

Di-electron production in dp collisions at $E_{\text{kin}} = 2.5 \text{ GeV}$

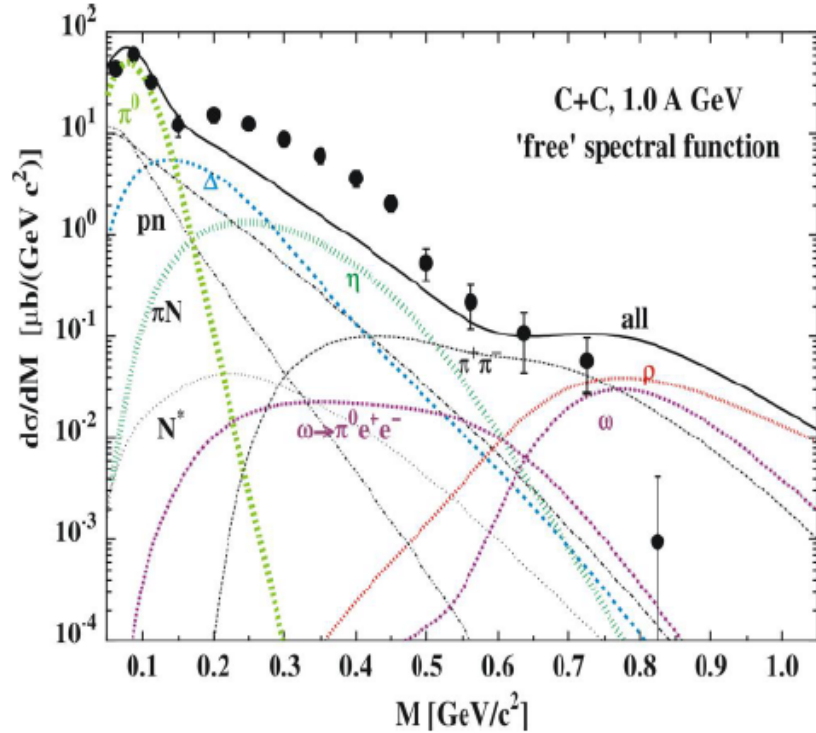
Jacek Biernat



DLS puzzle

Inclusive spectra

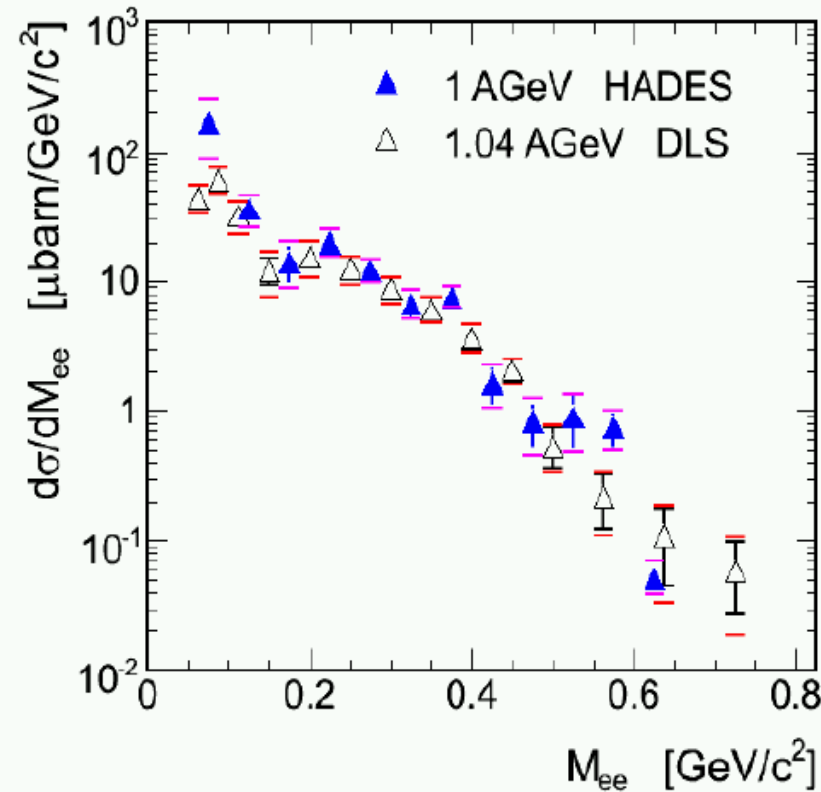
before HADES (<2003)



DLS data: R.J. Porter et al.: PRL. 79 (1997) 1229

Calculation: E.L.Bratkovskaya et al. PLB445 (1999) 265

2004 HADES re- measures C+C collisions



HADES data: Agakishiev et al., PLB 663 (2008) 43

Main contributing sources:

- **π⁰ Dalitz:**

$$\pi^0 \rightarrow \gamma e^+ e^-$$

- **η Dalitz:**

$$\eta \rightarrow \gamma e^+ e^-$$

- **Δ Dalitz:**

$$\Delta \rightarrow N e^+ e^-$$

- **Bremsstrahlung?**

$$NN \rightarrow NN e^+ e^-$$

- **Vector meson decay** $\rightarrow e^+ e^-$

✓ HADES fully confirms DLS results

- What about „excess”? Is it true in-medium effect?

- or not understood elementary process?

N+N and pi-N collisions in HADES

p+p collisions @

- $T_k = 1,25$ GeV
- $T_k = 2,2$ GeV
- $T_k = 3,5$ GeV

di-electron production

- Inclusive $e+e^-$ from pp and np

Phys.Rev. C85 (2012) 054005

Eur.Phys.J. A48 (2012) 64

Phys. Lett. B690 (2010) 118.

- **$np \rightarrow npe+e^-$**

- $pp \rightarrow ppe+e^-$

Eur.Phys.J. A50 (2014) 82

- $np \rightarrow de+e^-$ H. Kuc Thesis

- $\pi^-p \rightarrow ne+e^-$

quasi-free n+p collisions @

- $T_k = 1,25$ GeV

one-pion production

- $np \rightarrow np\pi$

- $pp \rightarrow pp\pi$

Eur.Phys.J. A50 (2014) 82

Eur.Phys.J. A48 (2012) 74

two-pion production

- $np \rightarrow np \pi+\pi^-$

Phys. Lett B 750, 12 (2015)

- $pp \rightarrow pp \pi+\pi^-$

on-going

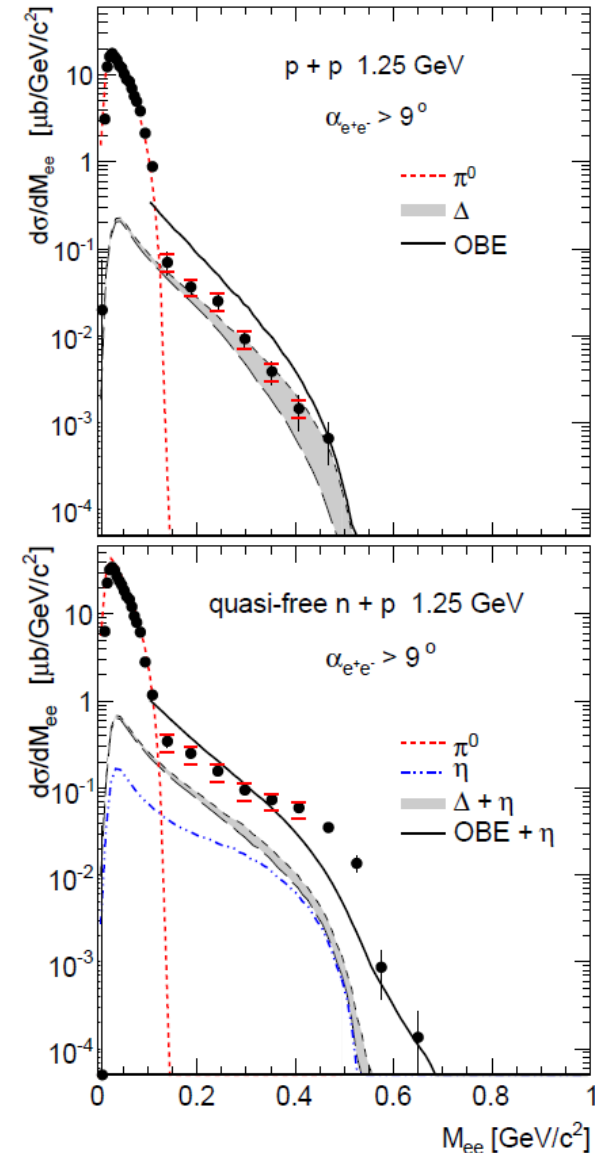
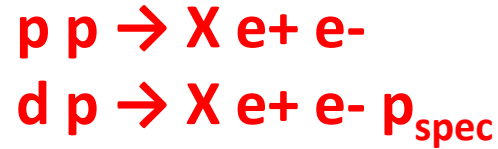
- $np \rightarrow d\pi+\pi^-$

- $\pi^-p \rightarrow n \pi+\pi^-$

Data taken in summer 2014

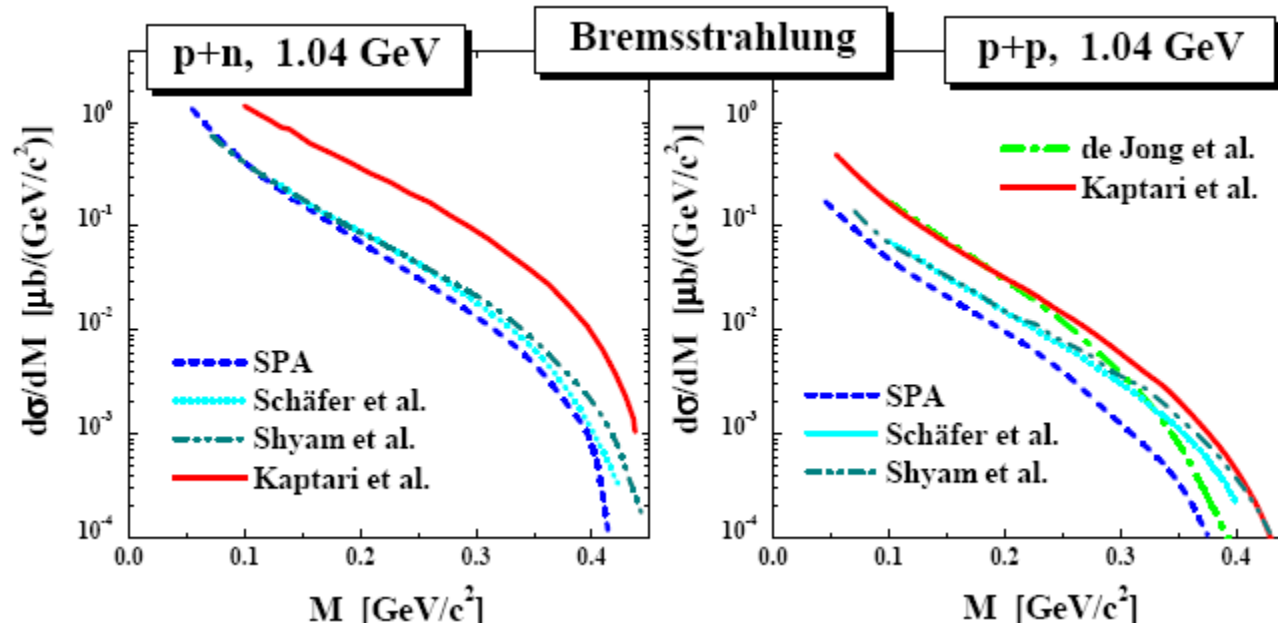
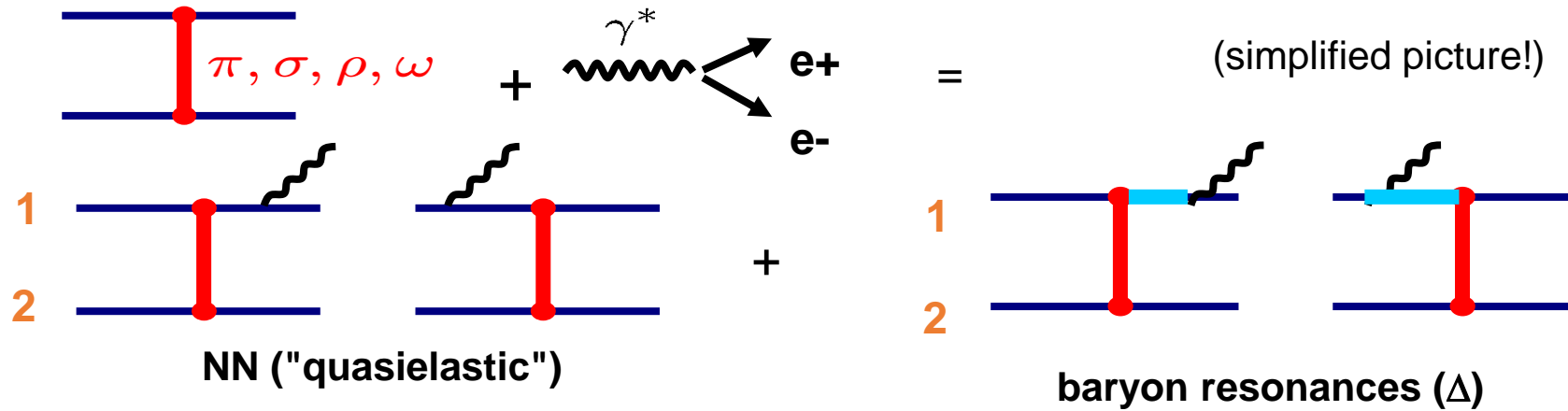
What have we learnt from inclusive spectra

- The pp spectra are well described by resonance model ($N_{\Delta}=3/2 N_{\pi^0}$) based on known cross sections. NOT described by OBE with increased bremsstrahlung contribution (see next slide)
- pn data are underestimated by the resonance model and also not described by OBE.
- general difference between pp and np reactions is the different Bremsstrahlung contribution and eta contribution. (OBE + η)
- none of the contributions could explain the enhancement in the di-lepton yield in np.



