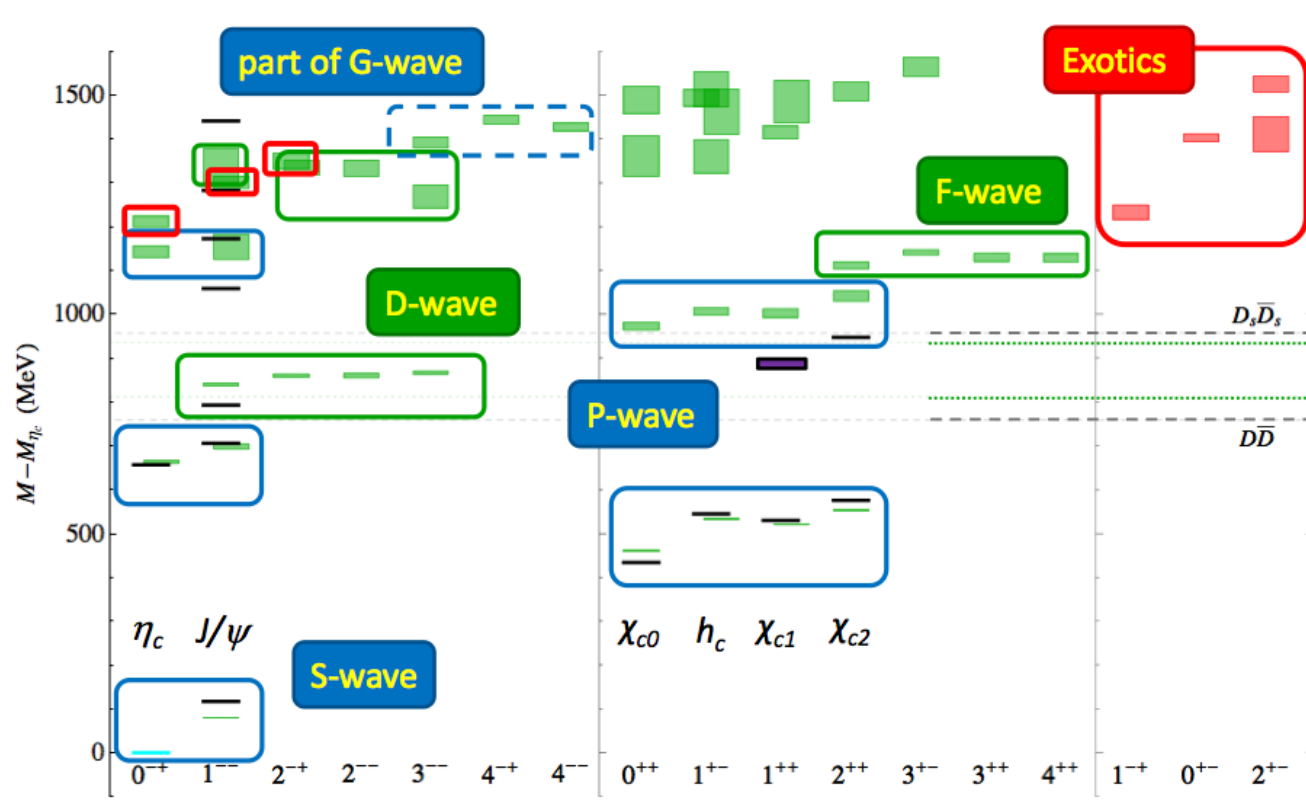
- $m_\pi \approx 400$  MeV,  $L \approx 2.9$  fm,  $N_f = 2+1$
- reliable  $J^{PC}$  determination
- identification with  $n^{2S+1}L_J$  multiplets using  $\langle O | n \rangle$
- green: lat, black: exp

Hybrids:

large overlap with  $O = \bar{q} F_{ij} q$   
 gluonic tensor  $F_{ij} = [D_i, D_j]$



# cc spectrum: single hadron approximation



[HSC, L. Liu et al: 1204.5425, JHEP]

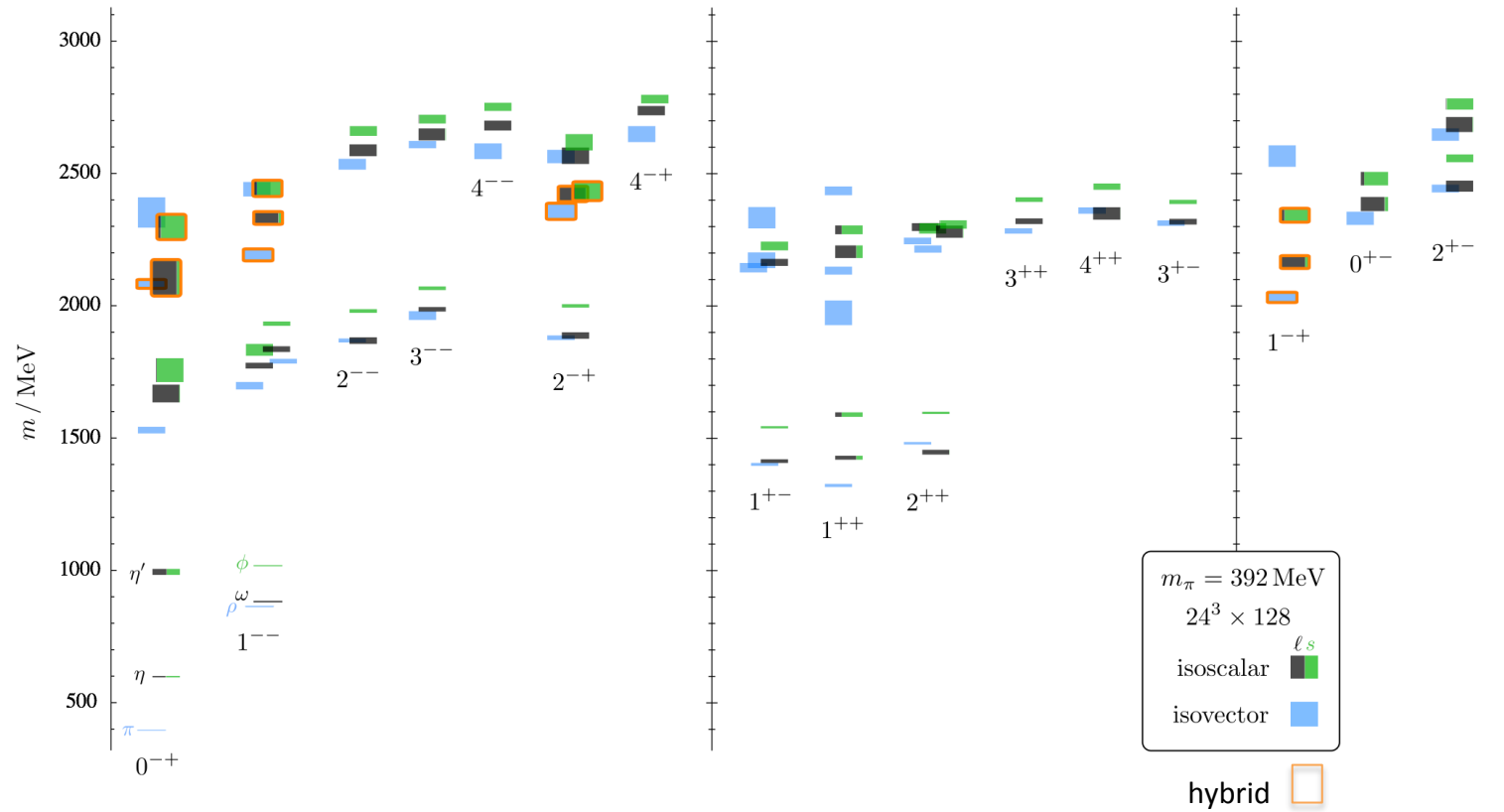
- $m_\pi \approx 400$  MeV,  $L \approx 2.9$  fm,  $N_f = 2+1$
- reliable  $J^{PC}$  determination
- identification with  $n^{2S+1}L_J$  multiplets using  $\langle O | n \rangle$
- green: lat, black: exp

Hybrids:

some of them have exotic  $J^{PC}$   
large overlap with  $O = \underline{q} F_{ij} q$



# Isoscalar mesons: single hadron approximation



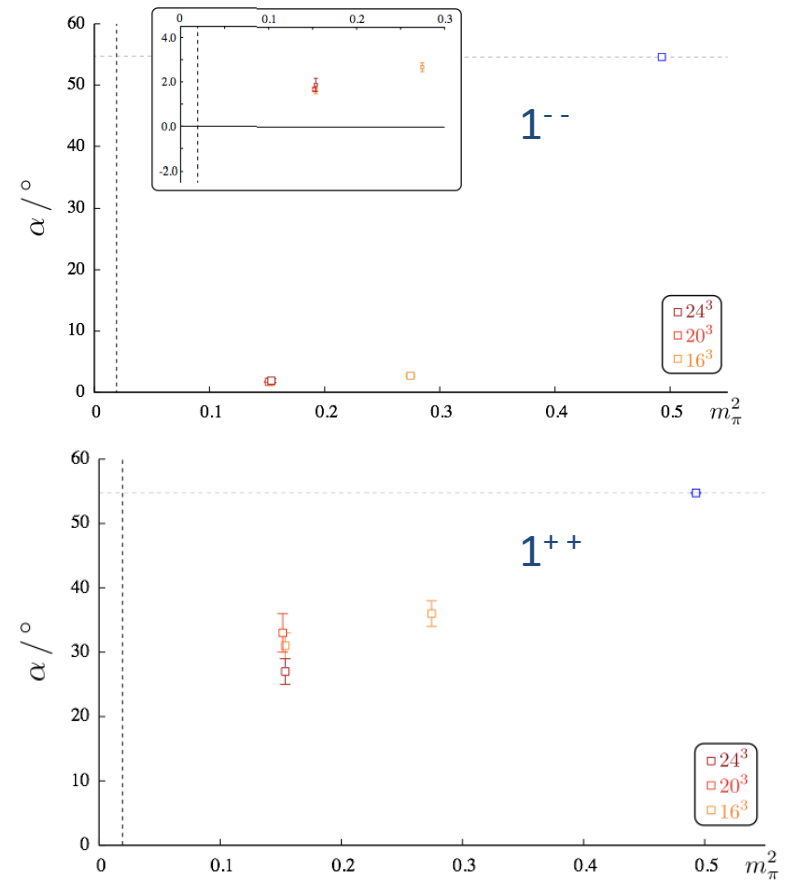
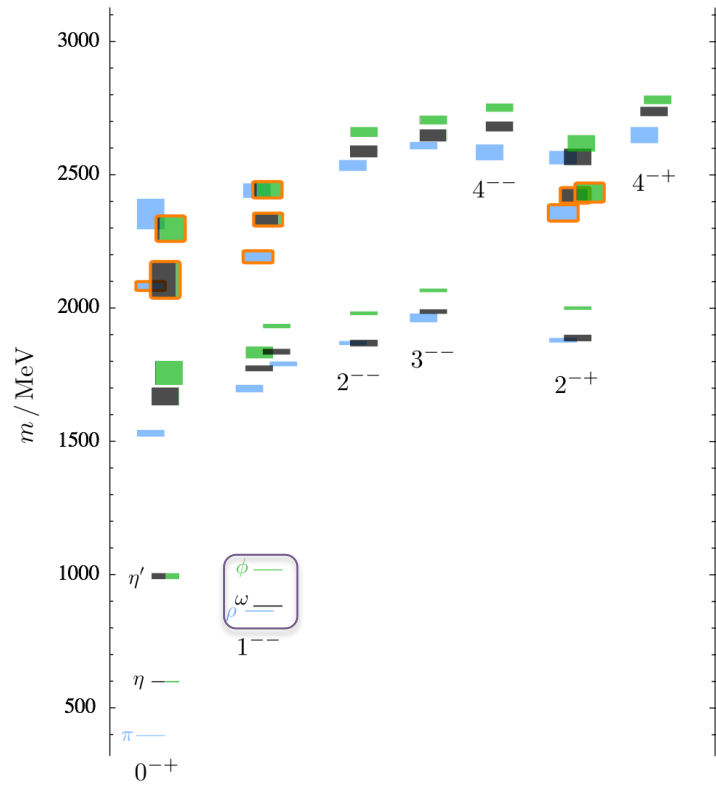
$$\begin{pmatrix} |a\rangle \\ |b\rangle \end{pmatrix} = \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} |\ell\rangle \\ |s\rangle \end{pmatrix}$$

$$|\ell\rangle \equiv \frac{1}{\sqrt{2}} (|u\bar{u}\rangle + |d\bar{d}\rangle)$$

$$|s\rangle \equiv |s\bar{s}\rangle$$

[HSC : Dudek, Edward, Guo, Thomas: 1309.2608, PRD]

