

Looking for chiral anomaly in $K\gamma \rightarrow K\pi$ reactions

Phys. Rev. **D93**, 094029 (2016); 1512.04438

M. I. Vysotsky, E. V. Zhemchugov

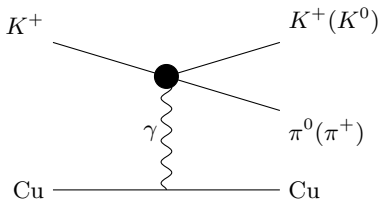
A. I. Alikhanov Institute for Theoretical and Experimental Physics
Moscow, Russia

14th International Workshop on Meson Production,
Properties and Interaction

June 2–7, 2016
Kraków, Poland

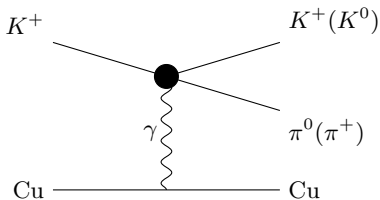
Institute for High-Energy Physics
Protvino, Russia
OKA Detector

Current experiment
 $E_K = 17.7$ GeV.

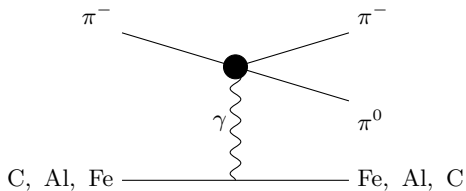


Institute for High-Energy Physics
Protvino, Russia
OKA Detector

Current experiment
 $E_K = 17.7$ GeV.



[Yu. M. Antipov *et. al.*, Phys. Rev. **D36**, 21 (1987)]
 $E_\pi = 40$ GeV.



$$A(\pi^- \gamma \rightarrow \pi^- \pi^0) = h(s, t, u) \cdot \varepsilon^{\mu\alpha\beta\gamma} A_\mu \partial_\alpha \pi^- \partial_\beta \pi^+ \partial_\gamma \pi^0$$

Sutherland-Veltman (chiral symmetry): $h(0, 0, 0) \equiv h(0) = 0$

Relation to the $\pi^0 \rightarrow \gamma\gamma$ process¹

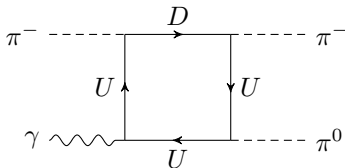
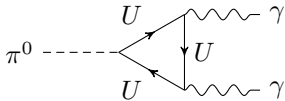
Wess-Zumino anomaly²

Direct calculation of the box

$$h(0) = \frac{e}{4\pi^2 F_\pi^3} = 9.8 \text{ GeV}^{-3}$$

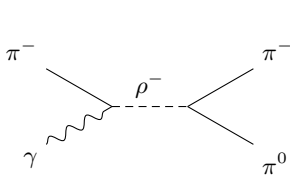
($F_\pi = 92.2 \text{ MeV}$ from $\pi \rightarrow \ell\nu$ decay)

$h(0) \neq 0 \Rightarrow$ chiral anomaly

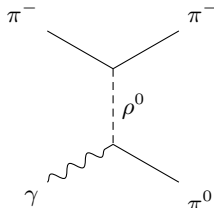


¹[Terent'ev, JETP Letters **14**, 94 (1971)]

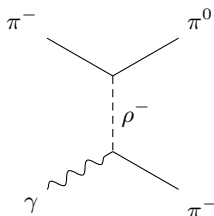
²[Wess, Zumino, Phys. Lett. **37B**, 95 (1971)]



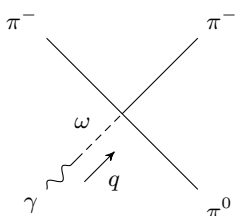
(s)



(t)



(u)



$$A(\pi^- \gamma \rightarrow \pi^- \pi^0) = h(s, t, u) \cdot \varepsilon^{\mu\alpha\beta\gamma} A_\mu \partial_\alpha \pi^- \partial_\beta \pi^+ \partial_\gamma \pi^0$$

$$h(s, t, u) = h(0) \left\{ 1 + \frac{2f_{\rho\pi\pi} f_{\rho\pi\gamma}}{m_\rho^2 h(0)} \left[\frac{s}{m_\rho^2 - s} + \frac{t}{m_\rho^2 - t} + \frac{u}{m_\rho^2 - u} \right] + \frac{f_{\omega\gamma} f_{\omega 3\pi}}{m_\omega^2 h(0)} \frac{q^2}{m_\omega^2 - q^2} \right\}$$