N-N Bremsstrahlung

- Strong + electromagnetic process (OBE models)



• E.L Bratkovskaya & W. Cassing:
arXiv: 0712.0635v1

- bremsstrahlung OBE calculations:

Kaptari & Kämpfer, NPA 764 (2006) 338:

new OBE calculation:

**pn bremsstrahlung 4
larger than in earlier
(<2000) calculations !**

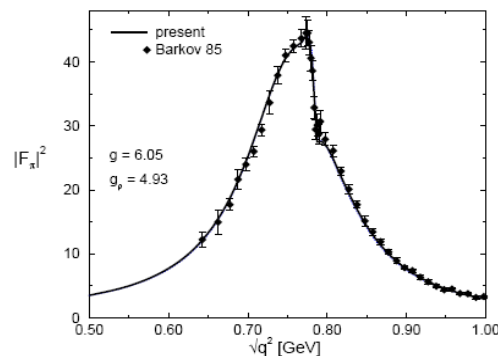
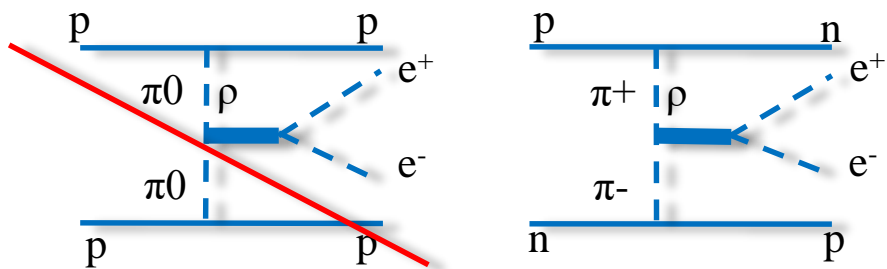
Possible explanation of e⁺e⁻ excess in np (I)

Possible explanation:

e⁺e⁻ excess in np

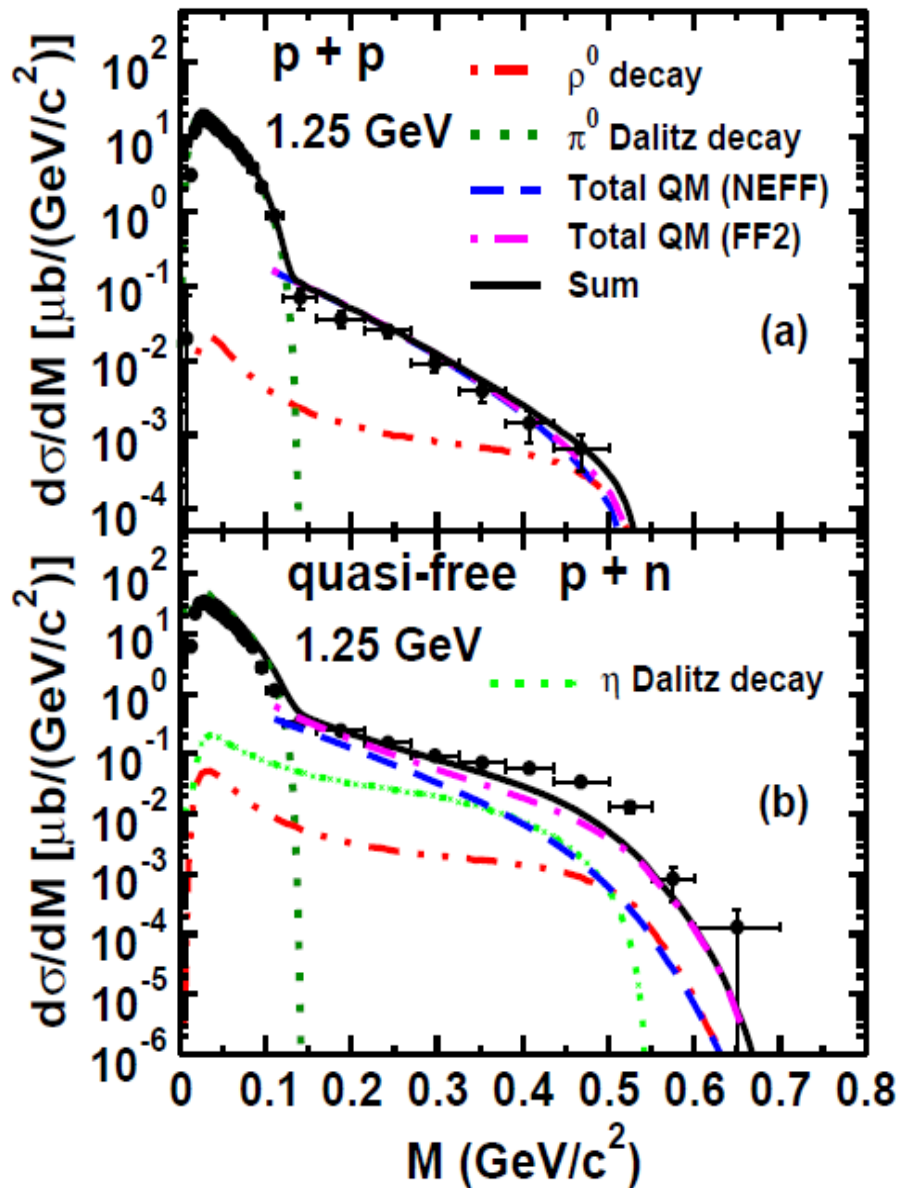
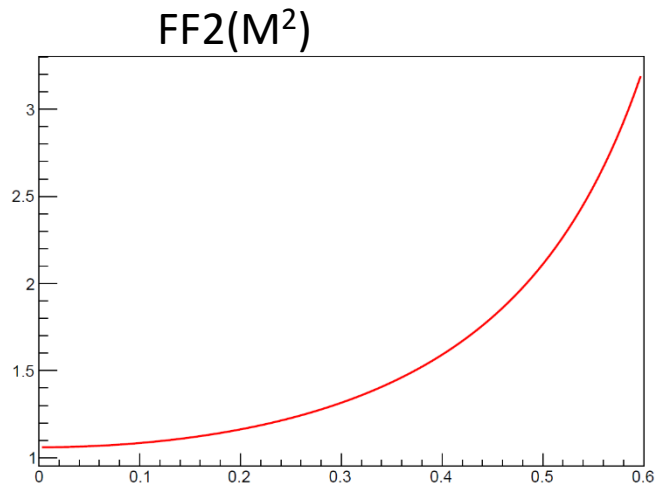
Introducing charged pion FF ?

R. Shyam , U. Mosel, Phys.Rev. C82 (2010) 062201



FF2 ➡
$$F_{\pi}(M^2) = \frac{0.4}{1 - M^2/\lambda^2} + \left(\frac{0.6}{1 - M^2/2m_{\rho}^2} \right) \left(\frac{m_{\rho}^2}{m_{\rho}^2 - M^2 - im_{\rho}\Gamma_{\rho}(M^2)} \right)$$

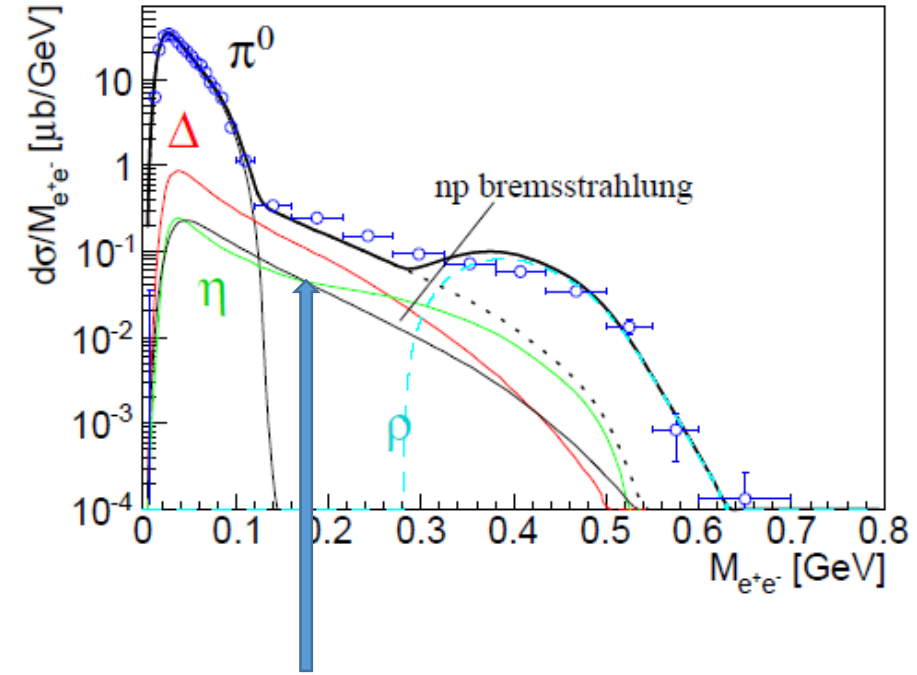
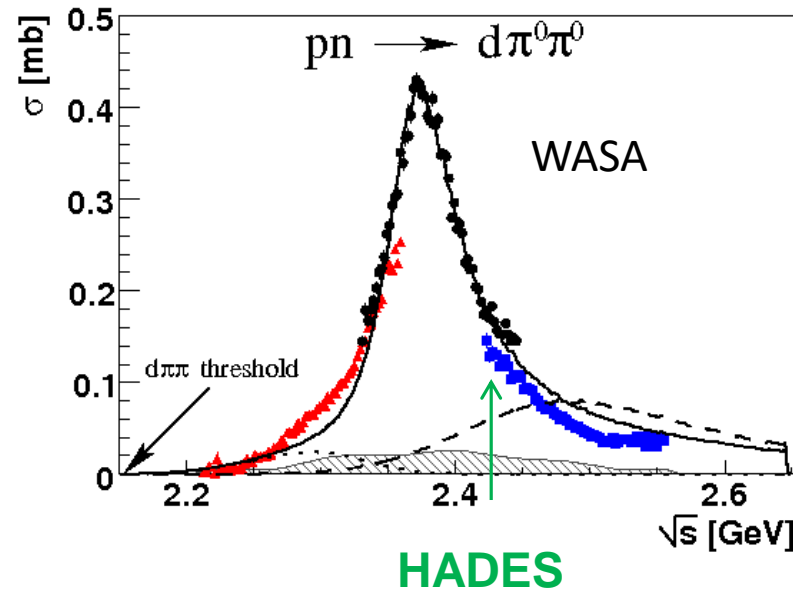
$m_{\rho} = 0.760 \text{ GeV}/c^2$
 $\lambda = 1.9 \text{ GeV}/c^2$



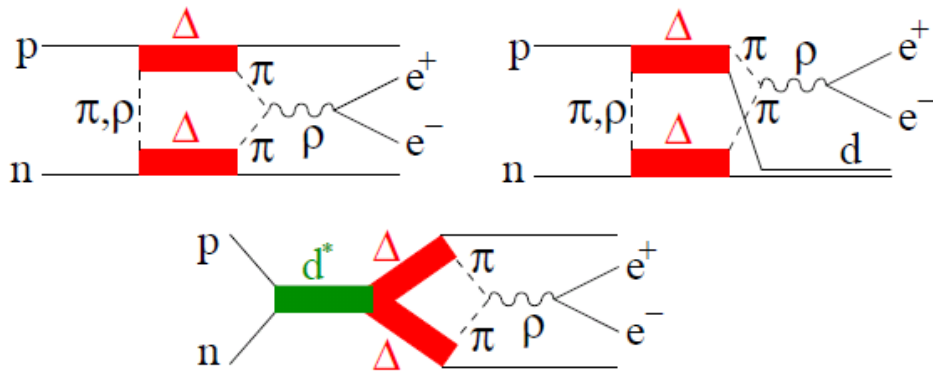
Possible explanation of e+e- excess in np (II)

off-shell ρ contribution in $\Delta\Delta$ interactions

M. Bashkanov and H. Clement
Eur.Phys.J. A50 (2014) 107



$$pn \rightarrow \Delta\Delta \rightarrow pn[\pi^+\pi^-]_{I=L=1} \rightarrow pn\rho^0 \rightarrow pne^+e^-,$$



d* info:
 $I(J^P) = 0(3^+)$
 $M = 2.37 \text{ GeV}/c^2$
 $\Gamma = 70 \text{ MeV}$