# Isoscalar mesons: mixing angle



$$\begin{pmatrix} |\mathbf{a}\rangle \\ |\mathbf{b}\rangle \end{pmatrix} = \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} |\ell\rangle \\ |s\rangle \end{pmatrix}$$

$$\begin{aligned} |\ell\rangle &\equiv \frac{1}{\sqrt{2}} (|u\bar{u}\rangle + |d\bar{d}\rangle) \\ |s\rangle &\equiv |s\bar{s}\rangle \end{aligned}$$

[HSC : Dudek, Edward, Guo, Thomas: 1309.2608, PRD]

## States near threshold: rigorous treatment

note: most of interesting states are found near threshold:

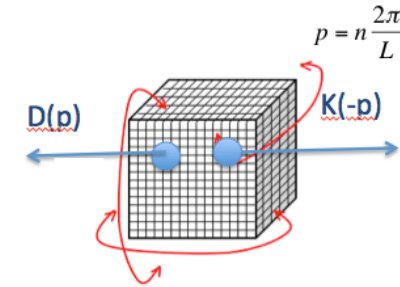
$$D_{s0}^*(2317), X(3872), Z_c^+, Z_b^+$$

# Discrete energy spectrum from correlators

$$\langle C \rangle \propto \int DG Dq D\bar{q} C(q, \bar{q}, G) e^{i S_{QCD} / \hbar}, \quad S_{QCD} = \int d^4 x L_{QCD}$$

Example: meson channel with given  $J^{PC}$

$$\mathcal{O} = \bar{q} \Gamma q, \quad \bar{q} \Gamma' q, \quad (\bar{q} \Gamma_1 q)(\bar{q} \Gamma_2 q), \dots$$



$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^+(0) | 0 \rangle$$

$$= \sum_n \langle 0 | \mathcal{O}_i | n \rangle e^{-E_n t} \langle n | \mathcal{O}_j^+ | 0 \rangle = \sum_n Z_i^n Z_j^{n*} e^{-E_n t} \quad Z_i^n = \langle 0 | \mathcal{O}_i | n \rangle$$

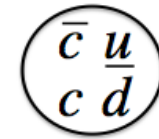
All physical states with given  $J^{PC}$  appear as energy levels  $E_n$  in principle : single particle, two-particle,...

channel	:	"eigenstates"
$J^{PC} = 0^{++}, \bar{s}c$	:	$D_{s0}^*(2317), DK$
$J^{PC} = 1^{+-}, \bar{c}c\bar{d}u$	:	$Z_c^+, J/\psi \pi^+, \dots$
$J^{PC} = 1^{--}, \bar{s}u$	:	$K^*(892), K\pi$

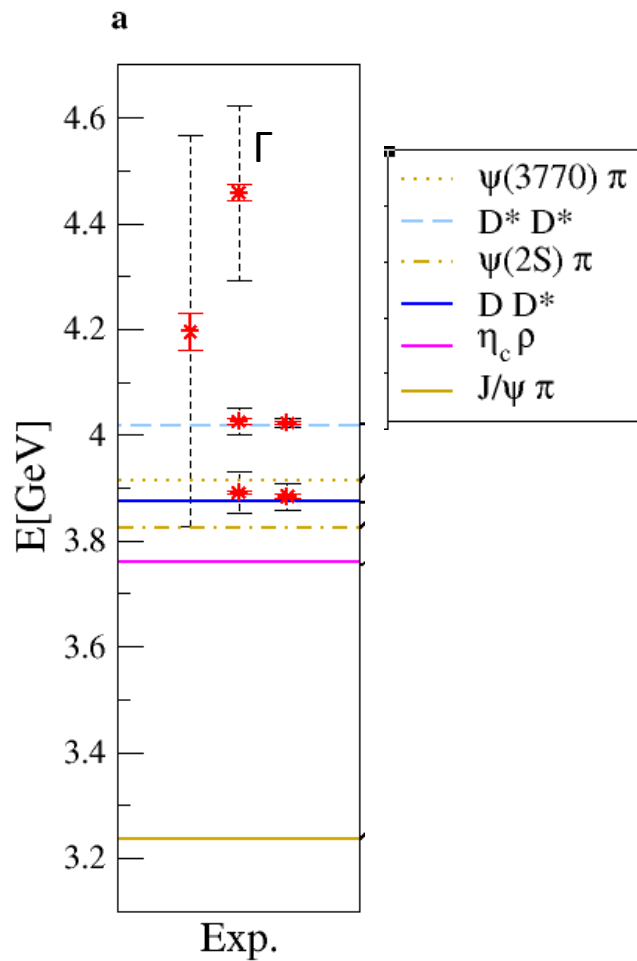
In experiment: these correspond to two-meson decay products with continuous spectrum.

On lattice: these are discrete due to finite box and periodic BC.

# Charged charmonium $Z_c^+$ : experimental status

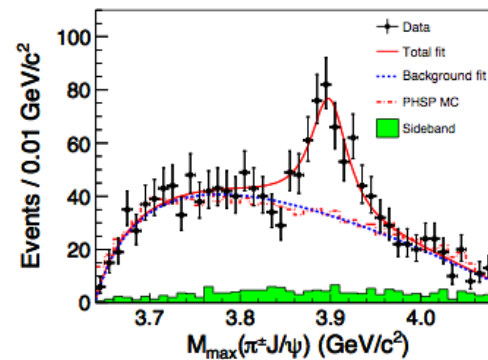


candidates with preferred  
 $|G=1^+, J^{PC}=1^{+-}$

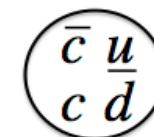
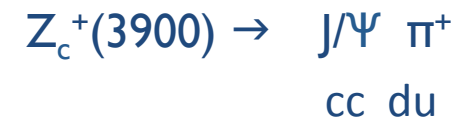


[review: Brambilla et al., 1404.3723]

particle	C	J <sup>P</sup>	decay	year	coll
$Z^+(4430)$	-	1+	$\psi(2S) \pi^+$	2008	Belle, BABAR, LHCb
$Z_c^+(3900)$	-	?	$J/\psi \pi^+$	2013	BESIII, Belle, CLEOc
$Z_c^+(3885)$	-	1+	$(DD^*)^+$	2013	BESIII
$Z_c^+(4020)$	-	?	$h_c(1P) \pi^+$	2013	BESIII
$Z_c^+(4025)$	-	?	$(D^* D^*)^+$	2013	BES III
$Z^+(4200)$	-	1+	$J/\psi \pi^+$	2014	Belle
$Z^+(4050)$	+	?	$\chi_{c1} \pi^+$	2008	Belle
$Z^+(4250)$	+	?	$\chi_{c1} \pi^+$	2008	Belle



[BESIII, 2013, 1303.5949, PRL]

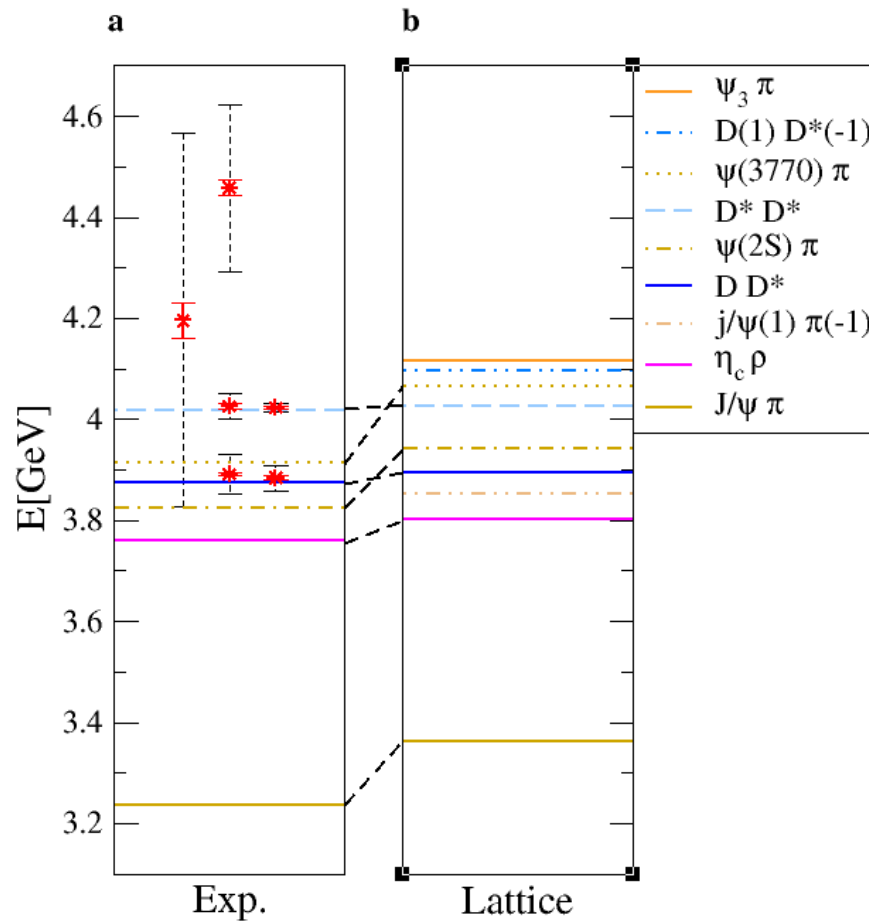
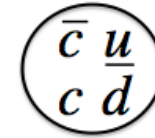




# Previous searches for $Z_c^+$ from lattice

- Search in  $J^{PC}=1^{+-}$  channel for  $m < 4$  GeV:  
no  $Z_c^+$  candidate found  
[S.P. & L. Leskovec, 1308.2097, PLB]
- Search for resonance in  $D\bar{D}^*$  scattering with  $J^{PC}=1^{+-}$  near threshold  $E \sim 3.9$  GeV  
no  $Z_c^+$  candidate found  
[Y. Chen et al, 1403.1318, PRD]