[Terent'ev, Phys. Lett. **38B**, 419 (1972)]

$$A(\pi^- \gamma \rightarrow \pi^- \pi^0) = h(s, t, u) \cdot \varepsilon^{\mu\alpha\beta\gamma} A_\mu \partial_\alpha \pi^- \partial_\beta \pi^+ \partial_\gamma \pi^0$$

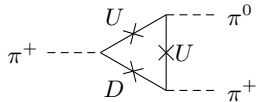
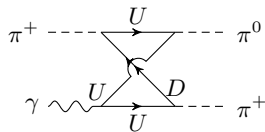
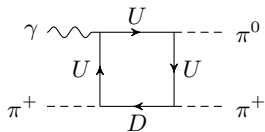
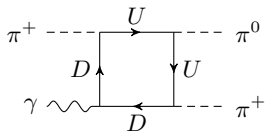
$$h(s, t, u) = h(0) \left\{ 1 + \frac{2f_{\rho\pi\pi} f_{\rho\pi\gamma}}{m_\rho^2 h(0)} \left[\frac{s}{m_\rho^2 - s} + \frac{t}{m_\rho^2 - t} + \frac{u}{m_\rho^2 - u} \right] + \frac{f_{\omega\gamma} f_{\omega 3\pi}}{m_\omega^2 h(0)} \frac{q^2}{m_\omega^2 - q^2} \right\}$$

$h(0)$ values

Theory	9.8 GeV^{-3}
Experiment at LO (1987)	$12.9 \pm 0.9 \pm 0.5 \pm 1.0 \text{ GeV}^{-3}$
Experiment at NNLO + EMC (2001)	$10.7 \pm 1.2 \text{ GeV}^{-3}$

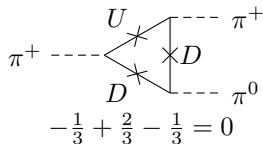
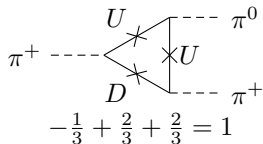
Update from the COMPASS Collaboration?

$$\pi^+ \gamma \rightarrow \pi^+ \pi^0$$

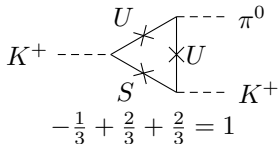


$$-\frac{1}{3} + \frac{2}{3} + \frac{2}{3} = 1$$

$$\pi^+ \gamma \rightarrow \pi^+ \pi^0$$

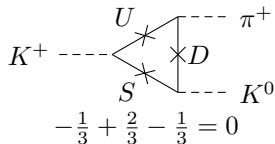


$$K^+ \gamma \rightarrow K^+ \pi^0$$



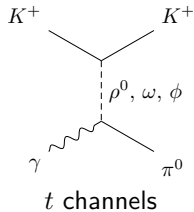
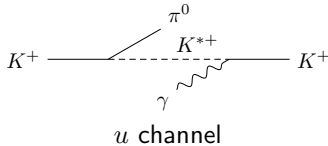
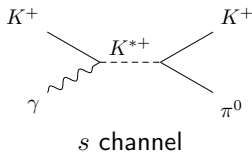
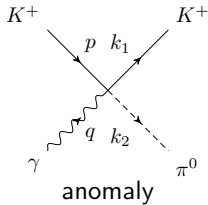
neutral pion production

$$K^+ \gamma \rightarrow K^0 \pi^+$$

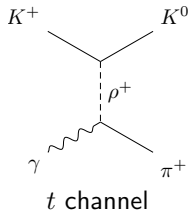
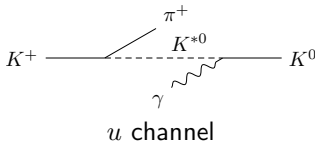
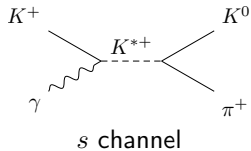