**0.15 GeV < M_{e+e-} < 0.3 GeV
 still a slightly underestimated
 region maybe due to**

$$pn \rightarrow d^* \rightarrow [pn]_{I=0} e^+ e^-$$

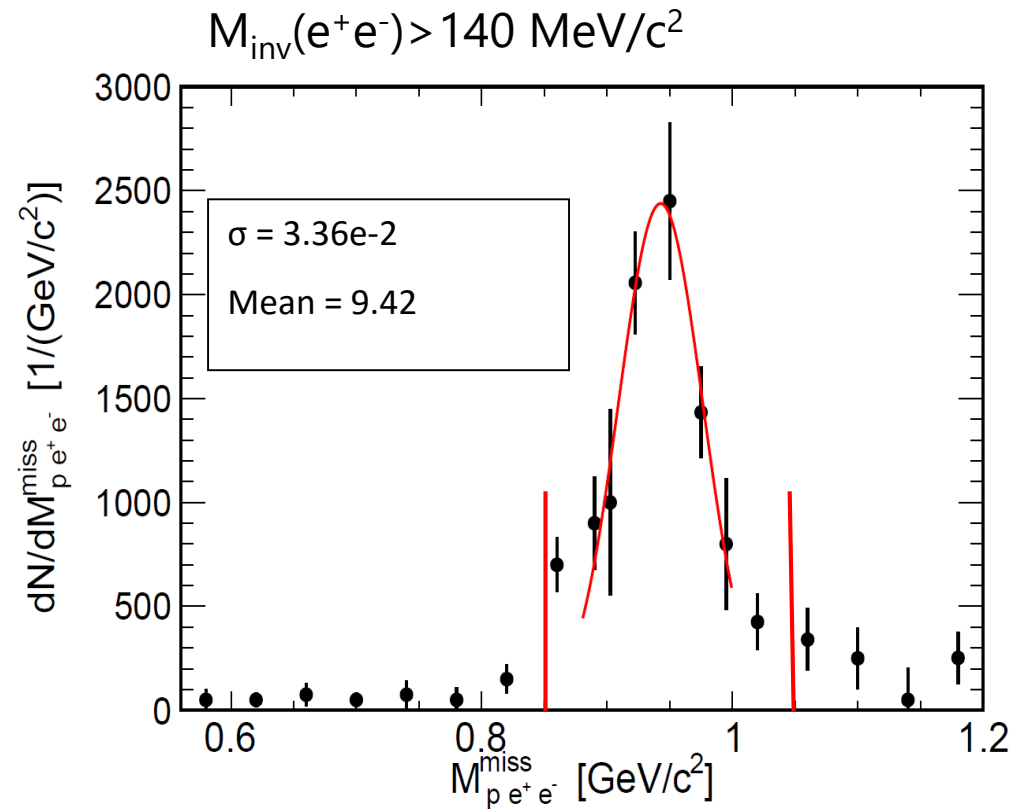
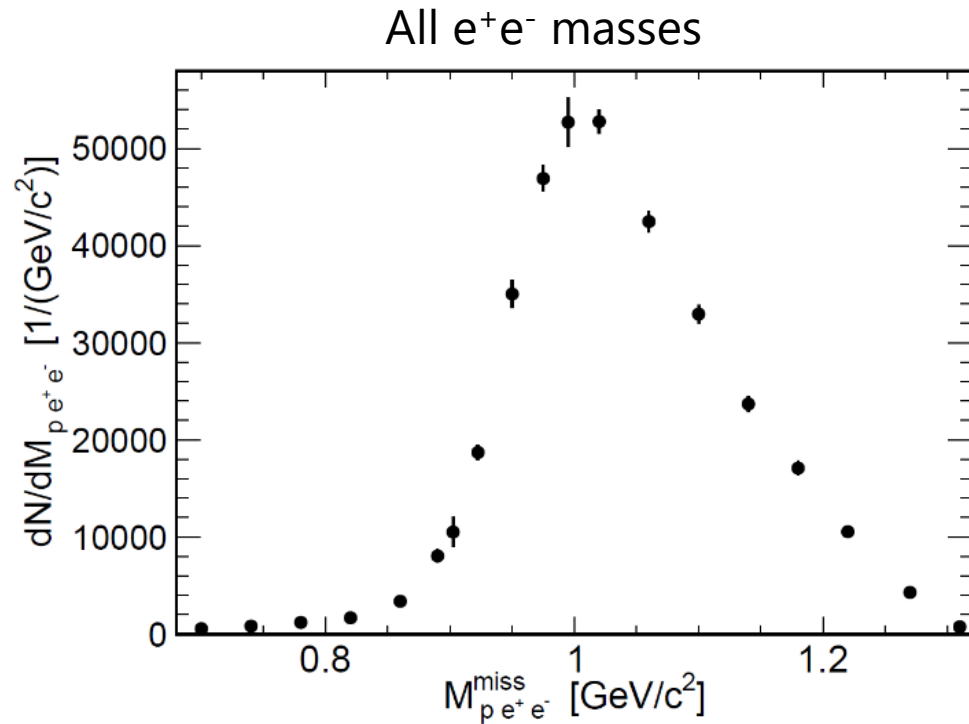
$$pn \rightarrow d^* \rightarrow de^+ e^-$$

$$|\mathcal{M}(\pi^+\pi^- \rightarrow \rho^0 \rightarrow e^+e^-)|^2 = \frac{m_\rho^2 \Gamma_{\pi^+\pi^-} \Gamma_{e^+e^-}}{(s - m_\rho^2)^2 + m_\rho^2 \Gamma_\rho^2}$$

Transition form $\pi^+\pi^-$ to e^+e^-

Identification of exclusive channel $(np)_{p_spec} \rightarrow (npe^+e^-)_{p_spec}$

- Three particle trigger selected (in case of np)
- Tagged proton in Forward Wall detector (spectator)
- 3 particles (proton, e^+e^-) identified in HADES
- selection via missing mass window

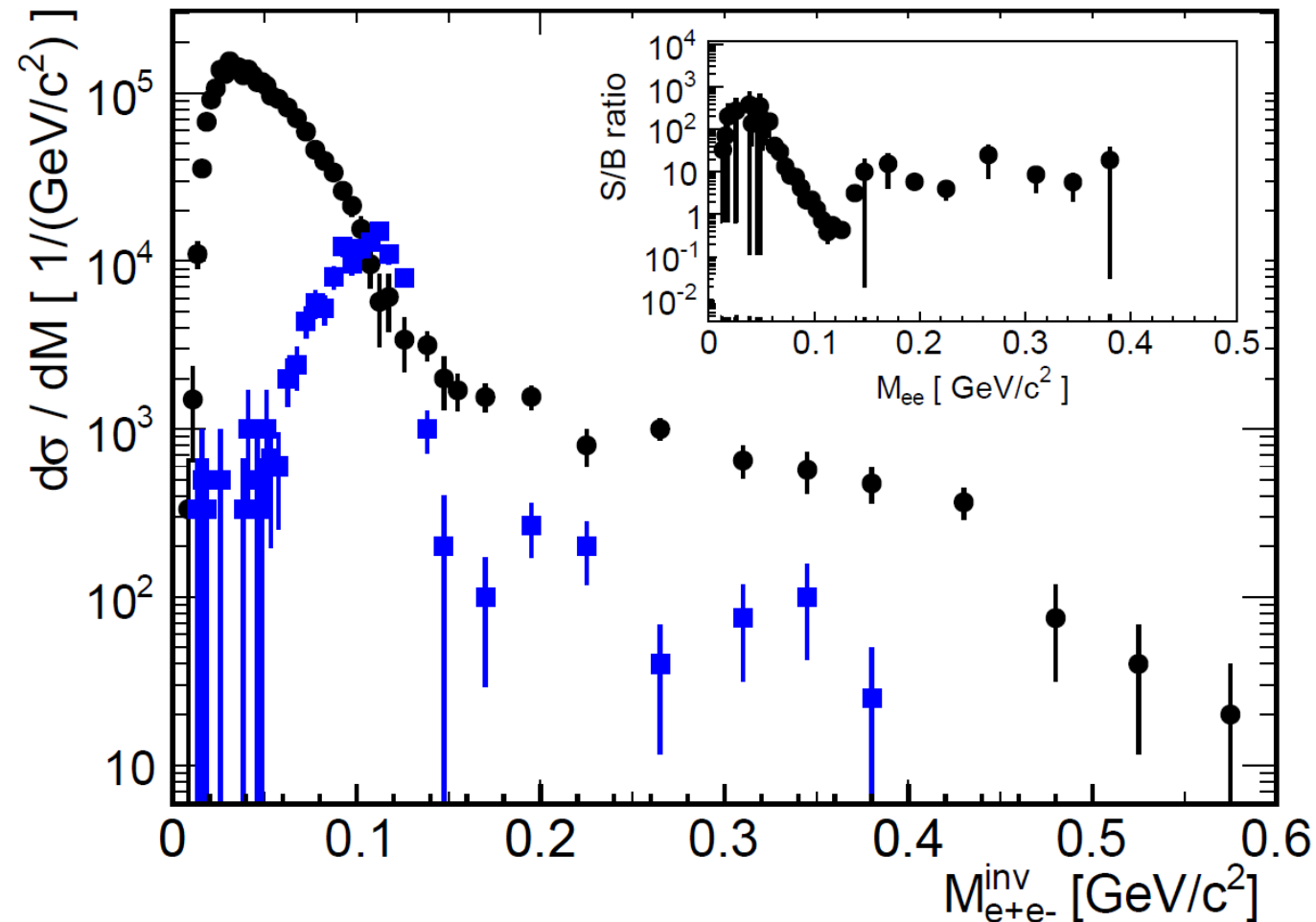


Unlike-sign combinatorial background estimation

The unlike-sign combinatorial background can be estimated by the reconstructed like-sign distribution.

$$N_{\text{CB}} = 2 * \sqrt{N_{++} * N_{--}}$$
$$N_{\text{sig_reco}} = N_{\text{sig}} - N_{\text{CB}}$$

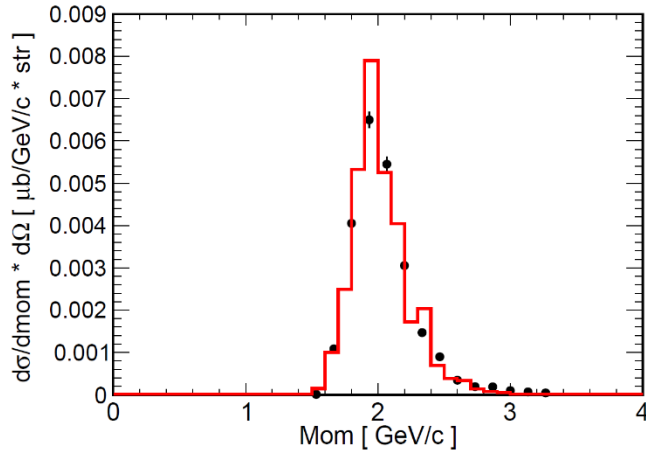
Above 140 MeV/c² background is negligible



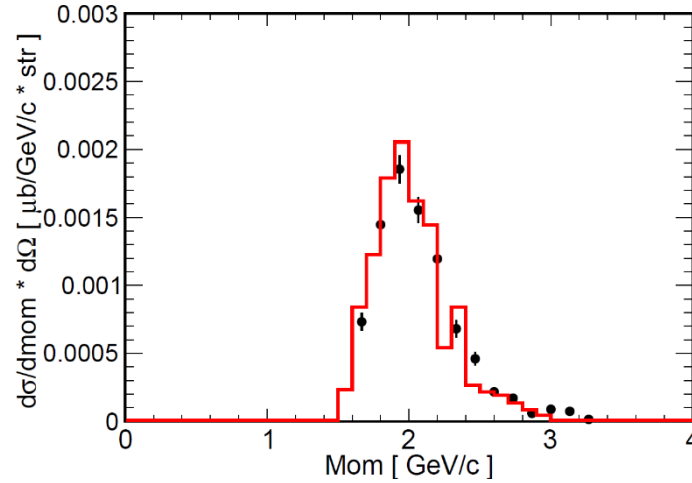
Comparison of spectator momentum distributions with simulation

$\theta < 2$ deg

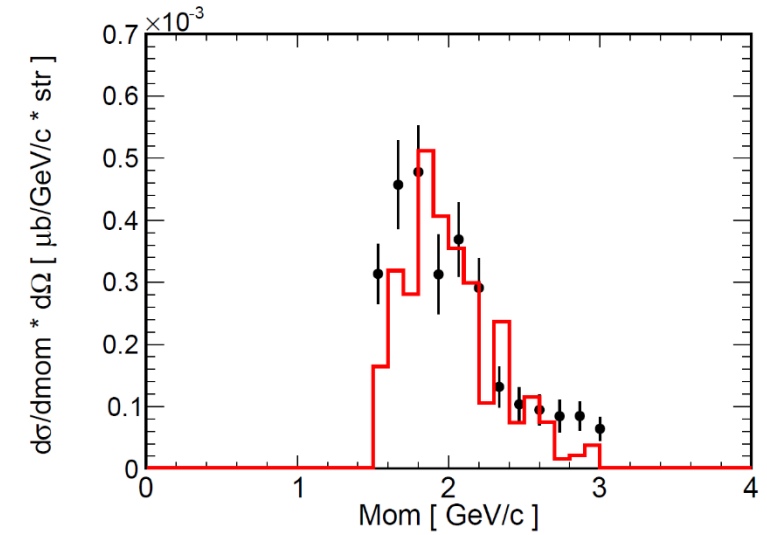
$M_{inv}(e^+e^-) < 140 \text{ MeV}/c^2$



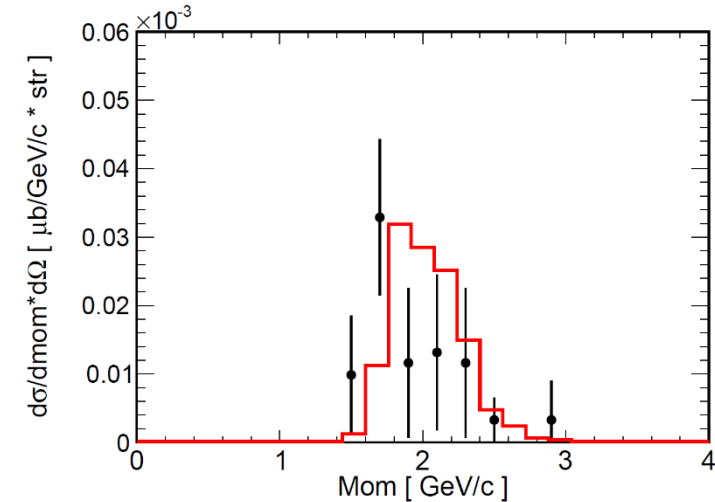
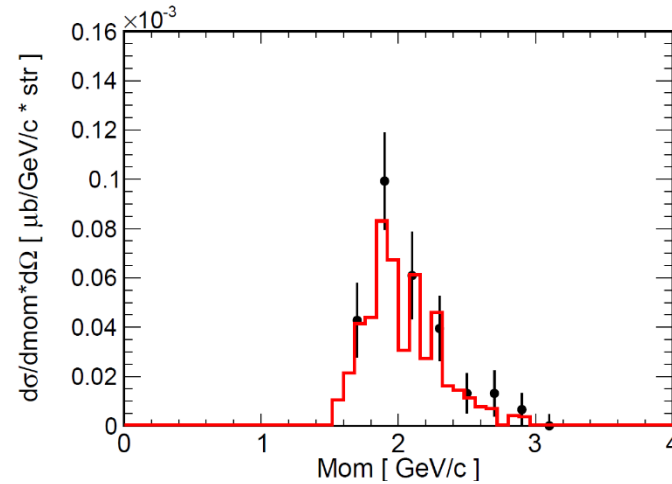
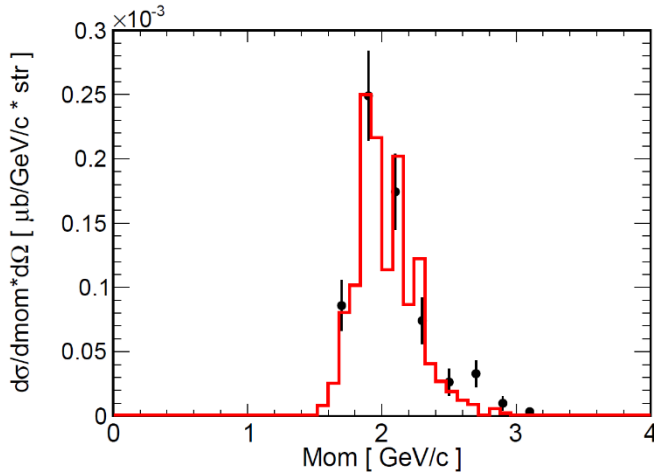
$2 < \theta < 4$