# First evidence for $Z_c^+$ from lattice: $I^G=1^+, J^{PC}=1^{+-}$



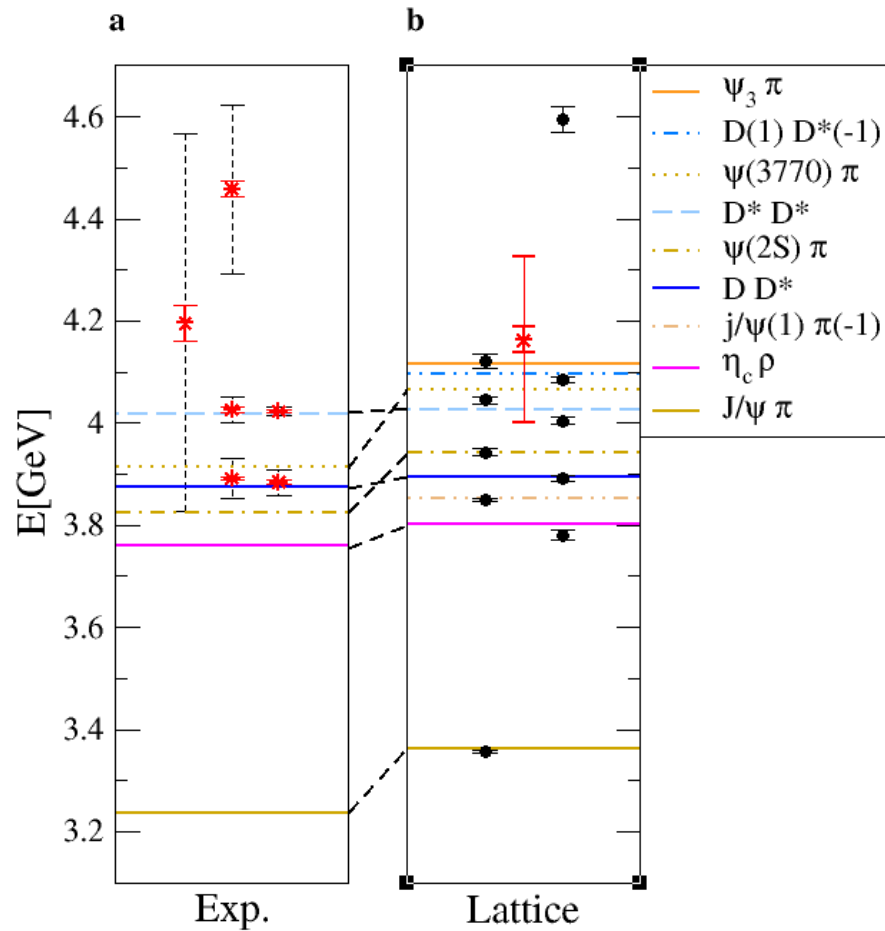
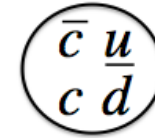
Lattice:

Energies of all two-particle states with  $E < 4.3$  GeV

[S.P., Lang, Leskovec, Mohler, 1405.7623]

$m_\pi \approx 266$  MeV,  $L \approx 2$  fm,  $N_f=2$

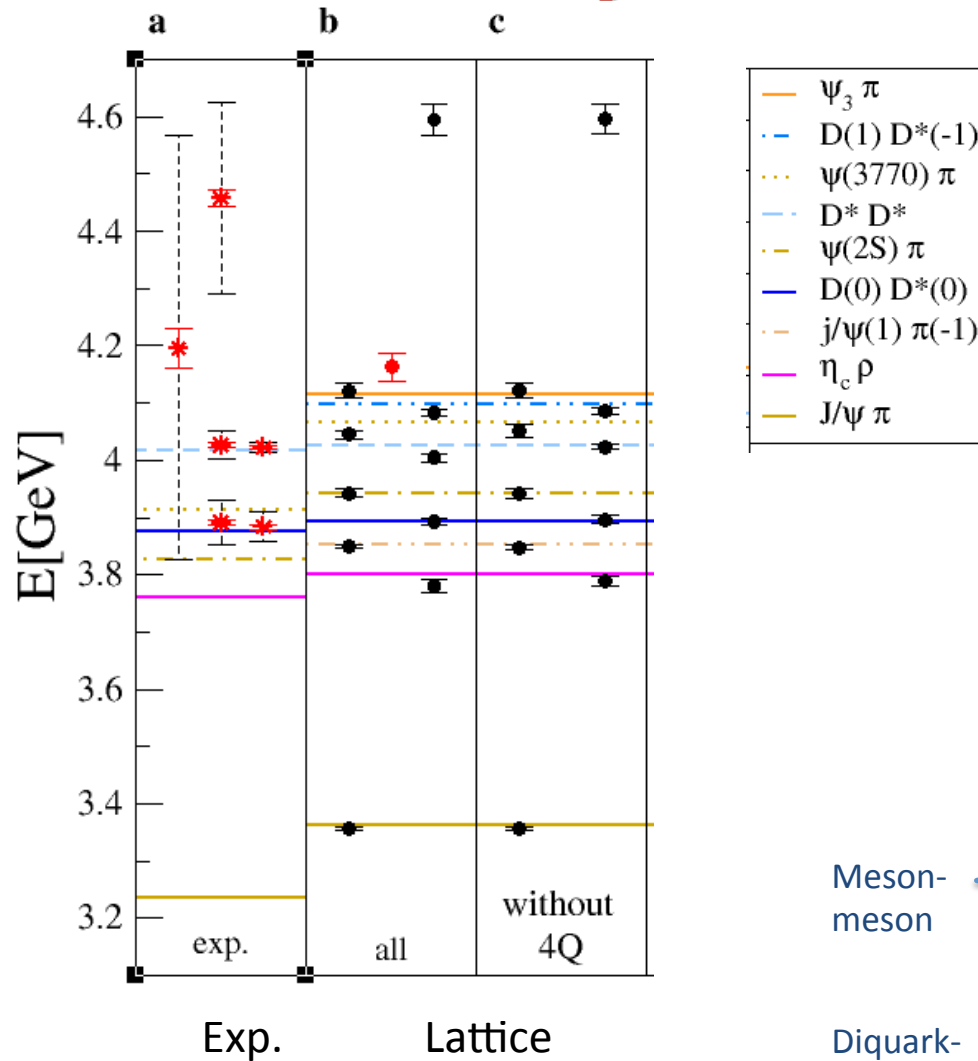
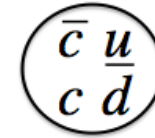
# First evidence for $Z_c^+$ from lattice: $I^G=1^+, J^{PC}=1^{+-}$



[S.P., Lang, Leskovec, Mohler, 1405.7623]

$m_\pi \approx 266$  MeV,  $L \approx 2$  fm,  $N_f=2$

# First evidence for $Z_c^+$ from lattice: $I^G=1^+, J^{PC}=1^{+-}$



[S.P., Lang, Leskovec, Mohler, 1405.7623]

Meson-meson

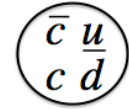
$$\begin{cases} \mathcal{O}_1^{\psi(0)\pi(0)} = \bar{c}\gamma_i c(0) \bar{d}\gamma_5 u(0) \\ \mathcal{O}^{\eta_c(0)\rho(0)} = \bar{c}\gamma_5 c(0) \bar{d}\gamma_i u(0), \\ \mathcal{O}_1^{D(0)D^*(0)} = \bar{c}\gamma_5 u(0) \bar{d}\gamma_i c(0) + \{\gamma_5 \leftrightarrow \gamma_i\} \end{cases}$$

and 11 others ..

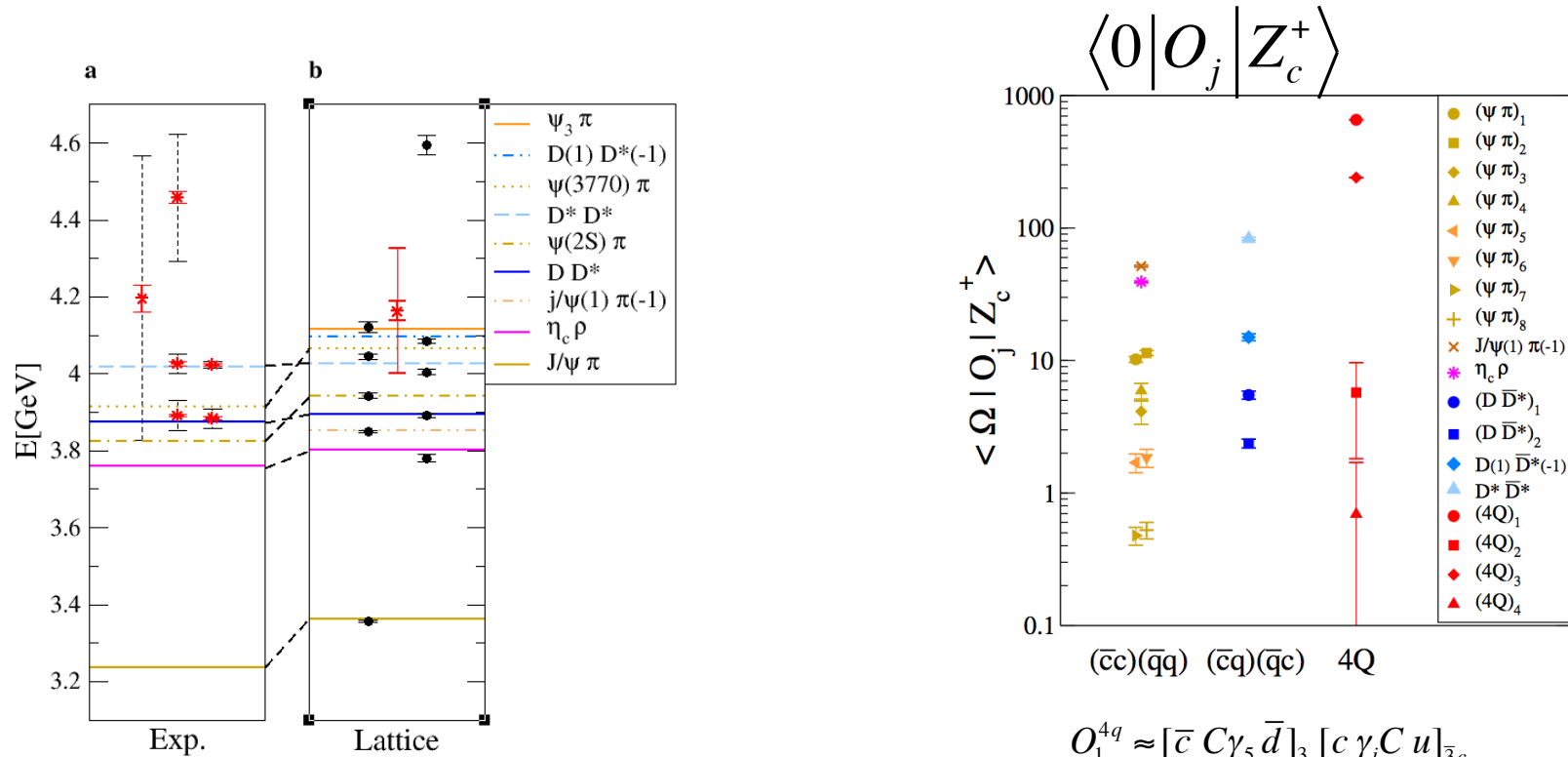
Diquark-antidiquark

$$\begin{cases} \mathcal{O}_1^{4q} \approx [\bar{c} C \gamma_5 \bar{d}]_{3_c} [c \gamma_i C u]_{\bar{3}_c} \\ \mathcal{O}_2^{4q} \approx [\bar{c} C \bar{d}]_{3_c} [c \gamma_i \gamma_5 C u]_{\bar{3}_c} \end{cases}$$

and 2 others ..