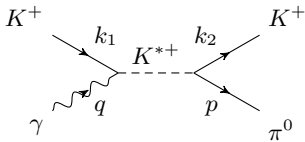
charged pion production

$$K^+ \gamma \rightarrow K^+ \pi^0$$



$$K^+ \gamma \rightarrow K^0 \pi^+$$





s-channel amplitude:

$$A_s^{(0)}(K^+\gamma \rightarrow K^+\pi^0) = -\frac{2f_{K^{*+}K^+\gamma}f_{K^{*+}K^+\pi^0}}{s - m_{K^{*+}}^2 + i\sqrt{s}\Gamma_{K^{*+}}(s)} \varepsilon^{\alpha\beta\gamma\delta} \epsilon_\alpha p_\beta k_{1\gamma} k_{2\delta}$$

$$A_s(K^+\gamma \rightarrow K^+\pi^0) = A_s^{(0)}(K^+\gamma \rightarrow K^+\pi^0) - A_s^{(0)}(K^+\gamma \rightarrow K^+\pi^0)|_{s=0}$$

Cross section:

$$\begin{aligned} \frac{d\sigma(K^+\gamma \rightarrow K^+\pi^0)}{dt} &= \frac{1}{27\pi} \left(t + \frac{(st - m_{K^+}^2 m_{\pi^0}^2)(t - m_{\pi^0}^2)}{(s - m_{K^+}^2)^2} \right) \\ &\times \left| \frac{e}{4\pi^2 F_\pi^3} + \frac{2f_{K^{*+}K^+\gamma}f_{K^{*+}K^+\pi^0}}{m_{K^{*+}}^2 - s - i\sqrt{s}\Gamma_{K^{*+}}(s)} \cdot \frac{s}{m_{K^{*+}}^2} \right. \\ &+ \frac{2f_{K^{*+}K^+\gamma}f_{K^{*+}K^+\pi^0}}{m_{K^{*+}}^2 - u} \cdot \frac{u}{m_{K^{*+}}^2} + \frac{2f_{\rho^0\pi^0\gamma}f_{\rho^0K^+K^+}}{m_{\rho^0}^2 - t} \cdot \frac{t}{m_{\rho^0}^2} \\ &\left. + \frac{2f_{\omega\pi^0\gamma}f_{\omega K^+K^+}}{m_\omega^2 - t} \cdot \frac{t}{m_\omega^2} + \frac{2f_{\phi\pi^0\gamma}f_{\phi K^+K^+}}{m_\phi^2 - t} \cdot \frac{t}{m_\phi^2} \right|^2 \end{aligned}$$

$$f_{K^{*+}K^+\pi^0} = 3.10$$

$$f_{K^{*+}K^0\pi^+} = 4.38$$

$$f_{K^{*0}K^+\pi^+} = 4.41$$

$$f_{\rho^0 K^+ K^+} = 3.16$$

$$f_{\rho^+ K^+ K^0} = -4.47$$

$$f_{\omega K^+ K^+} = 3.16$$

$$f_{\phi K^+ K^+} = -4.47$$

$$f_{K^{*+}K^+\gamma} = 0.240 \text{ GeV}^{-1}$$

$$f_{K^{*0}K^0\gamma} = -0.385 \text{ GeV}^{-1}$$

$$f_{\rho^0\pi^0\gamma} = 0.252 \text{ GeV}^{-1}$$

$$f_{\rho^+\pi^+\gamma} = 0.219 \text{ GeV}^{-1}$$

$$f_{\omega\pi^0\gamma} = 0.696 \text{ GeV}^{-1}$$

$$f_{\phi\pi^0\gamma} = 0.040 \text{ GeV}^{-1}$$

Decay widths:

$$\Gamma(K^* \rightarrow K\pi) \implies |f_{K^*K\pi}|$$

$$\Gamma(K^* \rightarrow K\gamma) \implies |f_{K^*K\gamma}|$$

$$\Gamma(\phi \rightarrow K^+K^-) \implies |f_{\phi K^+K^+}|$$

$$\Gamma(\rho^+ \rightarrow \pi^+\gamma) \implies |f_{\rho^+\pi^+\gamma}|$$

$$\Gamma(\rho^0 \rightarrow \pi^0\gamma) \implies |f_{\rho^0\pi^0\gamma}|$$

$$\Gamma(\omega \rightarrow \pi^0\gamma) \implies |f_{\omega\pi^0\gamma}|$$

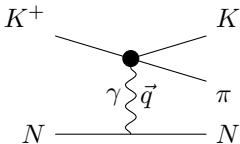
$$\Gamma(\phi \rightarrow \pi^0\gamma) \implies |f_{\phi\pi^0\gamma}|$$