$4 < \theta < 6$



$M_{inv}(e^+e^-) > 140 \text{ MeV}/c^2$



Very good agreement in all mass range !

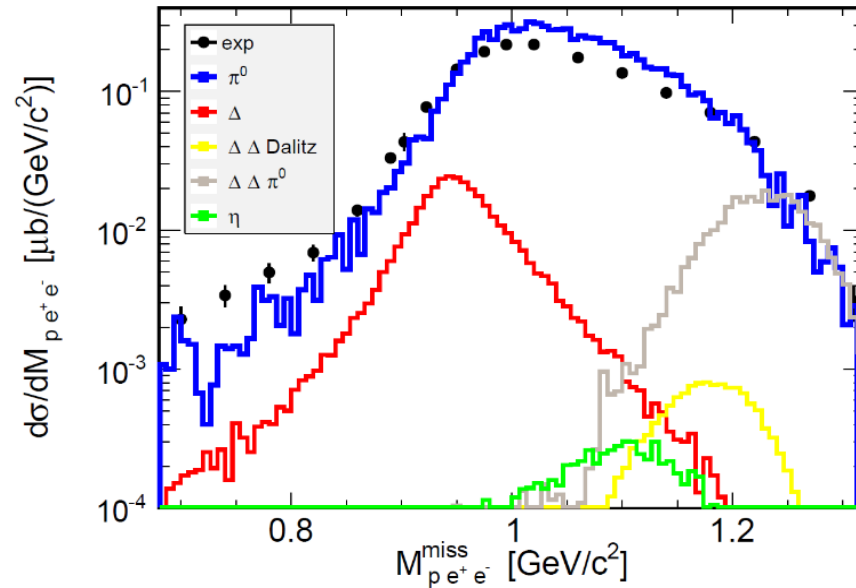
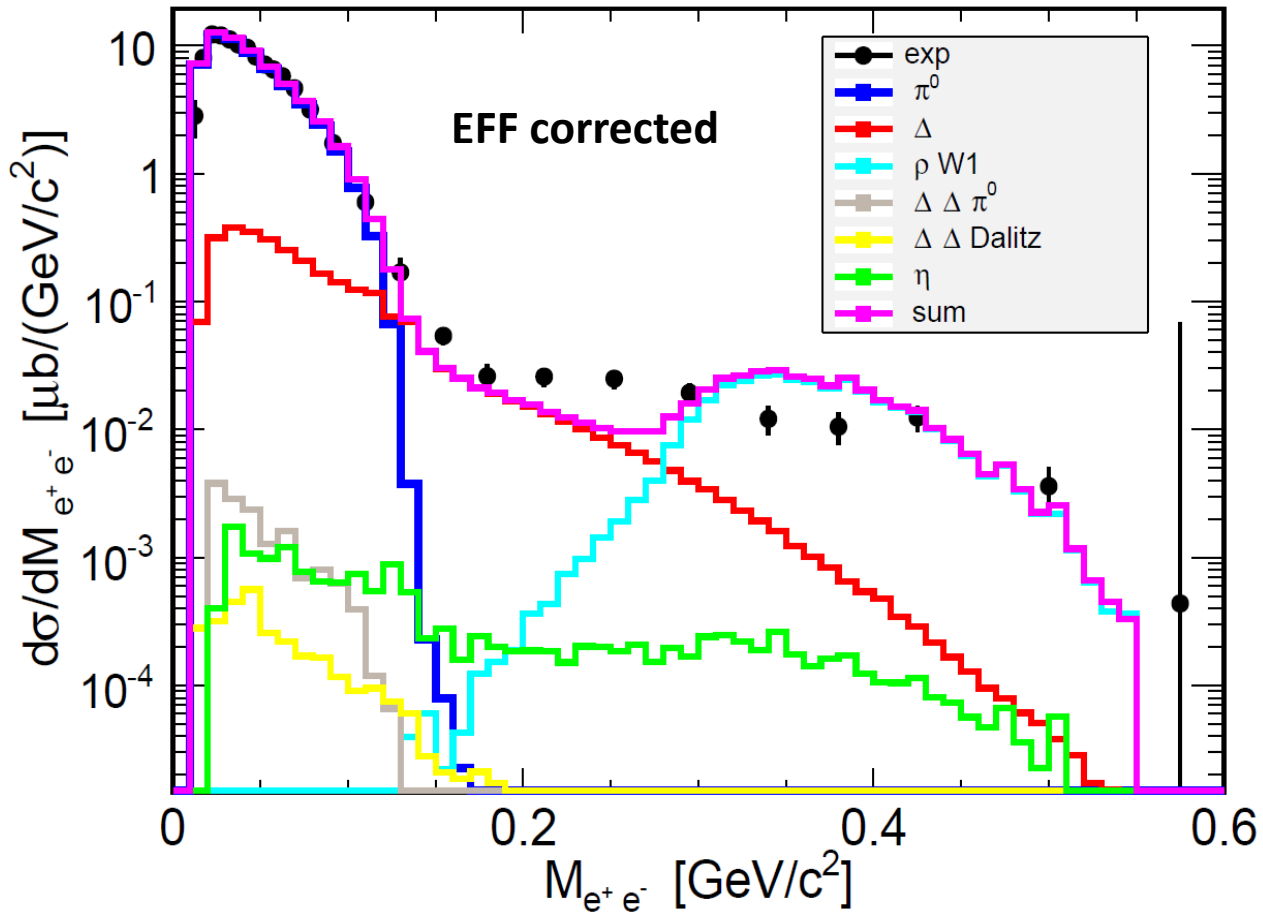
Comparison to models

Resonance model + rho contribution from Clement & Bashkanov:

Obtained form authors in a event by event form.

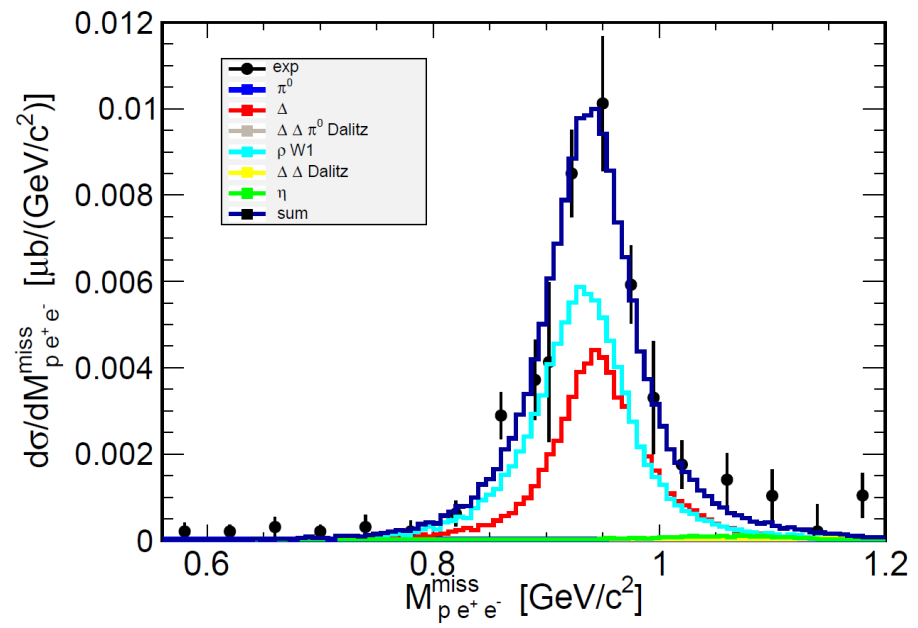
Total exclusive cross section is 210 μb :

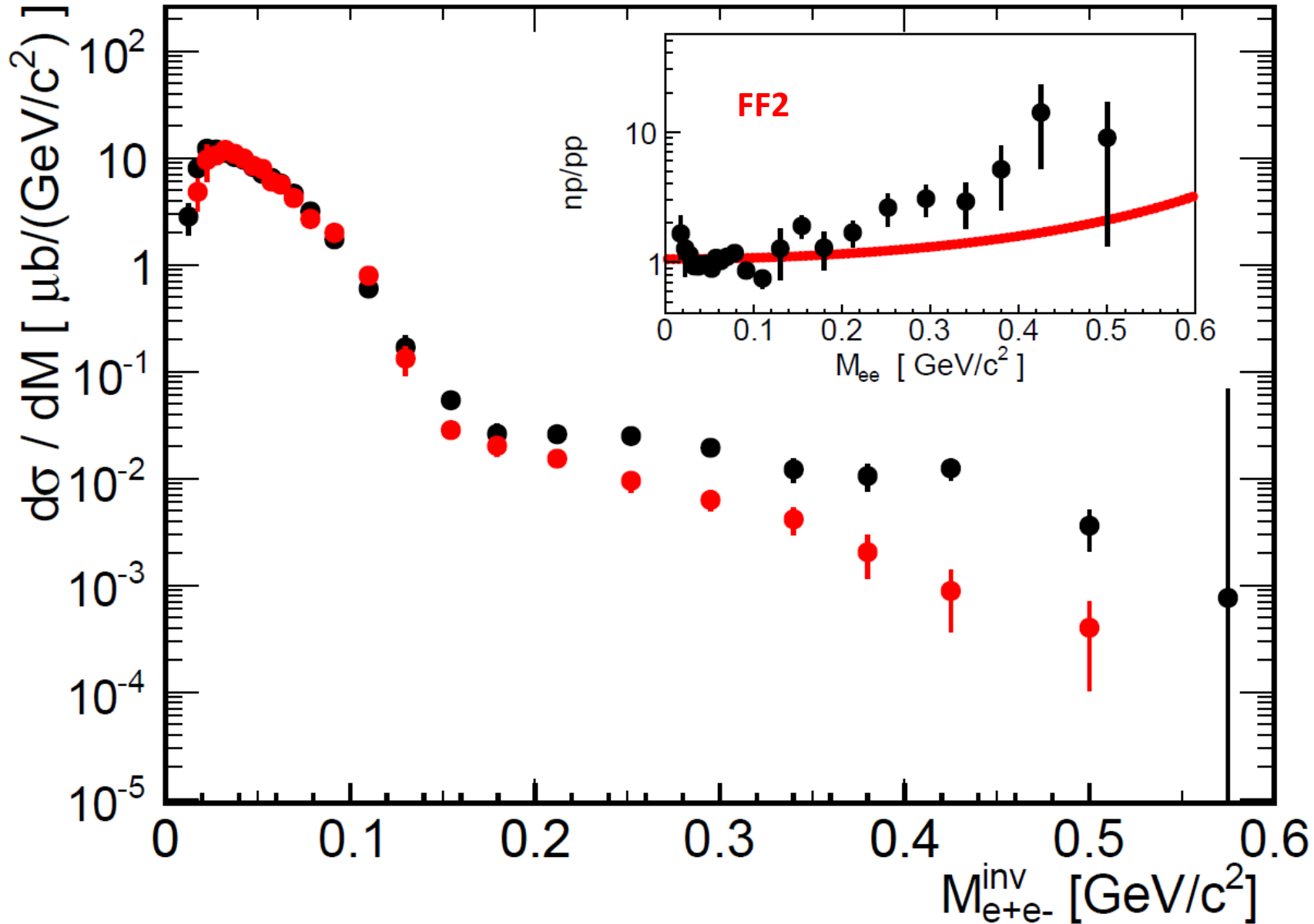
1. $np \rightarrow \Delta\Delta \rightarrow npp$ $\sigma = 170 \mu\text{b}$
2. $np \rightarrow d^* \rightarrow npp$ $\sigma = 40 \mu\text{b}$



All e^+e^- masses

$M_{\text{inv}}(e^+e^-) > 140 \text{ MeV}/c^2$





pp data scaled to the same π^0 cross section as in **np** data set.