# First evidence for $Z_c^+$ from lat: $I^G=1^+, J^{PC}=1^+$



## Nearby experimental candidates:

$Z_c^+(4020)$ ,  $\Gamma=7.9 \pm 3.7$  MeV BESIII 2013

$Z_c^+(4025)$ ,  $\Gamma=24.8 \pm 9.5$  MeV BESIII 2013

$Z_c^+(4200)$ ,  $\Gamma=370 \pm 110$  MeV Belle, Moriond 2014

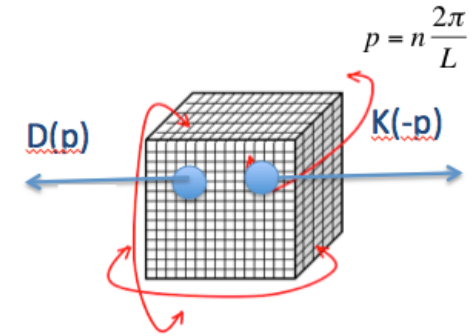
Lattice ( $m_\pi=266$  MeV,  $N_f=2$ ) :

$m(Z_c^+) = 4.16$  GeV

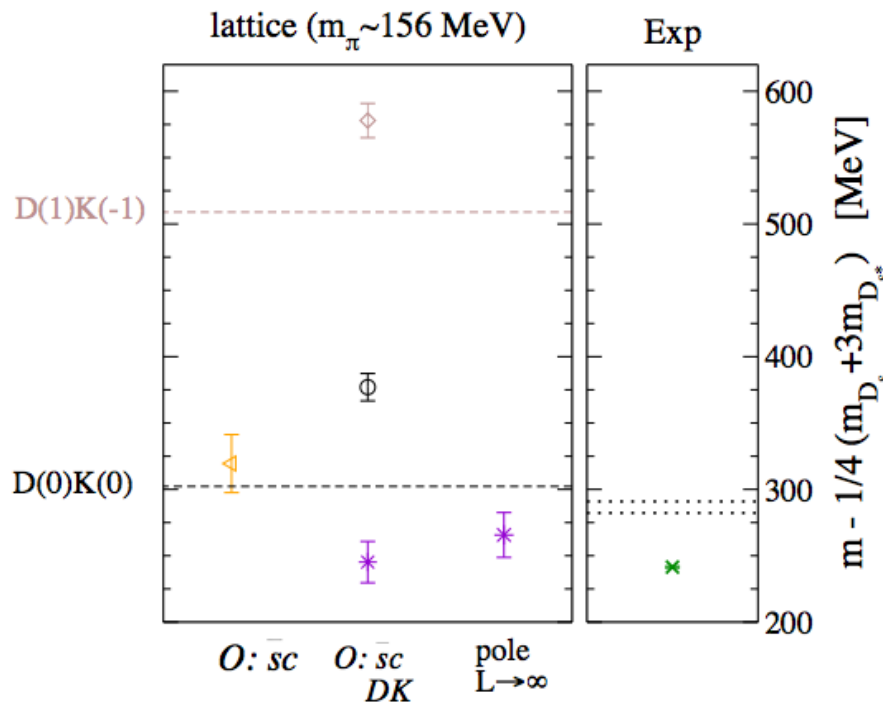
$\pm 0.163$  GeV  $\pm O(\Gamma)$

[S.P., Lang, Leskovec, Mohler, 1405.7623]

# $D_{s0}^*(2317)$ and DK scattering : $J^P=0^+$



$$\mathcal{O} : \bar{s} c, \quad DK \approx [\bar{d}\gamma_5 c] [\bar{s}\gamma_5 d]$$



Rigorous relation [M. Luscher, 1991]:

$E \rightarrow \delta(E)$  phase shift for DK scattering in s-wave

$$p \cot \delta(p) = \frac{1}{a_0} + \frac{1}{2} r_0 p^2$$

pole position and mass of  $D_{s0}^*(2317)$

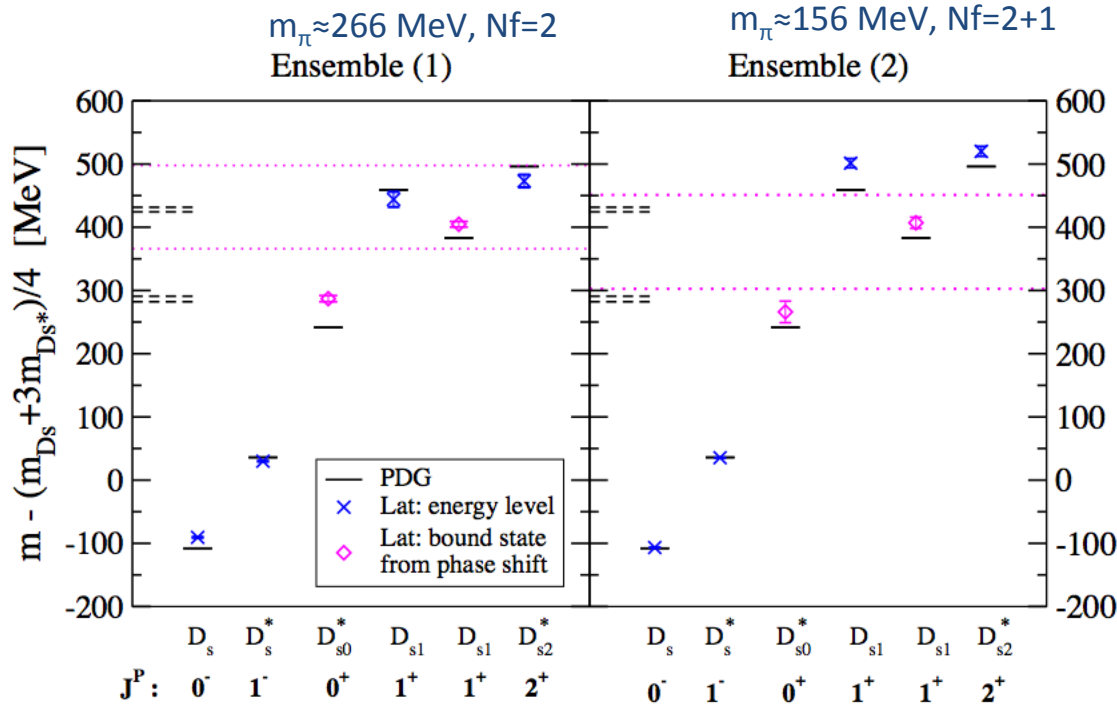
$$S \propto \frac{1}{\cot \delta - i} = \infty \quad \cot \delta(p_{BS}) = i$$

$$m_{D_{s0}}^{lat, L \rightarrow \infty} = E_D(p_{BS}) + E_K(p_{BS})$$

[D. Mohler, C. Lang, L. Leskovec, S.P., R. Woloshyn: PRL 2013]

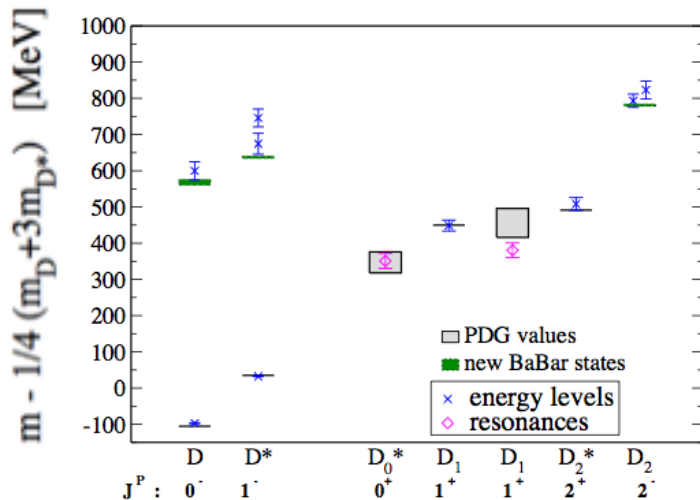
$m_\pi \approx 156$  MeV,  $L \approx 2.9$  fm,  $N_f = 2+1$

# Charmed meson spectrum: rigorous treatment



## Ds mesons (near-threshold)

[C. Lang, L. Leskovec, D. Mohler, S.P., R. Woloshyn: PRL 2013, 1403.8103]



## D mesons (resonances)

[D. Mohler, S.P., R. Woloshyn: PRD 2013]

## Charmed scalar meson "puzzle" revisited

• why do these scalar partners have mass so close ?

$$D_0^*(2400): M \approx 2318 \text{ MeV} \quad \Gamma \approx 267 \text{ MeV} \quad \bar{c}u \text{ or } \bar{c}u\bar{s}s ?$$

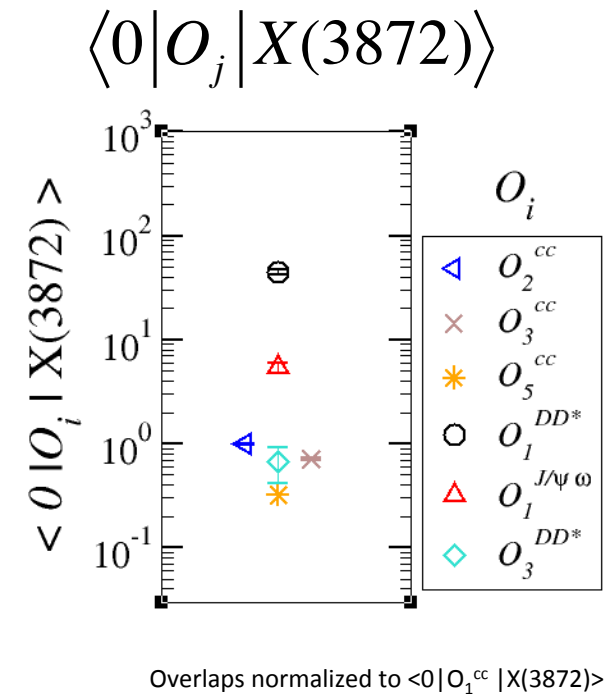
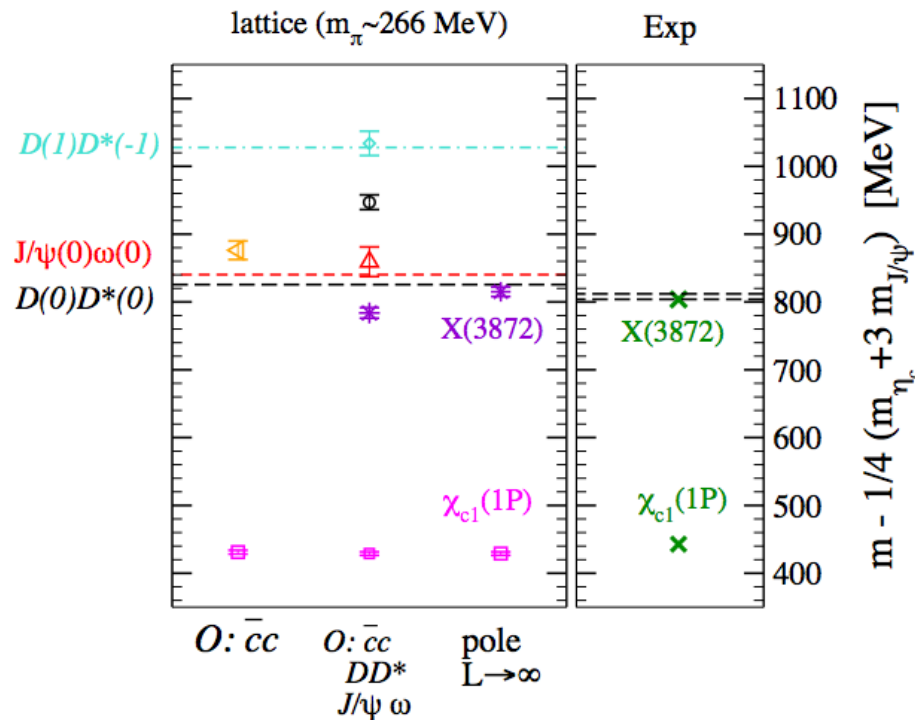
$$D_{s0}(2317): M \approx 2318 \text{ MeV} \quad \Gamma \approx 0 \text{ MeV} \quad \bar{c}s \text{ or } \bar{c}s[\bar{u}u + \bar{d}d] ?$$