$SU(3)$ symmetry:

$$\begin{aligned}\sqrt{2}f_{K^{*+}K^+\pi^0} &= f_{K^{*+}K^0\pi^+} = f_{K^{*0}K^+\pi^+} = -f_{\rho^+K^+K^0} \\ &= \sqrt{2}f_{\rho^0 K^+ K^+} = \sqrt{2}f_{\omega K^+ K^+} = -f_{\phi K^+ K^+}\end{aligned}$$

$$f_{K^{*+}K^+\gamma} = f_{\rho^+\pi^+\gamma} = f_{\rho^0\pi^0\gamma} = \frac{1}{3}f_{\omega\pi^0\gamma} = -\frac{1}{2}f_{K^{*0}K^0\gamma}$$

The sign of the anomaly term is unknown.



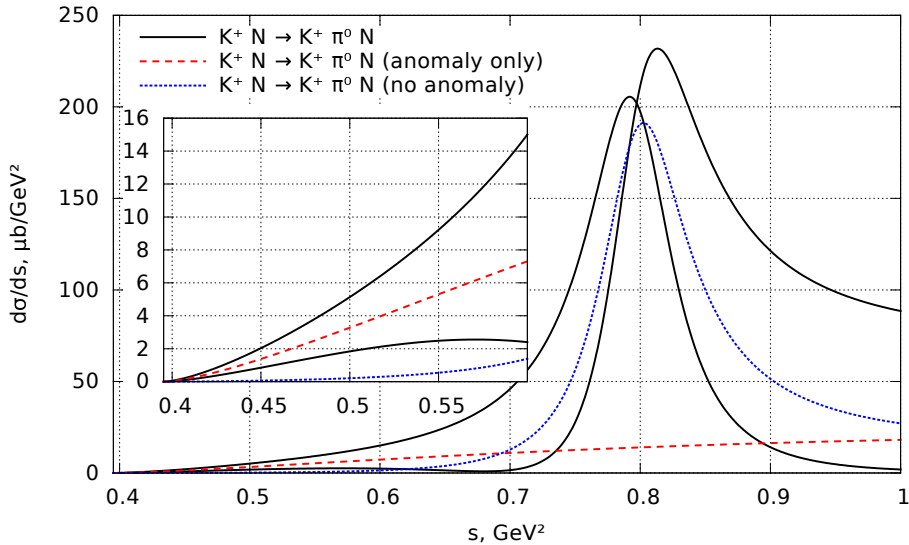
Weizsacker-Williams equivalent photons approximation:

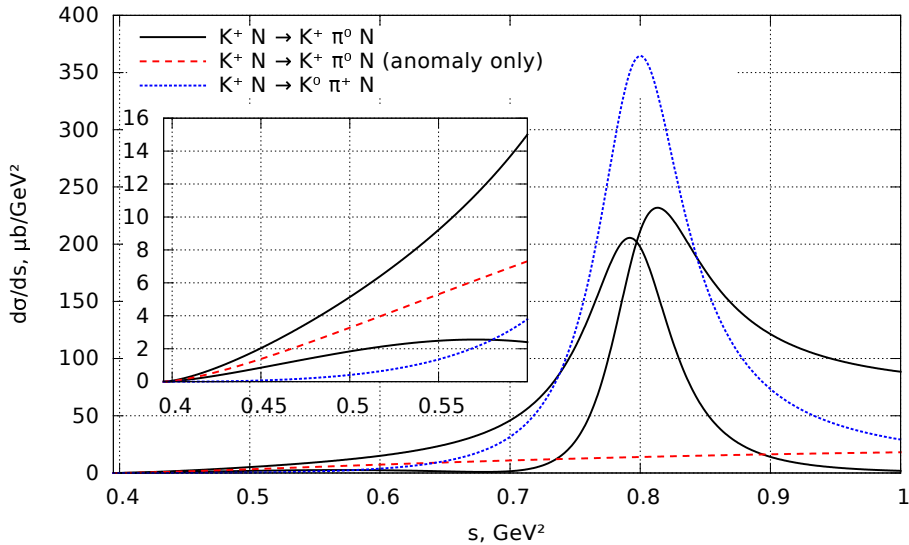
$$\frac{d\sigma(K^+N \rightarrow K\pi N)}{dt ds dq_{\perp}^2} = \frac{Z^2\alpha}{\pi(s - m_{K^+}^2)} \frac{q_{\perp}^2}{\left(q_{\perp}^2 + \left(\frac{s - m_{K^+}^2}{2E_K}\right)^2\right)^2} \frac{d\sigma(K^+\gamma \rightarrow K\pi)}{dt} |F(\vec{q}^2)|^2$$