np excess above **pp** higher than Shyam/Mosel calculations with charged pion FF

angular distributions of proton in the center of mass

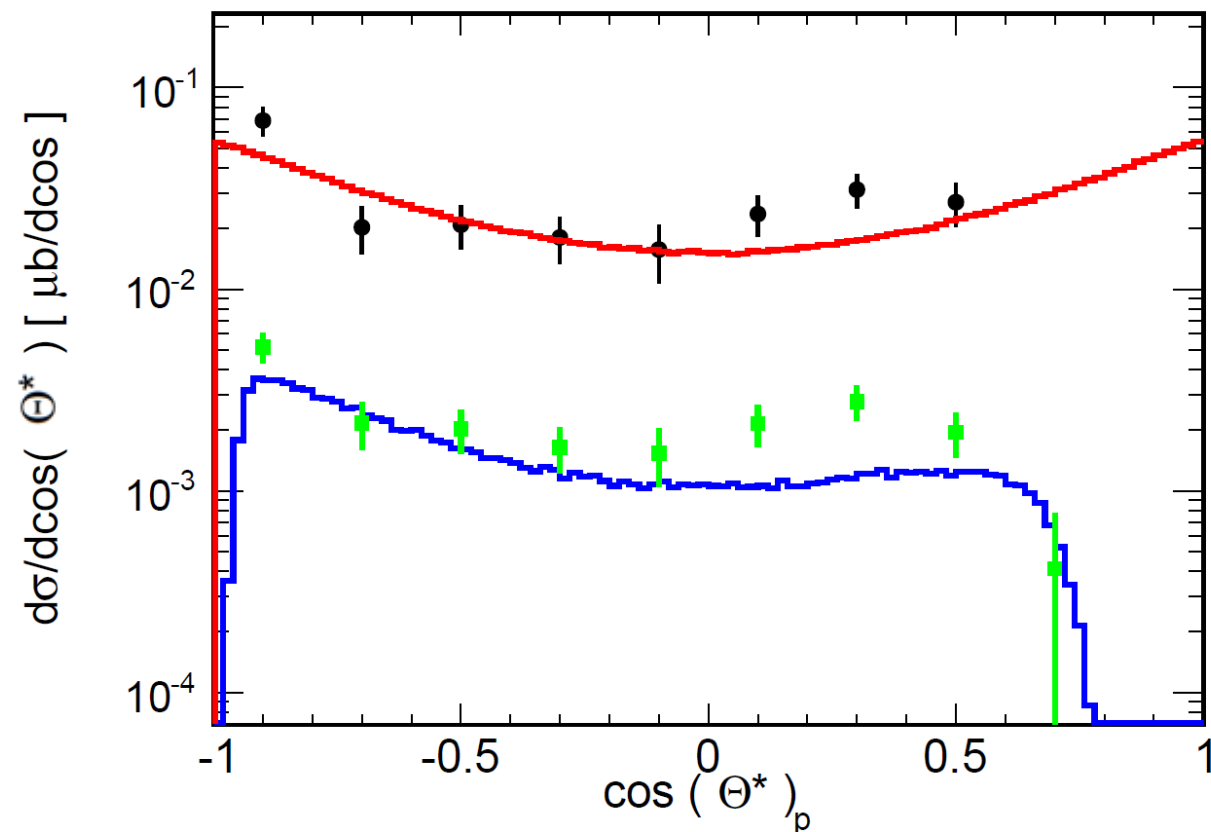
$140 < M_{\text{inv}}(e^+e^-) < 280 \text{ MeV}/c^2$

Data corrected to 4π

Pluto simulation ($\Delta \rightarrow p e^+ e^-$)

Data in acceptance (EFF corrected)

Sim in acceptance (EFF corrected)



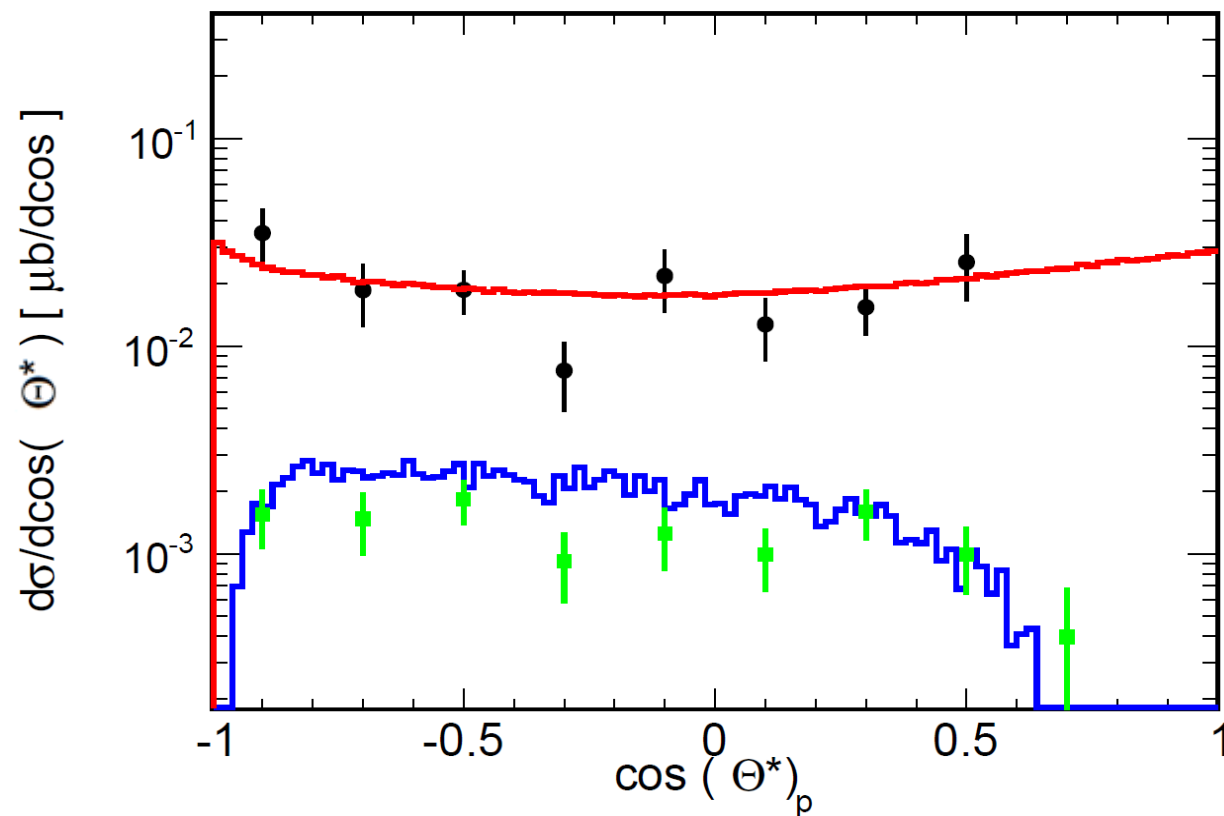
$M_{\text{inv}}(e^+e^-) > 280 \text{ MeV}/c^2$

Data corrected to 4π

Bashkanov & Clement

Data in acceptance (EFF corrected)

Sim in acceptance (EFF corrected)



angular distributions of virtual photon (γ^*) in the center of mass

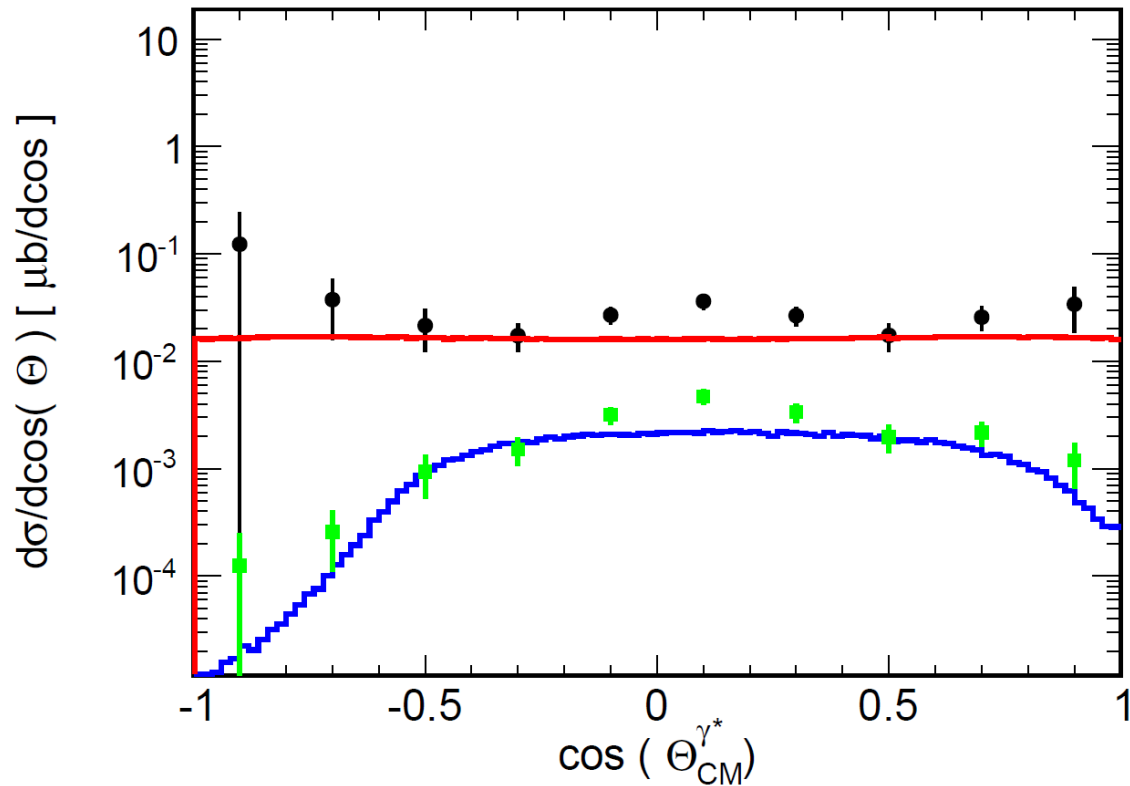
$140 < M_{\text{inv}}(e^+e^-) < 280 \text{ MeV}/c^2$

Data corrected to 4π

Pluto simulation ($\Delta \rightarrow p e^+ e^-$)

Data in acceptance (EFF corrected)

Sim in acceptance (EFF corrected)



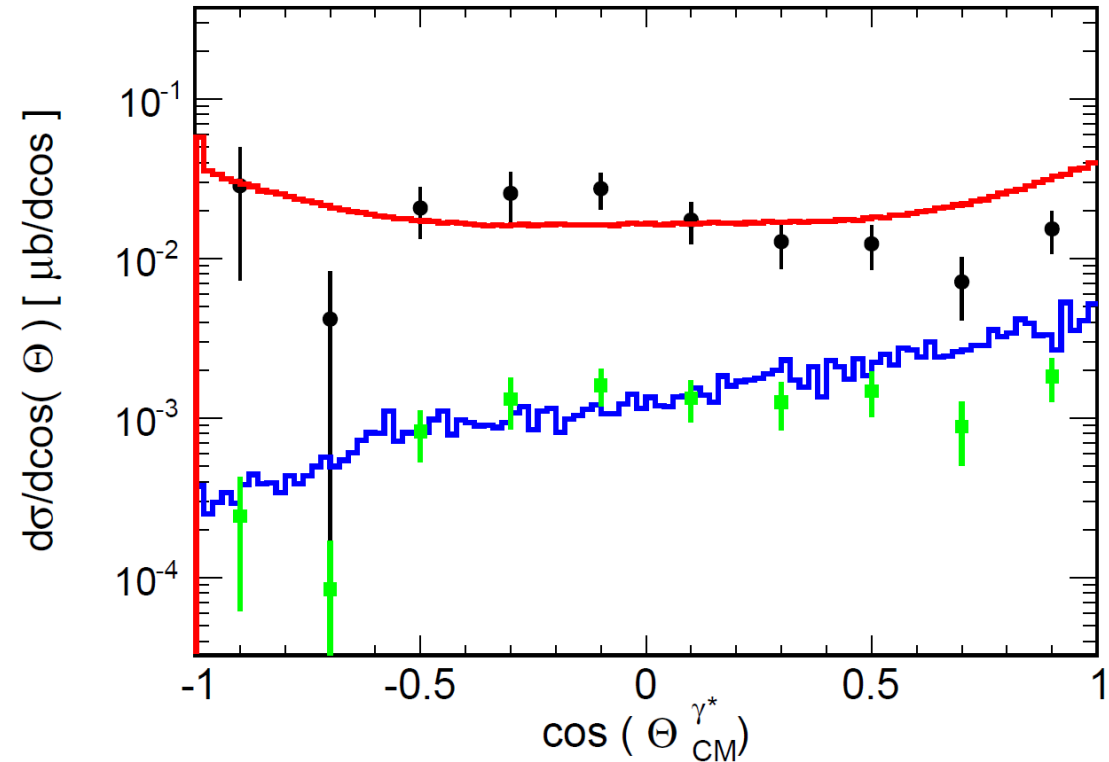
$M_{\text{inv}}(e^+e^-) > 280 \text{ MeV}/c^2$