1) is  $D_0^*$  mass pushed up : valence  $\bar{s}s$  pair ?? ✗

2) is  $D_{s0}^*$  mass pushed down : effect of DK threshold ?? ✓

# First evidence for X(3872) from lattice : $J^{PC}=1^{++}, I=0$

$\mathcal{O} : \bar{c}c, DD^*, J/\psi\omega$



Overlaps normalized to  $\langle 0 | O_1^{cc} | X(3872) \rangle$

S. P. and L. Leskovec : 1307.5172, PRL

$m_{\pi} \approx 266$  MeV,  $L \approx 2$  fm,  $N_f=2$

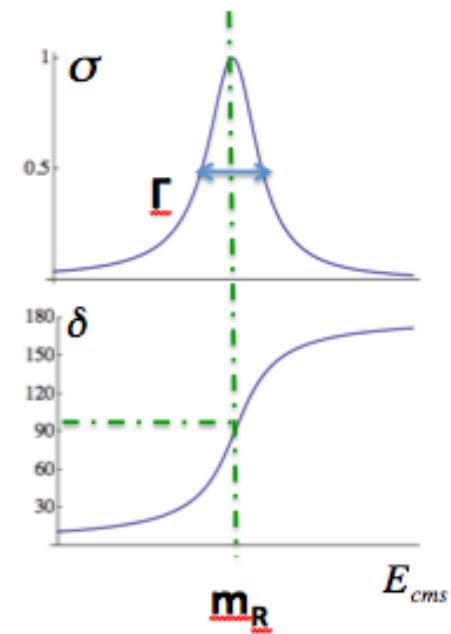
X(3872)	$m - (m_{D0} + m_{D0^*})$
lat	$- 11 \pm 7$ MeV
exp	$- 0.14 \pm 0.22$ MeV

lat: simulations on larger L required

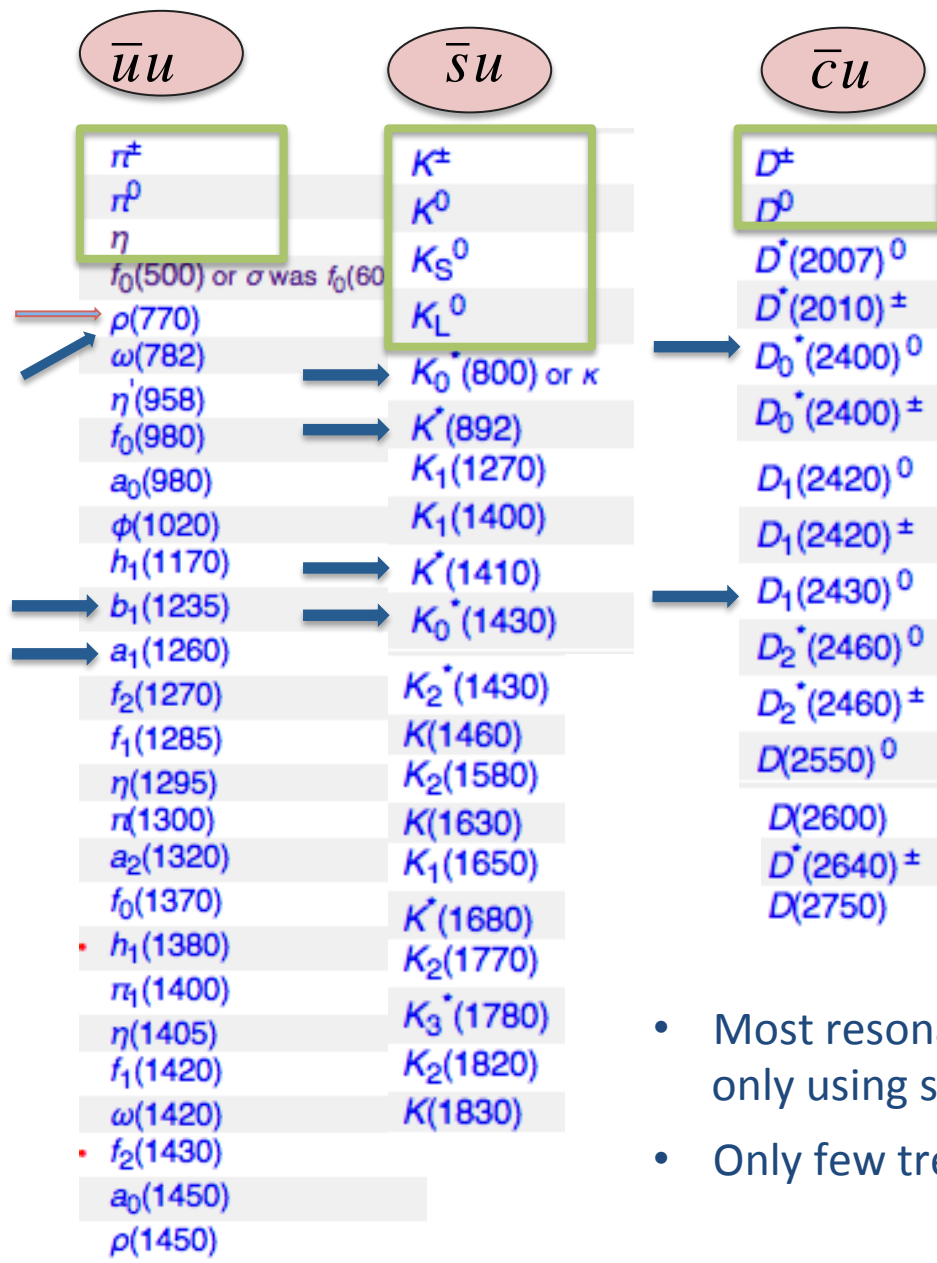
exp: Tomaradze et al., 1212.4191 24



# States above threshold - resonances: rigorous treatment



# Almost all mesons are resonances (decay strongly)



stable on strong decay: ab-initio OK

others decay strongly; hadronic resonances

$\rightarrow$  only resonance simulated properly by several collab. (first simulation: CP-PACS 2007)

$\rightarrow$  SP and coll. (Lang, Mohler, Leskovec, Woloshyn) (2011- 2014)

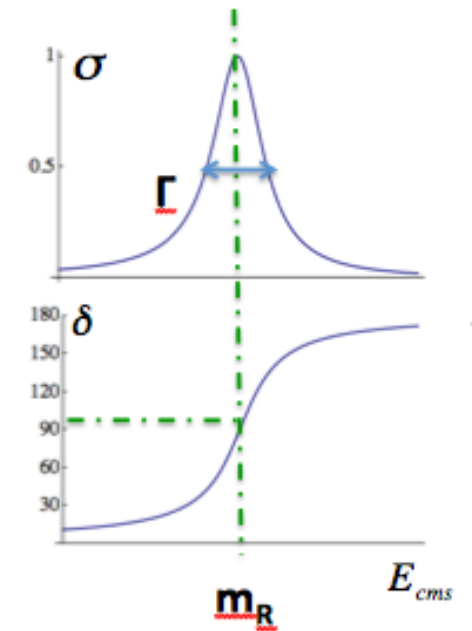
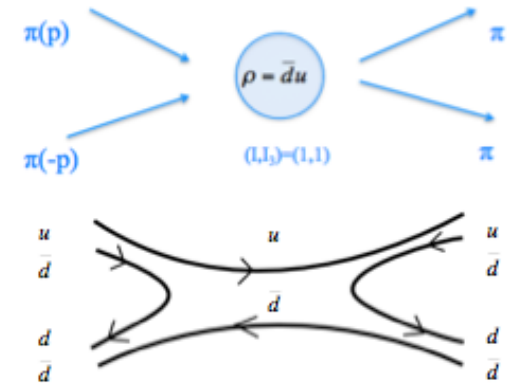
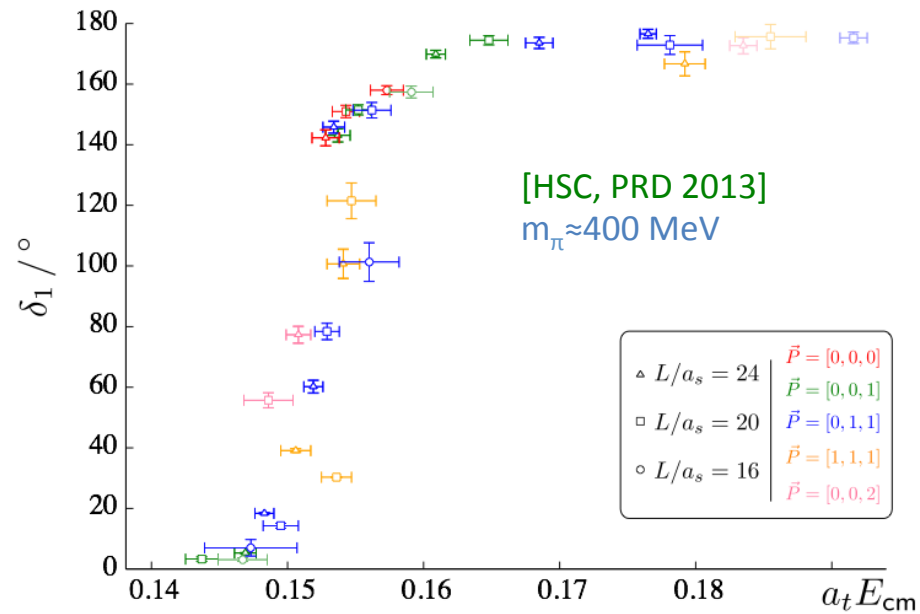
- Most resonances simulated only using single-hadron app.
- Only few treated rigorously

# $\rho$ resonance

$P \neq 0$ :  $s = E^2 - P^2$ , Luscher-type relation  $s \rightarrow \delta(s)$

$m_1 = m_2$

[Rummukainen, Gotlieb 1995]