$$F(\vec{q}^2) = \exp\left(-\frac{\langle r^2 \rangle \vec{q}^2}{6}\right)$$

$$\frac{d\sigma(K^+N \rightarrow K\pi N)}{dt ds} = \frac{Z^2\alpha}{\pi} \frac{E_1(a) - 1}{s - m_{K^+}^2} \frac{d\sigma(K^+\gamma \rightarrow K\pi)}{dt}$$

$$E_1(a) = \int_a^{\infty} \frac{e^{-z}}{z} dz, \quad a = \frac{1}{3} r_0^2 A^{2/3} \left(\frac{s - m_{K^+}^2}{2E_K}\right)^2$$

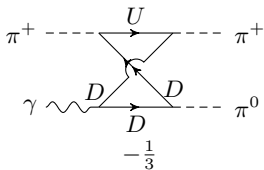
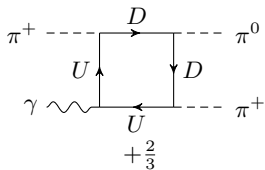
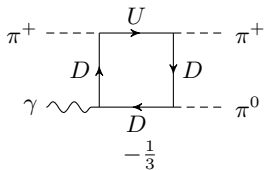
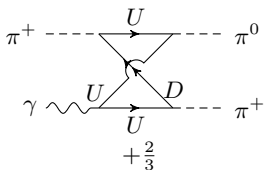
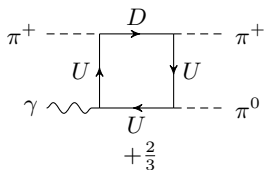
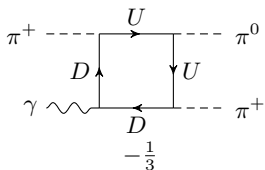




Conclusions

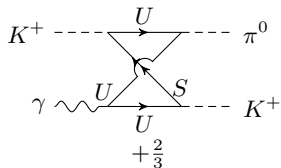
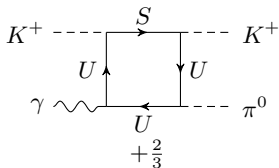
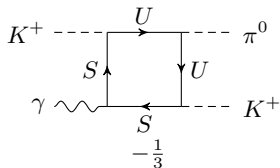
- ▶ A theoretical prediction has been made for the cross sections of $K^+\gamma \rightarrow K^+\pi^0$ and $K^+\gamma \rightarrow K^0\pi^+$ reactions at low energies. For the anomalous reaction, we predict two possible values depending on the a priori unknown sign of the interference term, which should be resolved by the experiment.
- ▶ It is possible to observe the chiral anomaly through comparison of cross section of $K^+ Cu \rightarrow K^+\pi^0 Cu$ reaction with that of $K^+ Cu \rightarrow K^0\pi^+ Cu$ reaction at $\sqrt{s} \lesssim 0.6 \text{ GeV}^2$. The point is that only the first one has the anomaly which manifests itself as an increase in the cross section at low \sqrt{s} .
- ▶ Luminosity of $60\mu\text{b}^{-1}$ at $0.4 < s < 0.6 \text{ GeV}^2$ is planned to be collected in the Protvino experiment. In this case expected observations are ≈ 10 events of $K^0\pi^+$ production and either ≈ 20 or ≈ 70 events of $K^+\pi^0$ production, depending on the sign of the interference term.

$$\pi^+ \gamma \rightarrow \pi^+ \pi^0$$



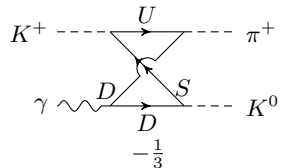
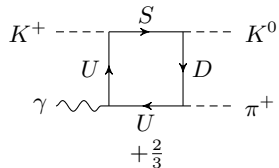
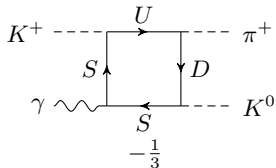
Total: 1

$$K^+ \gamma \rightarrow K^+ \pi^0$$



Total: 1

$$K^+ \gamma \rightarrow K^0 \pi^+$$



Total: 0