Data corrected to 4π

Bashkanov & Clement

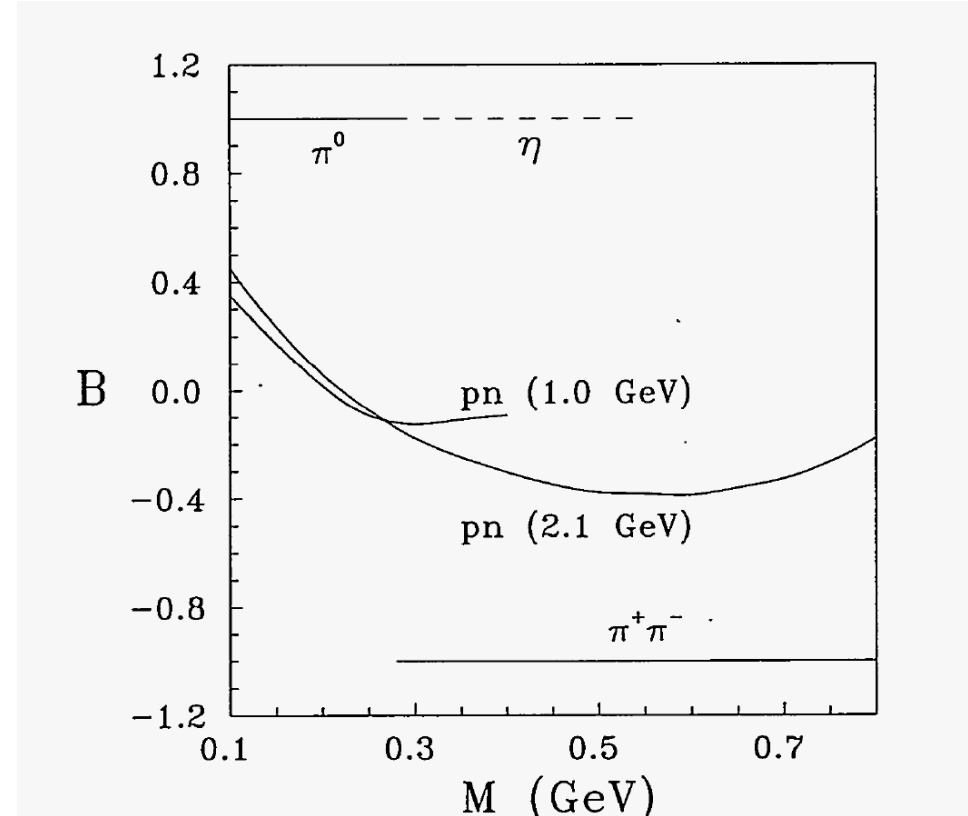
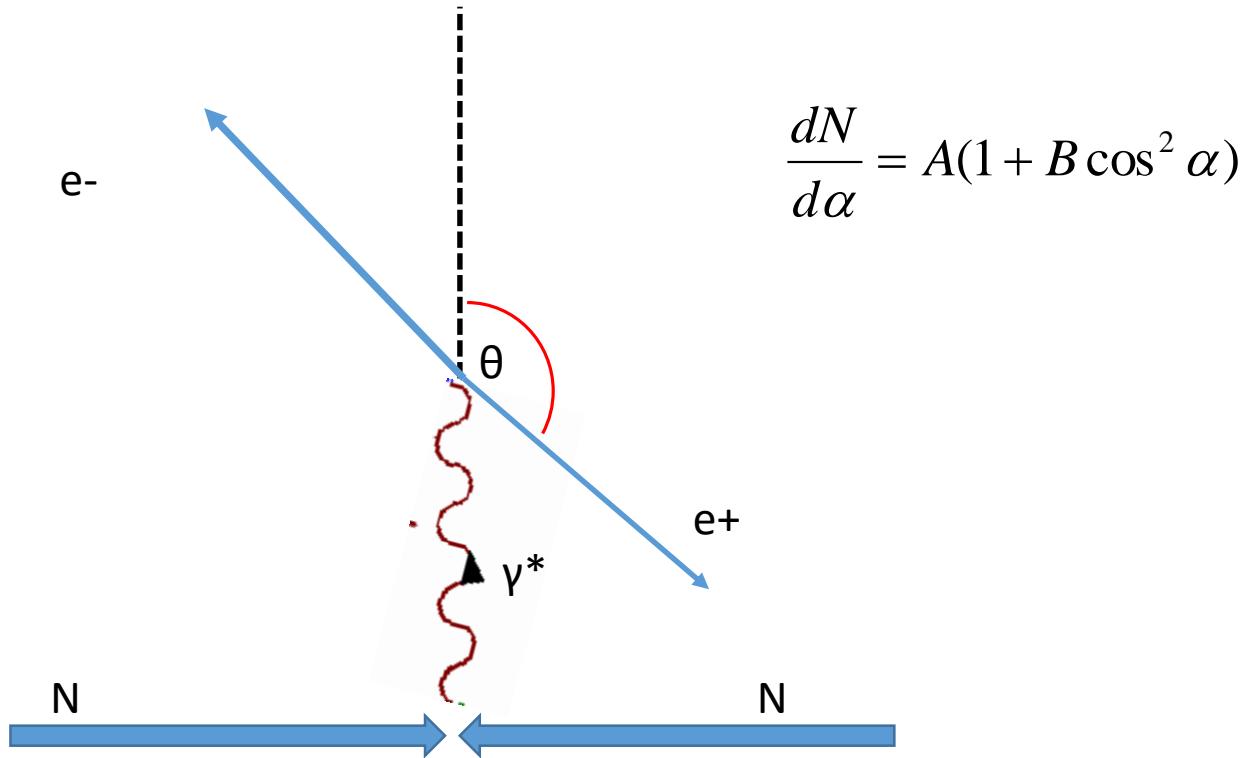
Data in acceptance (EFF corrected)

Sim in acceptance (EFF corrected)



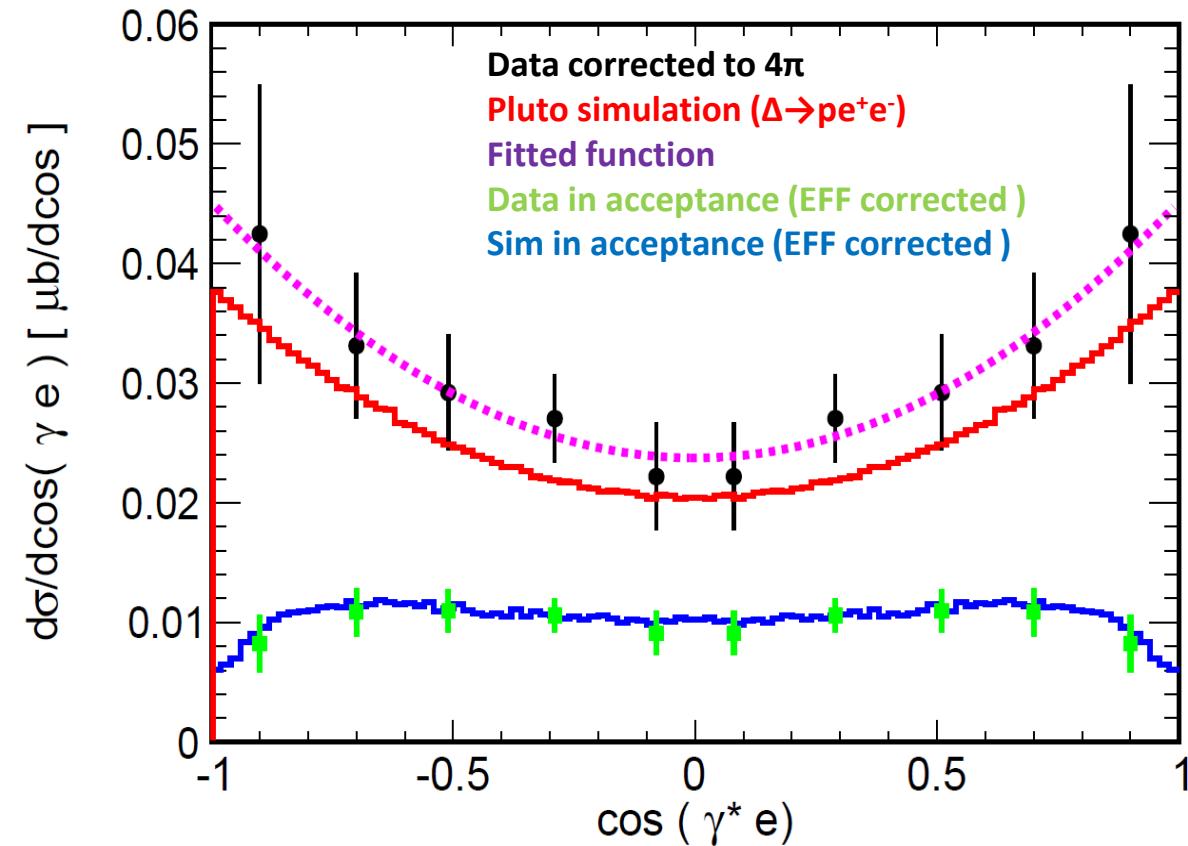
Pseudo- Helicity

- Pseudo- Helicity is defined as the angle between the lepton and the virtual photon in the virtual photon rest frame (leptons are boosted directly to γ^* rest frame)
- Two regions of interest selected
- Data extrapolated to 4π



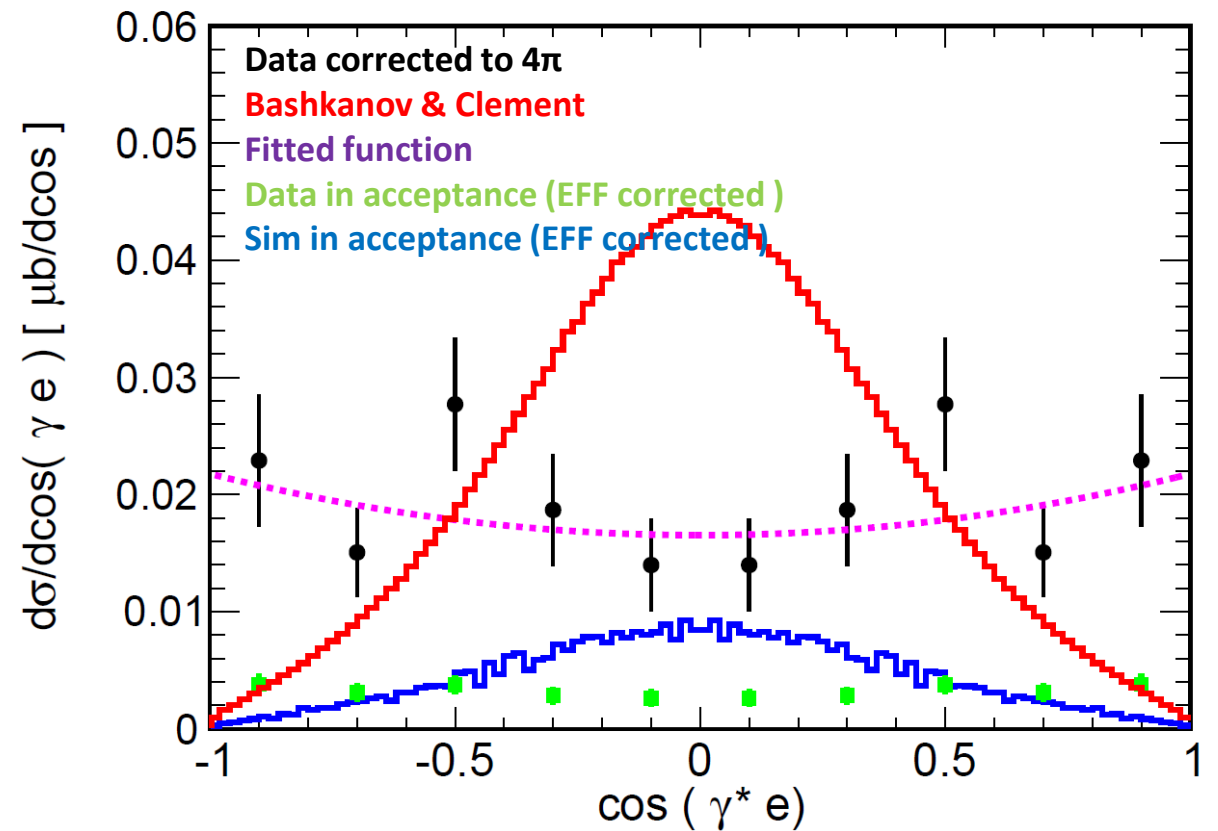
Pseudo-Helicity

$140 < M_{\text{inv}}(e^+e^-) < 280 \text{ MeV}/c^2$



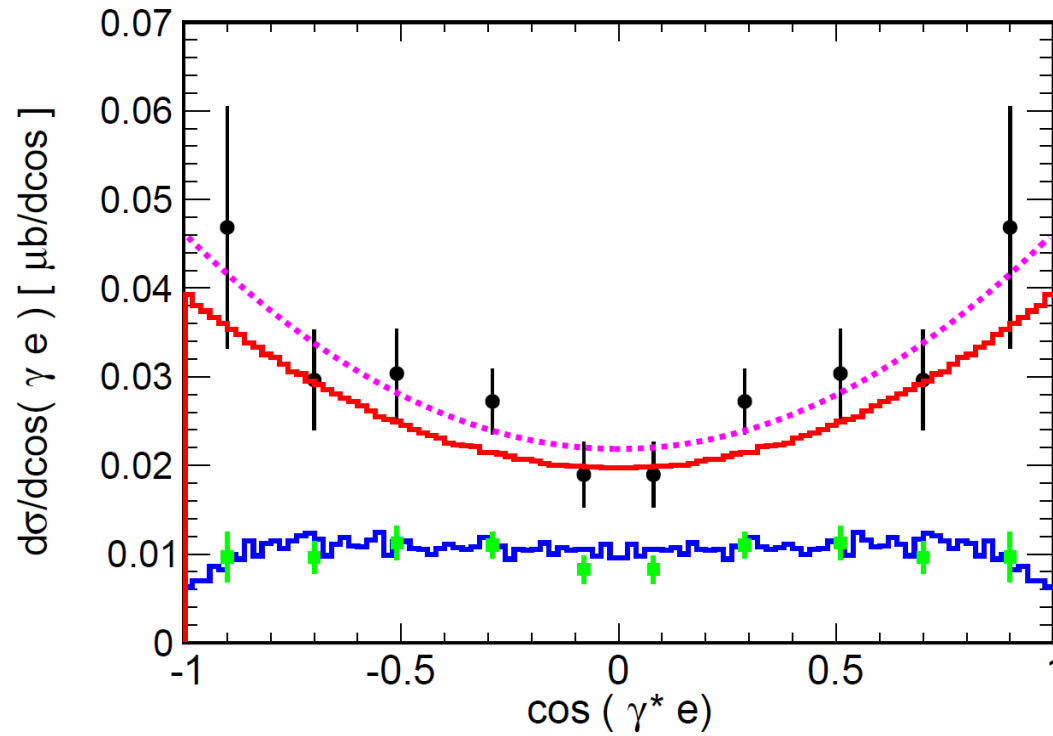
| $140 \text{ MeV}/c^2 < M < 280 \text{ MeV}/c^2$ | Anisotropy parameter (B) |
|-------------------------------------------------|--------------------------|
| Simulation | 0.77 ± 0.006 |
| Experiment | 0.9 ± 0.36 |