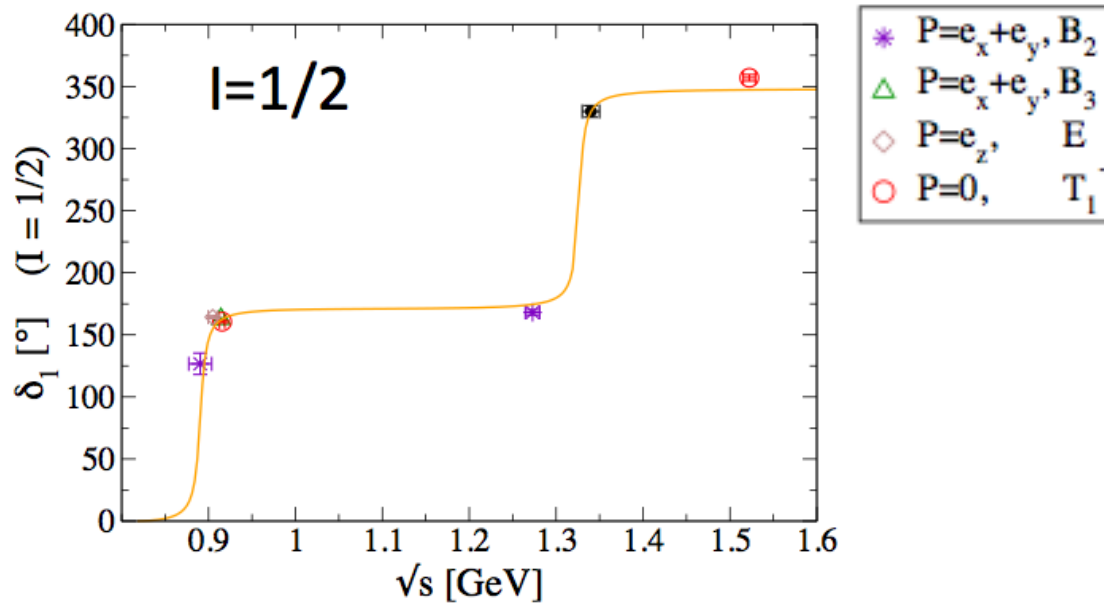


Simulation also by CP-PACS, PACS-CS, QCDSF, ETMC

[Lang, Mohler, S.P., Vidmar, PRD 2011]  $m_\pi \approx 266$  MeV

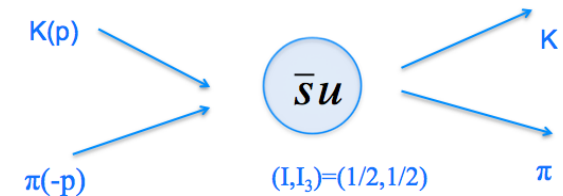
# $K^*(892)$ resonance: first lat determination of the width

$K\pi$ ,  $I=1/2$ : p-wave phase shift



[S.P., Lang, Leskovec, Mohler,  
1307.0736, PRD 2013]

$m_\pi \approx 266$  MeV



$$\text{BW : } \delta = \text{acot} \frac{m_R^2 - E_{\text{cms}}^2}{m_R \Gamma}$$

$$\Gamma[K^* \rightarrow K\pi] = \frac{g^2 p^{*3}}{6\pi s}$$

	$m_{K^*(892)}$ [MeV]	$g_{K^*(892)}$ [no unit]
lat	$891 \pm 14$	$5.7 \pm 1.6$
exp	$891.66 \pm 0.26$	$5.72 \pm 0.06$

# D-meson resonance masses and widths

$$\Gamma(E) \equiv g^2 \frac{P}{E^2} \quad \text{g is compared to exp instead of } \Gamma \quad (\Gamma \text{ depends on phase sp. and } m_\pi)$$

$J^P=0^+ : D \pi$

$J^P=1^+ : D^* \pi$

(analysis of spectrum in this case is based on an assumption given in paper below)

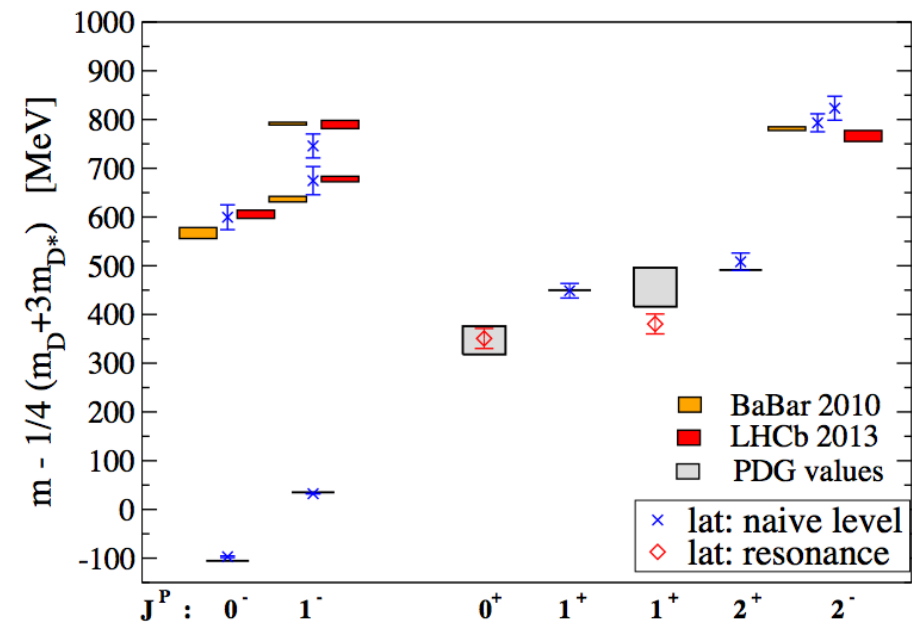
$D_0^*(2400)$	$m - 1/4(m_D+3m_{D^*})$	$g$
lat	$351 \pm 21 \text{ MeV}$	$2.55 \pm 0.21 \text{ GeV}$
exp	$347 \pm 29 \text{ MeV}$	$1.92 \pm 0.14 \text{ GeV}$

$D_1(2430)$	$m - 1/4(m_D+3m_{D^*})$	$g$
lat	$381 \pm 20 \text{ MeV}$	$2.01 \pm 0.15 \text{ GeV}$
exp	$456 \pm 40 \text{ MeV}$	$2.50 \pm 0.40 \text{ GeV}$

first lattice result for strong decay width of a hadron containing charm quark

[D. Mohler, S.P., R. Woloshyn: 1208.4059, PRD]

- $m_\pi \approx 266 \text{ MeV}$ ,  $L \approx 2 \text{ fm}$ ,  $N_f=2$



# Lightest axial resonances $a_1(1260)$ and $b_1(1235)$

Simulating scattering:

$\rho\pi$  in  $1^{++}$  channel to extract  $a_1$

$\omega\pi$  in  $1^{+-}$  channel to extract  $b_1$

$$\Gamma(E) \equiv g^2 \frac{p}{E^2}$$

resonance	$a_1(1260)$			$b_1(1235)$	
quantity	$m_{a_1}^{\text{res}}$ [GeV]	$g_{a_1\rho\pi}$ [GeV]	$a_{l=0}^{\rho\pi}$ [fm]	$m_{b_1}^{\text{res}}$ [GeV]	$g_{b_1\omega\pi}$ [GeV]
lat	$1.435(53)^{(+0}_{-109)}$	1.71(39)	0.62(28)	$1.414(36)^{(+0}_{-83)}$	input
exp	1.230(40)	1.35(30)	-	1.2295(32)	0.787(25)

[Lang, Leskovec, Mohler, S.P. , 1401.2088, JHEP]

# Conclusions

Lattice QCD is the only non-perturbative theoretical tool that depends on the same parameters as QCD Lagrangian.

Available already for some time

- Precise result for states well below threshold

Important developments during past two years

- a) Extensive results for excited multiplets within single-hadron approximation
- b) Effect of threshold on near-threshold states: crucial for (c) and (d)
- c) First evidence for an exotic  $Z_c^+$  with two valence quarks and antiquarks
- d) First evidence for X(3872)
- e) First determination of  $\Gamma$  for strange and charmed meson resonances

# Thanks to my collaborators

Christian B. Lang (Graz)

Daniel Mohler (Fermilab)

Luka Leskovec (Ljubljana)

Richard Woloshyn (Triumf)



# Backup slides

[review: Brambilla et al., 1404.3723]

TABLE 10: Quarkonium-like states at the open flavor thresholds. For charged states, the  $C$ -parity is given for the neutral members of the corresponding isotriplets.

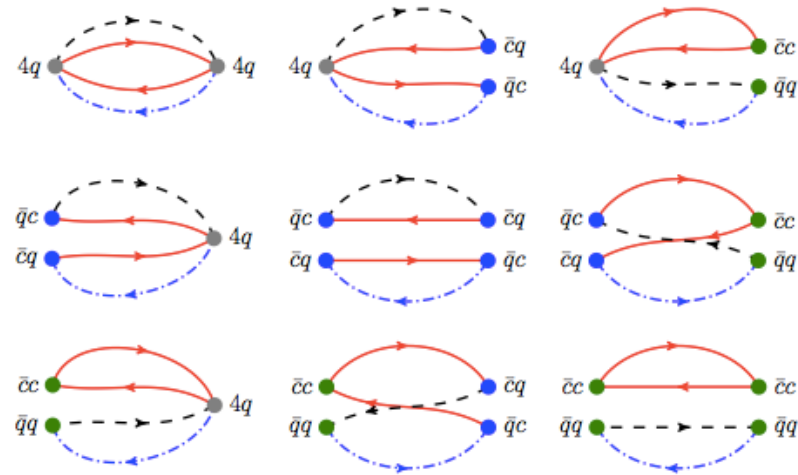
State	$M$ , MeV	$\Gamma$ , MeV	$J^{PC}$	Process (mode)	Experiment ( $\#\sigma$ )	Year	Status
$X(3872)$	$3871.68 \pm 0.17$	$< 1.2$	$1^{++}$	$B \rightarrow K(\pi^+\pi^-J/\psi)$	Belle [772, 992] ( $>10$ ), BaBar [993] (8.6)	2003	Ok
				$p\bar{p} \rightarrow (\pi^+\pi^-J/\psi) \dots$	CDF [994, 995] (11.6), D0 [996] (5.2)	2003	Ok
				$pp \rightarrow (\pi^+\pi^-J/\psi) \dots$	LHCb [997, 998] (np)	2012	Ok
				$B \rightarrow K(\pi^+\pi^-\pi^0J/\psi)$	Belle [999] (4.3), BaBar [1000] (4.0)	2005	Ok
				$B \rightarrow K(\gamma J/\psi)$	Belle [1001] (5.5), BaBar [1002] (3.5)	2005	Ok
					LHCb [1003] ( $> 10$ )		
				$B \rightarrow K(\gamma\psi(2S))$	BaBar [1002] (3.6), Belle [1001] (0.2)	2008	NC!
	LHCb [1003] (4.4)						
			$B \rightarrow K(D\bar{D}^*)$	Belle [1004] (6.4), BaBar [1005] (4.9)	2006	Ok	
$Z_c(3885)^+$	$3883.9 \pm 4.5$	$25 \pm 12$	$1^{+-}$	$Y(4260) \rightarrow \pi^-(D\bar{D}^*)^+$	BES III [1006] (np)	2013	NC!
$Z_c(3900)^+$	$3891.2 \pm 3.3$	$40 \pm 8$	$?^{? -}$	$Y(4260) \rightarrow \pi^-(\pi^+J/\psi)$	BES III [1007] (8), Belle [1008] (5.2)	2013	Ok
					T. Xiao <i>et al.</i> [CLEO data] [1009] ( $>5$ )		
$Z_c(4020)^+$	$4022.9 \pm 2.8$	$7.9 \pm 3.7$	$?^{? -}$	$Y(4260, 4360) \rightarrow \pi^-(\pi^+h_c)$	BES III [1010] (8.9)	2013	NC!
$Z_c(4025)^+$	$4026.3 \pm 4.5$	$24.8 \pm 9.5$	$?^{? -}$	$Y(4260) \rightarrow \pi^-(D^*\bar{D}^*)^+$	BES III [1011] (10)	2013	NC!
$Z_b(10610)^+$	$10607.2 \pm 2.0$	$18.4 \pm 2.4$	$1^{+-}$	$\Upsilon(10860) \rightarrow \pi(\pi\Upsilon(1S, 2S, 3S))$	Belle [1012–1014] ( $>10$ )	2011	Ok
				$\Upsilon(10860) \rightarrow \pi^-(\pi^+h_b(1P, 2P))$	Belle [1013] (16)	2011	Ok
				$\Upsilon(10860) \rightarrow \pi^-(B\bar{B}^*)^+$	Belle [1015] (8)	2012	NC!
$Z_b(10650)^+$	$10652.2 \pm 1.5$	$11.5 \pm 2.2$	$1^{+-}$	$\Upsilon(10860) \rightarrow \pi^-(\pi^+\Upsilon(1S, 2S, 3S))$	Belle [1012, 1013] ( $>10$ )	2011	Ok
				$\Upsilon(10860) \rightarrow \pi^-(\pi^+h_b(1P, 2P))$	Belle [1013] (16)	2011	Ok
				$\Upsilon(10860) \rightarrow \pi^-(B^*\bar{B}^*)^+$	Belle [1015] (6.8)	2012	NC!