$M_{\text{inv}}(e^+e^-) > 280 \text{ MeV}/c^2$

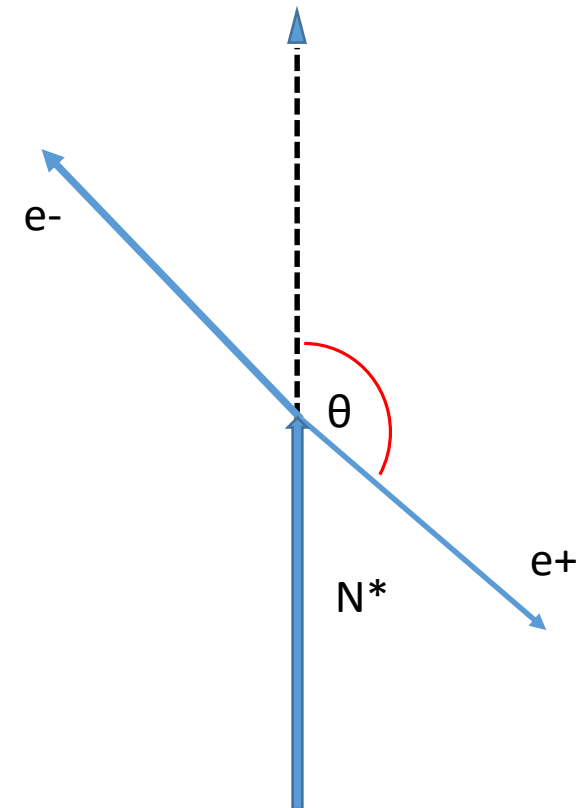


| $280 \text{ MeV}/c^2 < M$ | Anisotropy parameter (B) |
|---------------------------|--------------------------|
| Simulation | -1.30 ± 0.003 |
| Experiment | 0.15 ± 0.32 |

Helicity



- Since there is a confirmation of the major contribution of Δ in $e^+ e^-$ production in the range of $140 \text{ MeV}/c^2 < M < 280 \text{ MeV}/c^2$ Helicity has been calculated (boost to Δ reference frame)



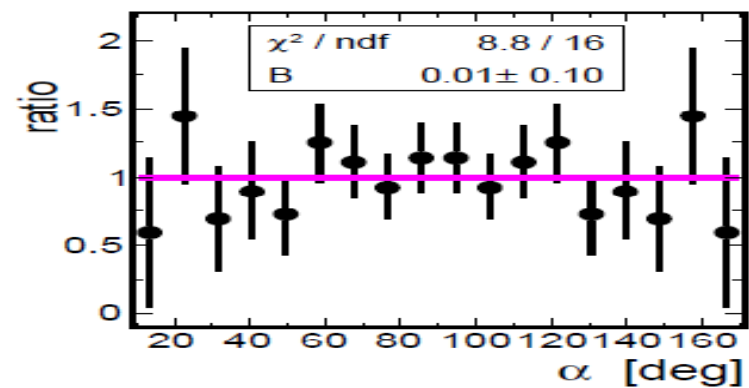
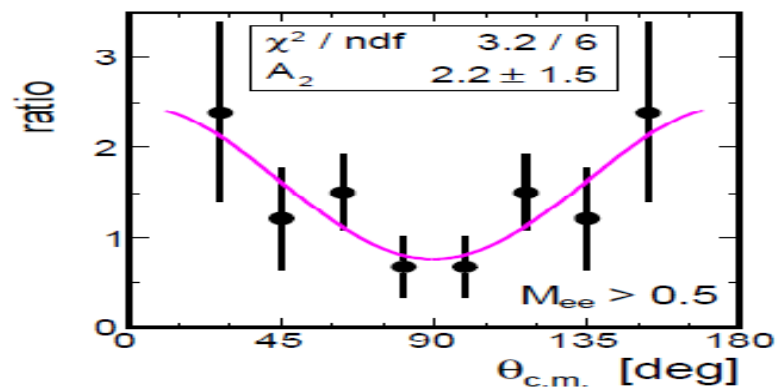
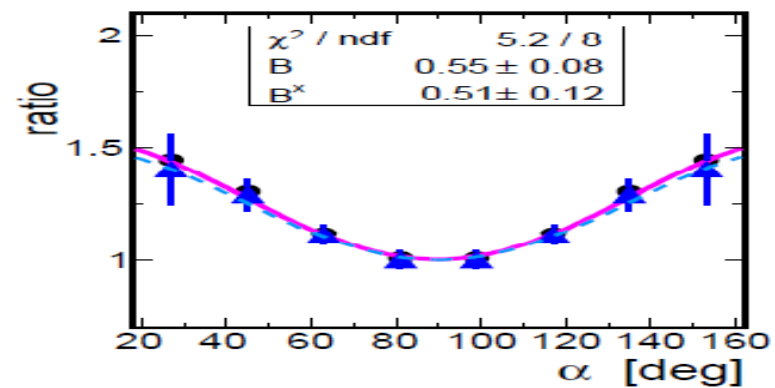
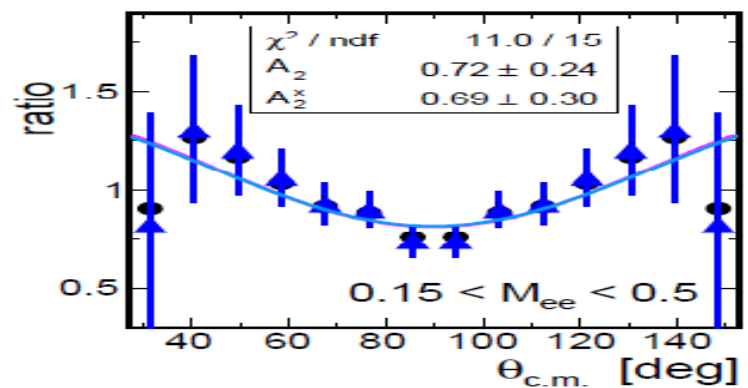
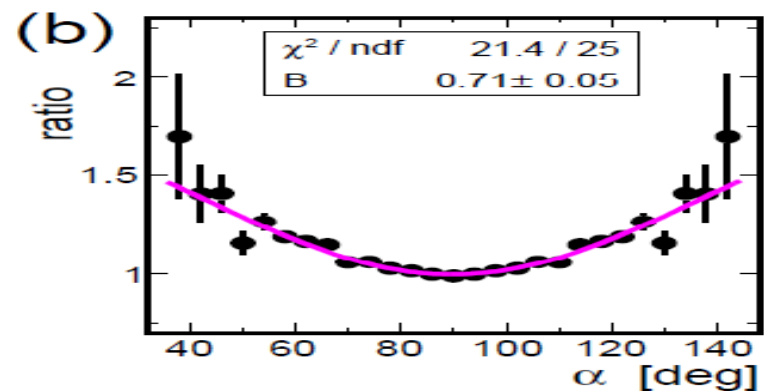
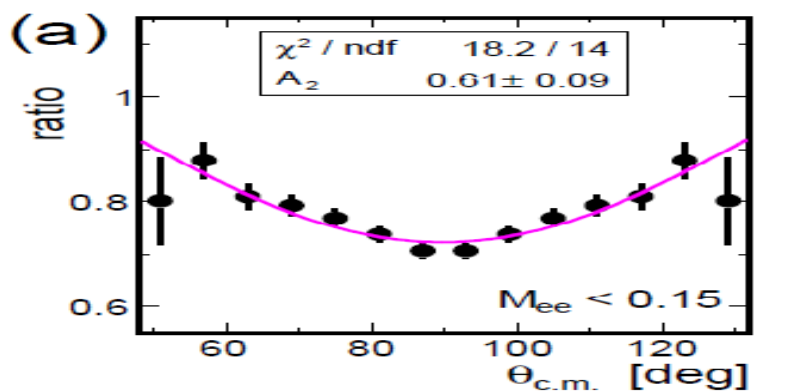
| $140 \text{ MeV}/c^2 < M < 280 \text{ MeV}/c^2$ | Anisotropy parameter (B) |
|-------------------------------------------------|--------------------------|
| Simulation | 1 ± 0.006 |
| Experiment | 1.1 ± 0.4 |

Conclusion

- excess of $e^+ e^-$ pairs in np over pp is a genuine feature of the exclusive channel
- Helicity distributions show an interesting pattern:
 - a) In mass region dominated with Δ , anisotropy is in agreement with expectation
 - b) In higher mass region (ρ - dominated) the distribution is isotropic \rightarrow similarity with Heavy Ion
- Model of Bashkanov overestimates the data by a factor of 2.
- Virtual photon distributions are isotropic
- Proton distributions obtained from the data are mostly described by the model
- charged pion FF in bremsstrahlung alone does not describe the ratio of np/pp

Backup

Results obtained from Ar-KCl run



The HADES spectrometer

- **Detector geometry**

full azimuthal range covered, 6 sectors
polar angle: $16^\circ < \theta < 84^\circ$

- **Tracking**

Superconducting coils, toroidal field
24 Mini Drift Chambers

- **Particle identification (e, p, K, π)**

RICH, MDC, TOF, TOFINO, Shower (RPC)

- **Resolutions**

$\Delta M\omega/M\omega \sim 2.1\%$ at ω peak

$\Delta p/p \sim 2-3\%$ for proton and π

- **Forward Wall:**

Plastic scintillators covering θ angles up to 7°
Detector dedicated to tag proton spectator

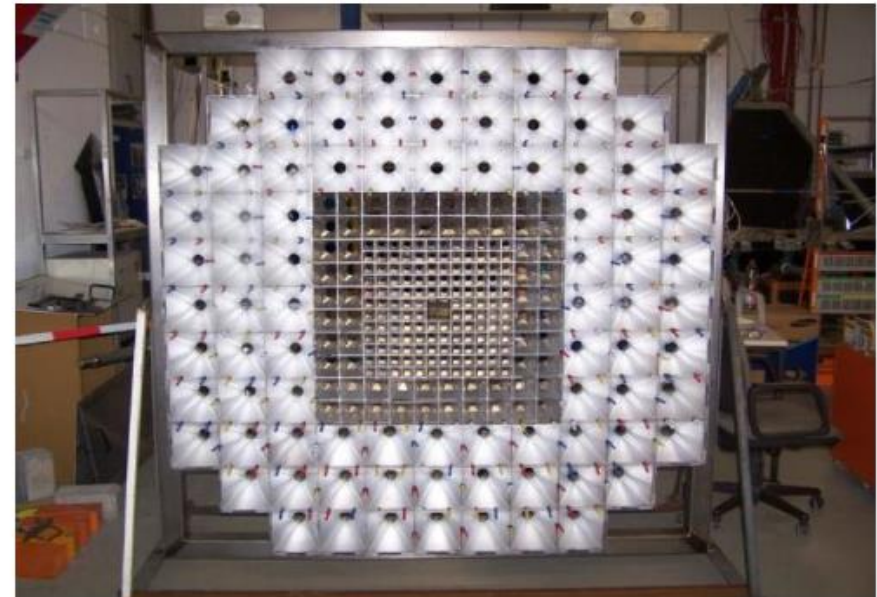
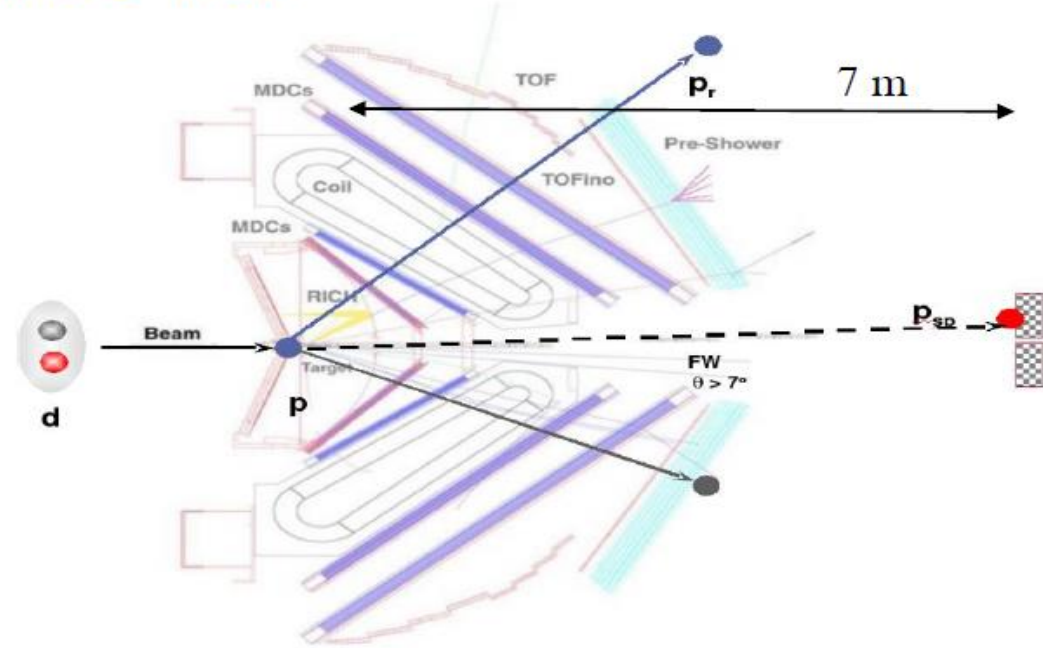
- **Cells in FW:**