[review: Brambilla et al., 1404.3723]

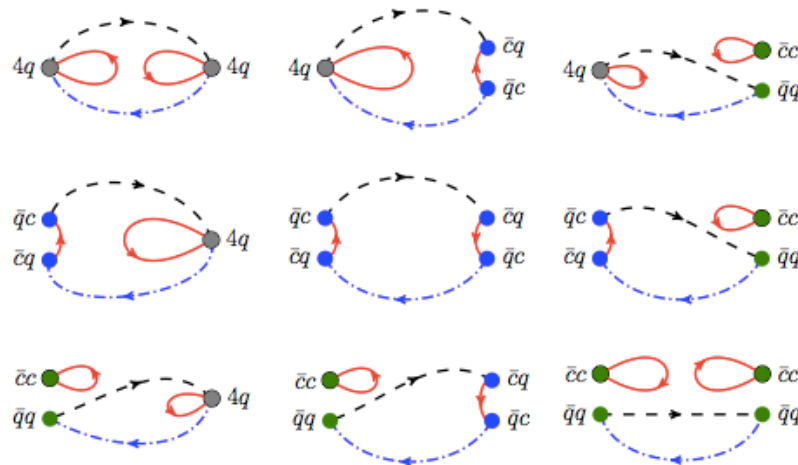
State	$M$ , MeV	$\Gamma$ , MeV	$J^{PC}$	Process (mode)	Experiment ( $\#\sigma$ )	Year	Status
$Y(3915)$	$3918.4 \pm 1.9$	$20 \pm 5$	$0/2^{7+}$	$B \rightarrow K(\omega J/\psi)$	Belle [1050] (8), BaBar [1000, 1051] (19)	2004	Ok
				$e^+e^- \rightarrow e^+e^-(\omega J/\psi)$	Belle [1052] (7.7), BaBar [1053] (7.6)	2009	Ok
$\chi_{c2}(2P)$	$3927.2 \pm 2.6$	$24 \pm 6$	$2^{++}$	$e^+e^- \rightarrow e^+e^-(D\bar{D})$	Belle [1054] (5.3), BaBar [1055] (5.8)	2005	Ok
$X(3940)$	$3942_{-8}^{+9}$	$37_{-17}^{+27}$	$?^{?+}$	$e^+e^- \rightarrow J/\psi(D\bar{D}^*)$	Belle [1048, 1049] (6)	2005	NC!
$Y(4008)$	$3891 \pm 42$	$255 \pm 42$	$1^{--}$	$e^+e^- \rightarrow (\pi^+\pi^- J/\psi)$	Belle [1008, 1056] (7.4)	2007	NC!
$\psi(4040)$	$4039 \pm 1$	$80 \pm 10$	$1^{--}$	$e^+e^- \rightarrow (D^{(*)}\bar{D}^{(*)}(\pi))$	PDG [1]	1978	Ok
				$e^+e^- \rightarrow (\eta J/\psi)$	Belle [1057] (6.0)	2013	NC!
$Z(4050)^+$	$4051_{-43}^{+24}$	$82_{-55}^{+51}$	$?^{?+}$	$\bar{B}^0 \rightarrow K^-(\pi^+\chi_{c1})$	Belle [1058] (5.0), BaBar [1059] (1.1)	2008	NC!
$Y(4140)$	$4145.8 \pm 2.6$	$18 \pm 8$	$?^{?+}$	$B^+ \rightarrow K^+(\phi J/\psi)$	CDF [1060] (5.0), Belle [1061] (1.9), LHCb [1062] (1.4), CMS [1063] (>5)	2009	NC!
					D0 [1064] (3.1)		
$\psi(4160)$	$4153 \pm 3$	$103 \pm 8$	$1^{--}$	$e^+e^- \rightarrow (D^{(*)}\bar{D}^{(*)})$	PDG [1]	1978	Ok
				$e^+e^- \rightarrow (\eta J/\psi)$	Belle [1057] (6.5)	2013	NC!
$X(4160)$	$4156_{-25}^{+29}$	$139_{-65}^{+113}$	$?^{?+}$	$e^+e^- \rightarrow J/\psi(D^*\bar{D}^*)$	Belle [1049] (5.5)	2007	NC!
$Z(4200)^+$	$4196_{-30}^{+35}$	$370_{-110}^{+99}$	$1^{+-}$	$\bar{B}^0 \rightarrow K^-(\pi^+ J/\psi)$	Belle [1065] (7.2)	2014	NC!
$Z(4250)^+$	$4248_{-45}^{+185}$	$177_{-72}^{+321}$	$?^{?+}$	$\bar{B}^0 \rightarrow K^-(\pi^+\chi_{c1})$	Belle [1058] (5.0), BaBar [1059] (2.0)	2008	NC!
$Y(4260)$	$4250 \pm 9$	$108 \pm 12$	$1^{--}$	$e^+e^- \rightarrow (\pi\pi J/\psi)$	BaBar [1066, 1067] (8), CLEO [1068, 1069] (11)	2005	Ok
					Belle [1008, 1056] (15), BES III [1007] (np)		
				$e^+e^- \rightarrow (f_0(980)J/\psi)$	BaBar [1067] (np), Belle [1008] (np)	2012	Ok
				$e^+e^- \rightarrow (\pi^- Z_c(3900)^+)$	BES III [1007] (8), Belle [1008] (5.2)	2013	Ok
				$e^+e^- \rightarrow (\gamma X(3872))$	BES III [1070] (5.3)	2013	NC!
$Y(4274)$	$4293 \pm 20$	$35 \pm 16$	$?^{?+}$	$B^+ \rightarrow K^+(\phi J/\psi)$	CDF [1060] (3.1), LHCb [1062] (1.0), CMS [1063] (>3), D0 [1064] (np)	2011	NC!
$X(4350)$	$4350.6_{-5.1}^{+4.6}$	$13_{-10}^{+18}$	$0/2^{7+}$	$e^+e^- \rightarrow e^+e^-(\phi J/\psi)$	Belle [1071] (3.2)	2009	NC!
$Y(4360)$	$4354 \pm 11$	$78 \pm 16$	$1^{--}$	$e^+e^- \rightarrow (\pi^+\pi^-\psi(2S))$	Belle [1072] (8), BaBar [1073] (np)	2007	Ok
$Z(4430)^+$	$4458 \pm 15$	$166_{-32}^{+37}$	$1^{+-}$	$\bar{B}^0 \rightarrow K^-(\pi^+\psi(2S))$	Belle [1074, 1075] (6.4), BaBar [1076] (2.4)	2007	Ok
					LHCb [1077] (13.9)		
				$\bar{B}^0 \rightarrow K^-(\pi^+ J/\psi)$	Belle [1065] (4.0)	2014	NC!
$X(4630)$	$4634_{-11}^{+9}$	$92_{-32}^{+41}$	$1^{--}$	$e^+e^- \rightarrow (\Lambda_c^+ \bar{\Lambda}_c^-)$	Belle [1078] (8.2)	2007	NC!
$Y(4660)$	$4665 \pm 10$	$53 \pm 14$	$1^{--}$	$e^+e^- \rightarrow (\pi^+\pi^-\psi(2S))$	Belle [1072] (5.8), BaBar [1073] (5)	2007	Ok
$\Upsilon(10860)$	$10876 \pm 11$	$55 \pm 28$	$1^{--}$	$e^+e^- \rightarrow (B_{(s)}^{(*)}\bar{B}_{(s)}^{(*)}(\pi))$	PDG [1]	1985	Ok
				$e^+e^- \rightarrow (\pi\pi\Upsilon(1S, 2S, 3S))$	Belle [1013, 1014, 1079] (>10)	2007	Ok
				$e^+e^- \rightarrow (f_0(980)\Upsilon(1S))$	Belle [1013, 1014] (>5)	2011	Ok
				$e^+e^- \rightarrow (\pi Z_b(10610, 10650))$	Belle [1013, 1014] (>10)	2011	Ok
				$e^+e^- \rightarrow (\eta\Upsilon(1S, 2S))$	Belle [948] (10)	2012	Ok
				$e^+e^- \rightarrow (\pi^+\pi^-\Upsilon(1D))$	Belle [948] (9)	2012	Ok
$Y_b(10888)$	$10888.4 \pm 3.0$	$30.7_{-7.7}^{+8.9}$	$1^{--}$	$e^+e^- \rightarrow (\pi^+\pi^-\Upsilon(nS))$	Belle [1080] (2.3)	2008	NC!

# Wick contractions for $Zc+$

a



b

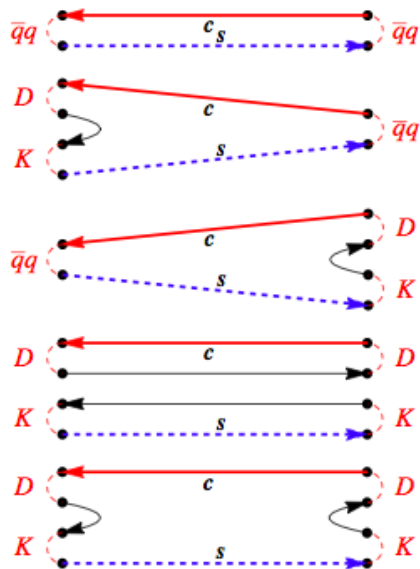


# $D_{s0}^*(2317)$ with $J^P=0^+$

$$\mathcal{O} = \bar{s} c$$

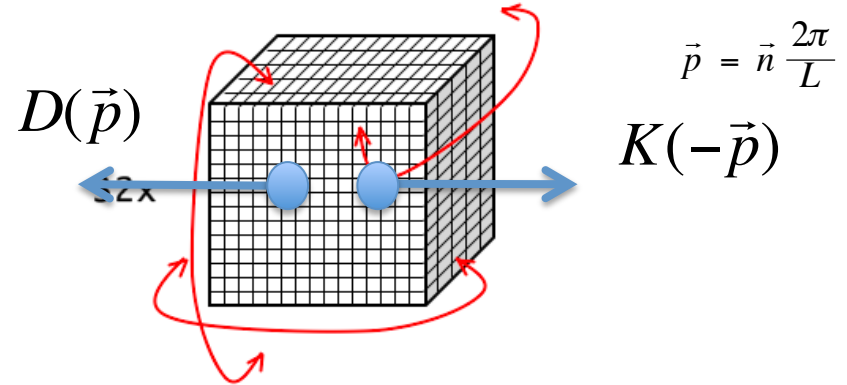
$$\mathcal{O} = DK \approx [\bar{d}\gamma_5 c] [\bar{s}\gamma_5 d]$$