140 small 4x4cm $\rightarrow (0^\circ < \theta < 2^\circ)$

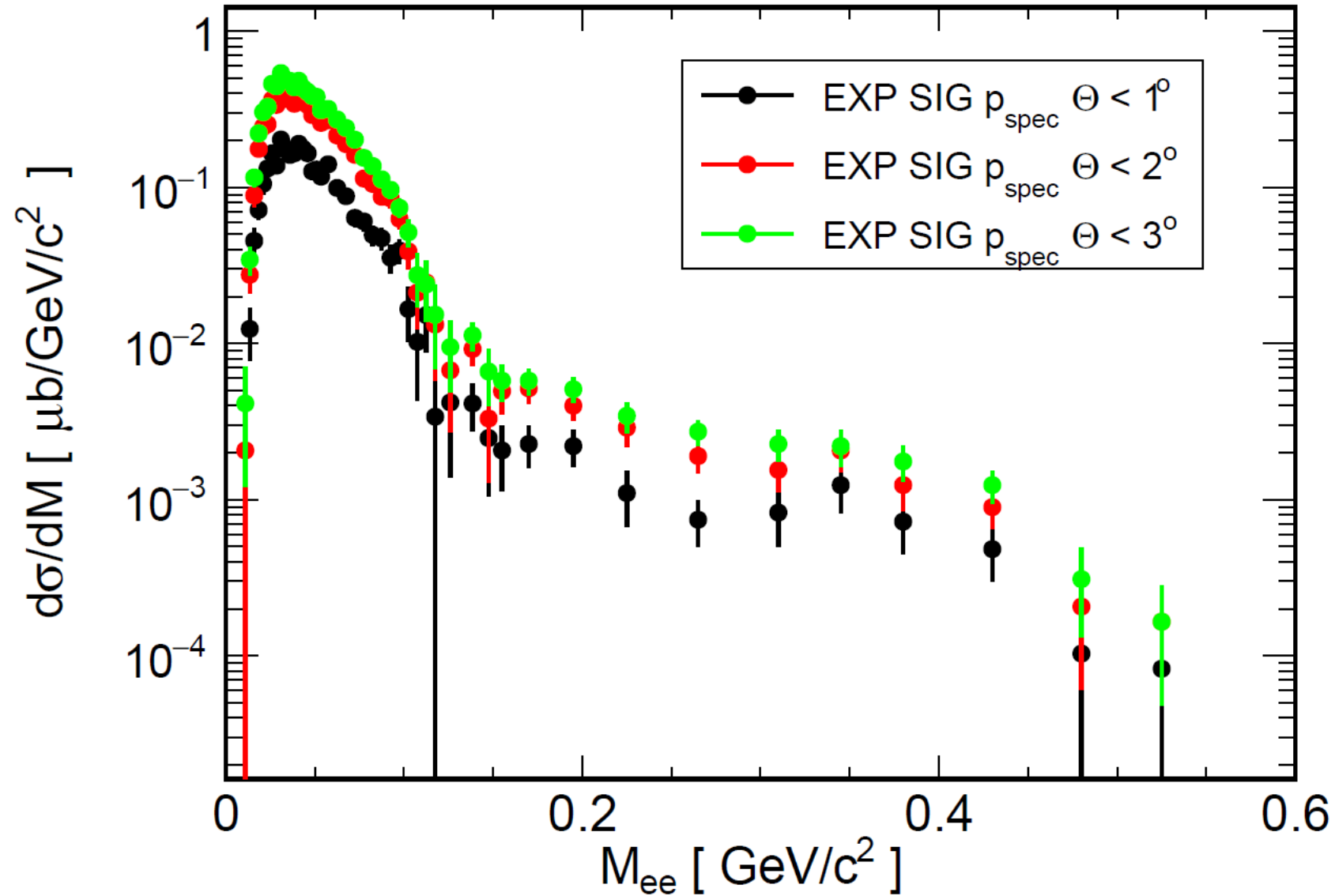
64 middle 8x8cm $\rightarrow (2^\circ < \theta < 3.3^\circ)$

84 large 16x16cm $\rightarrow (3.3^\circ < \theta < 7.2^\circ)$

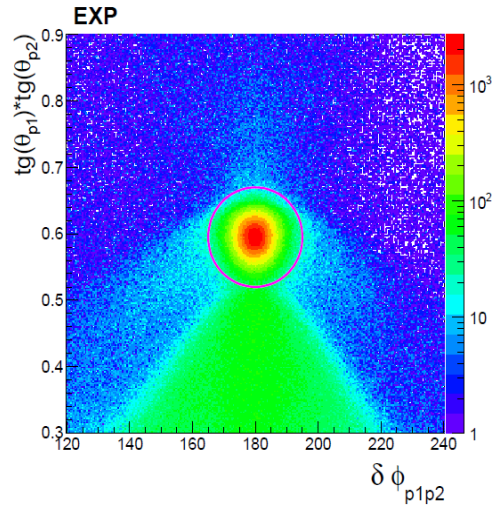
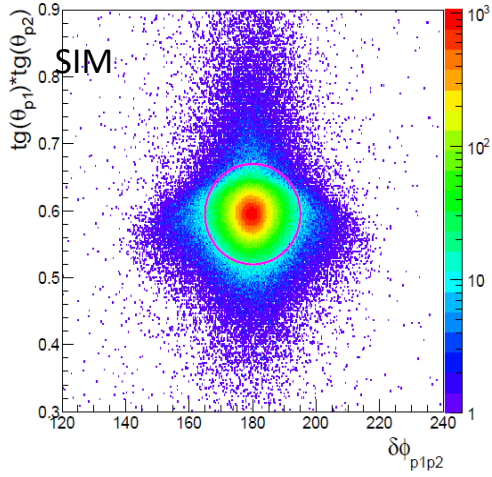
Designed for di-electron spectroscopy,
also suited for the charged hadron detection



Exclusive invariant mass distributions for various p_{spec} angles



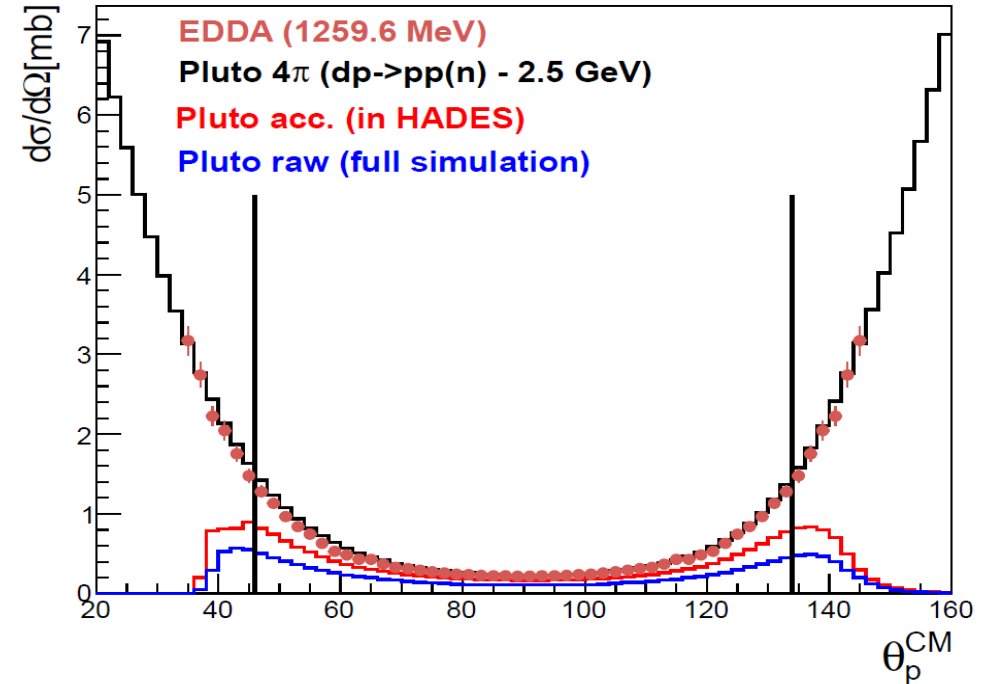
Normalization of HADES data in n-p collisions



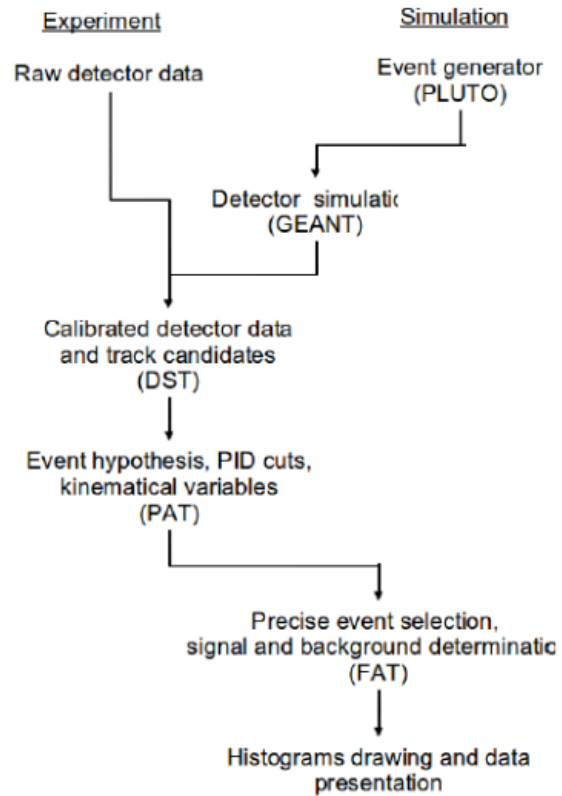
Selection of pp elastic events measured simultaneously by HADES based on angular correlation

- ✓ acceptance and efficiency corrections in the angular range $46^\circ < \Theta_{\text{CM}} < 134^\circ$
- ✓ normalization to the known cross section from the EDDA experiment in the same angular range

$K = \sigma_{\text{el}} / N_{\text{el}} = (2,95 \pm 0,25) \cdot 10^{-9} \text{ mb/counts}$
normalization factor
applied to the measured yield



Selected channels for simulation : resonance model



| lp. | channel | cr. sect. | br. ratio |
|-----|--------------------------------------------------------------------------------------------------------|-------------------|----------------------|
| 1. | $d p \Rightarrow p_{sp} n \Delta^+ \Rightarrow p_{sp} n p \pi^0 \rightarrow p_{sp} n p \gamma e^+ e^-$ | 3.67mb [59] | 0.012 |
| 2. | $d p \Rightarrow p_{sp} p \Delta^0 \Rightarrow p_{sp} p n \pi^0 \rightarrow p_{sp} p n \gamma e^+ e^-$ | 3.67mb [59] | 0.012 |
| 3. | $d p \Rightarrow n_{sp} p \Delta^+ \Rightarrow n_{sp} p p \pi^0 \rightarrow n_{sp} n p \gamma e^+ e^-$ | 3.67mb [59] | 0.012 |
| 4. | $d p \Rightarrow p_{sp} p \Delta^0 \rightarrow p_{sp} p n e^+ e^-$ | 5.54mb [59] | $4.82 \cdot 10^{-5}$ |
| 5. | $d p \Rightarrow p_{sp} n \Delta^+ \rightarrow p_{sp} n p e^+ e^-$ | 5.54mb [59] | $4.93 \cdot 10^{-5}$ |
| 6. | $d p \Rightarrow n_{sp} p \Delta^+ \rightarrow n_{sp} p p e^+ e^-$ | 5.54mb [59] | $4.94 \cdot 10^{-5}$ |
| 7. | $d p \Rightarrow p_{sp} n p \eta \rightarrow p_{sp} n p \gamma e^+ e^-$ | 13.6 μ b [38] | $5.86 \cdot 10^{-3}$ |
| 8. | $d p \Rightarrow p_{sp} d \eta \rightarrow p_{sp} d \gamma e^+ e^-$ | 23.9 μ b [38] | $5.82 \cdot 10^{-3}$ |
| 9. | $d p \Rightarrow n_{sp} p p \eta \rightarrow n_{sp} p p \gamma e^+ e^-$ | 2.33 μ b [38] | $5.84 \cdot 10^{-3}$ |
| 10. | $d p \Rightarrow p_{sp} n p e^+ e^-$ | 1.48 μ b [17] | 1 |
| 11. | $d p \Rightarrow p_{sp} p N^0(1520) \Rightarrow p_{sp} p n \rho \rightarrow p_{sp} p n e^+ e^-$ | 8.91 μ b [63] | $8.12 \cdot 10^{-4}$ |
| 12. | $d p \Rightarrow n_{sp} p N^+(1520) \Rightarrow n_{sp} p p \rho \rightarrow n_{sp} p p e^+ e^-$ | 8.91 μ b [63] | $8.12 \cdot 10^{-4}$ |
| 13. | $d p \Rightarrow p_{sp} d \rho \rightarrow p_{sp} d e^+ e^-$ | 6.40 μ b [63] | $8.12 \cdot 10^{-4}$ |
| 14. | $d p \Rightarrow p_{sp} d \gamma^* \Rightarrow p_{sp} d e^+ e^-$ | 41.7nb [60] | 1 |