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^+(0) | 0 \rangle$$



We use distillation method

[Peardon et al. 2009] to calculate  $C_{ij}$



$$\vec{p} = \vec{n} \frac{2\pi}{L}$$

Extract  $E_n$  from  $C_{ij}(t)$ : variational method

$$C_{ij}(t) = \sum_n A_n^{ij} e^{-E_n t}$$

due to strong int.

$$\vec{p} = \vec{n} \frac{2\pi}{L} \quad E(L) = \sqrt{m_D^2 + \vec{p}^2} + \sqrt{m_K^2 + (-\vec{p})^2} + \Delta E$$

Energy levels that appear in addition to these discrete two-particle states correspond to bound states or resonances

# Meson-meson scattering: extracting $\delta(p)$ from $E_n$ at $P=p_1+p_2=0$ [M. Luscher,1991]

- extract  $E_n(L)$
- $E_n$  renders  $p$  in the region "outside-interaction"

$$E = \sqrt{s} = \sqrt{m_1^2 + p^2} + \sqrt{m_2^2 + p^2}$$

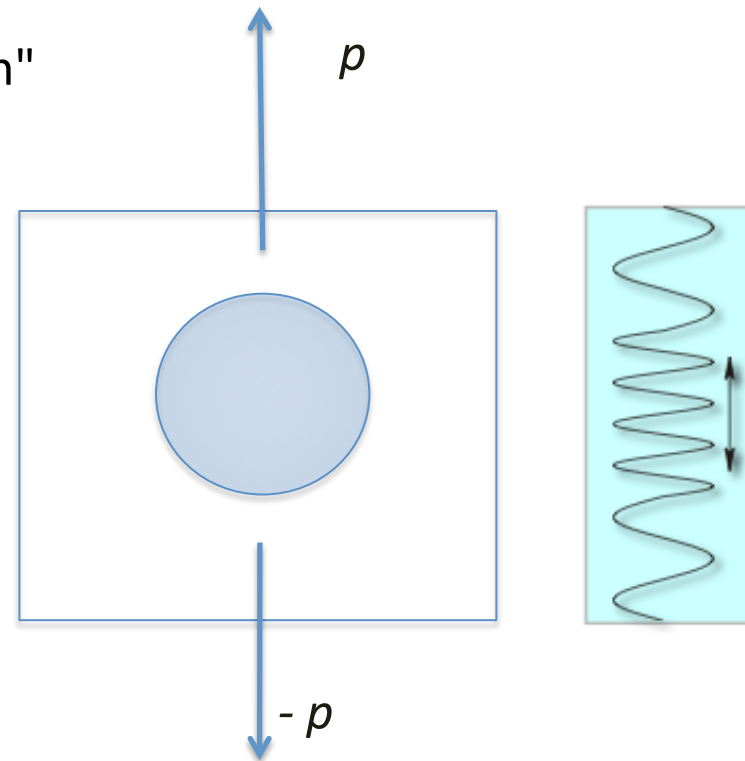
- $p$  contains info on  $\delta(p)$  for periodic BC

$$\det[\delta_{ll'}\delta_{mm'} - \tan\delta_l(p) M_{ll',mm'}(q^2)] = 0$$

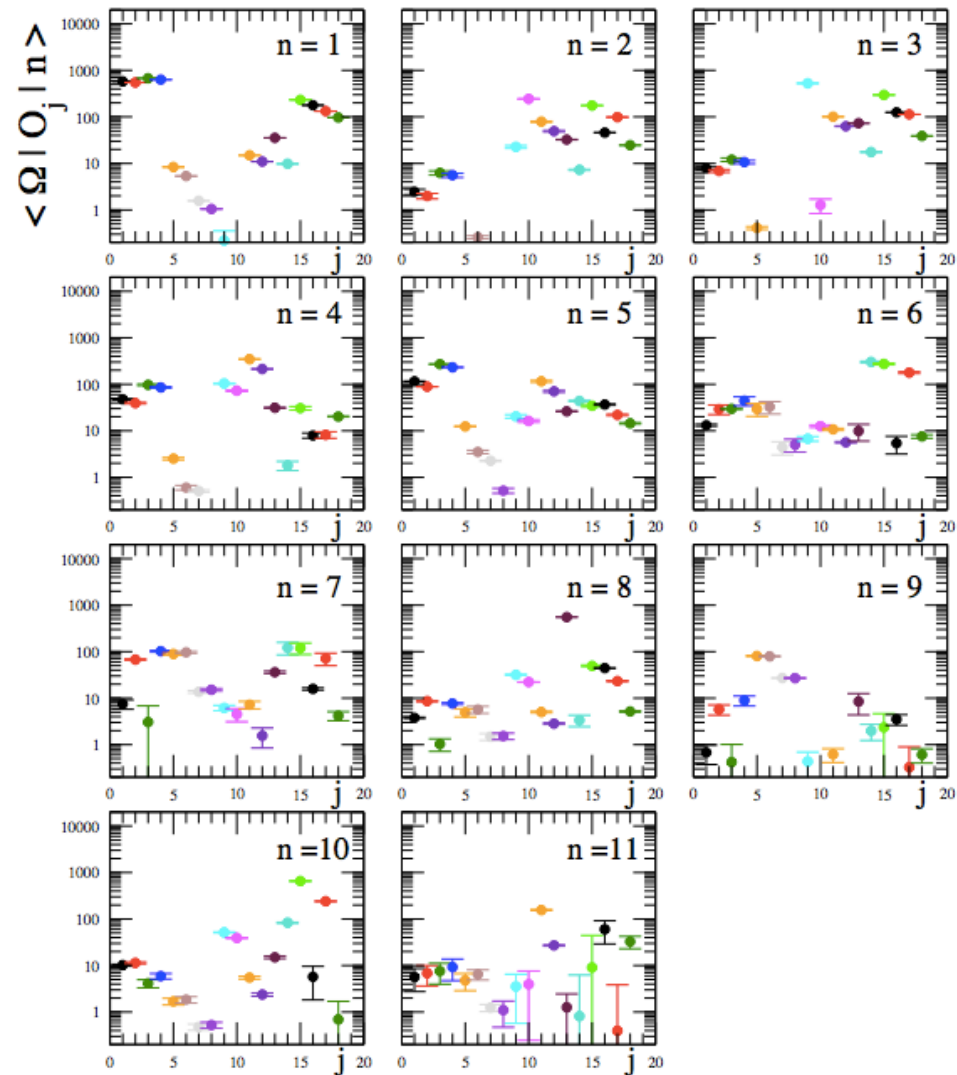
$$\tan\delta(p) = \frac{\pi^{3/2}q}{Z_{00}(1;q^2)}$$

$$Z_{00}(1;q^2) \equiv \sum_{\vec{n} \in N^3} \frac{1}{\vec{n}^2 - q^2}$$

$$q \equiv \frac{L}{2\pi} p$$



# Overlaps of all states in Zc+ channel



## Interpolators in Zc channel

$$\mathcal{O}_1 = \mathcal{O}_1^{\psi(0)\pi(0)} = \bar{c}\gamma_i c(0) \bar{d}\gamma_5 u(0),$$

$$\mathcal{O}_2 = \mathcal{O}_2^{\psi(0)\pi(0)} = \bar{c}\gamma_i \gamma_t c(0) \bar{d}\gamma_5 u(0),$$

$$\mathcal{O}_3 = \mathcal{O}_3^{\psi(0)\pi(0)} = \bar{c} \overleftarrow{\nabla}_j \gamma_i \overrightarrow{\nabla}_j c(0) \bar{d}\gamma_5 u(0),$$

$$\mathcal{O}_4 = \mathcal{O}_4^{\psi(0)\pi(0)} = \bar{c} \overleftarrow{\nabla}_j \gamma_i \gamma_t \overrightarrow{\nabla}_j c(0) \bar{d}\gamma_5 u(0),$$

$$\mathcal{O}_5 = \mathcal{O}_5^{\psi(0)\pi(0)} = |\epsilon_{ijk}| |\epsilon_{klm}| \bar{c}\gamma_j \overleftarrow{\nabla}_l \overrightarrow{\nabla}_m c(0) \bar{d}\gamma_5 u(0),$$

$$\mathcal{O}_6 = \mathcal{O}_6^{\psi(0)\pi(0)} = |\epsilon_{ijk}| |\epsilon_{klm}| \bar{c}\gamma_t \gamma_j \overleftarrow{\nabla}_l \overrightarrow{\nabla}_m c(0) \bar{d}\gamma_5 u(0),$$

$$\mathcal{O}_7 = \mathcal{O}_7^{\psi(0)\pi(0)} = R_{ijk} Q_{klm} \bar{c}\gamma_j \overleftarrow{\nabla}_l \overrightarrow{\nabla}_m c \bar{d}\gamma_5 u(0),$$

$$\mathcal{O}_8 = \mathcal{O}_8^{\psi(0)\pi(0)} = R_{ijk} Q_{klm} \bar{c}\gamma_t \gamma_j \overleftarrow{\nabla}_l \overrightarrow{\nabla}_m c \bar{d}\gamma_5 u(0).$$

$$\mathcal{O}_9 = \mathcal{O}^{\psi(1)\pi(-1)} = \sum_{e_k=\pm e_{x,y,z}} \bar{c}\gamma_i c(e_k) \bar{d}\gamma_5 u(-e_k),$$

$$\mathcal{O}_{10} = \mathcal{O}^{\eta_c(0)\rho(0)} = \bar{c}\gamma_5 c(0) \bar{d}\gamma_i u(0),$$

$$\mathcal{O}_{11} = \mathcal{O}_1^{D(0)D^*(0)} = \bar{c}\gamma_5 u(0) \bar{d}\gamma_i c(0) + \{\gamma_5 \leftrightarrow \gamma_i\},$$

$$\mathcal{O}_{12} = \mathcal{O}_2^{D(0)D^*(0)} = \bar{c}\gamma_5 \gamma_t u(0) \bar{d}\gamma_i \gamma_t c(0) + \{\gamma_5 \leftrightarrow \gamma_i\},$$

$$\mathcal{O}_{13} = \mathcal{O}^{D(1)D^*(-1)} = \sum_{e_k=\pm e_{x,y,z}} \bar{c}\gamma_5 u(e_k) \bar{d}\gamma_i c(-e_k) + \{\gamma_5 \leftrightarrow \gamma_i\},$$

$$\mathcal{O}_{14} = \mathcal{O}^{D^*(0)D^*(0)} = \epsilon_{ijk} \bar{c}\gamma_j u(0) \bar{d}\gamma_k c(0),$$

$$\mathcal{O}_{15} = \mathcal{O}_1^{4q} = N_L^3 \epsilon_{abc} \epsilon_{ab'c'} (\bar{c}_b C \gamma_5 \bar{d}_c c_{b'} \gamma_i C u_{c'} - \bar{c}_b C \gamma_i \bar{d}_c c_{b'} \gamma_5 C u_{c'})$$

$$\mathcal{O}_{16} = \mathcal{O}_2^{4q} = N_L^3 \epsilon_{abc} \epsilon_{ab'c'} (\bar{c}_b C \bar{d}_c c_{b'} \gamma_i \gamma_5 C u_{c'} - \bar{c}_b C \gamma_i \gamma_5 \bar{d}_c c_{b'} C u_{c'})$$

$$\mathcal{O}_{17} = \mathcal{O}_3^{4q} = \mathcal{O}_1^{4q} (N_v = 32),$$

$$\mathcal{O}_{18} = \mathcal{O}_4^{4q} = \mathcal{O}_2^{4q} (N_v = 32).$$