

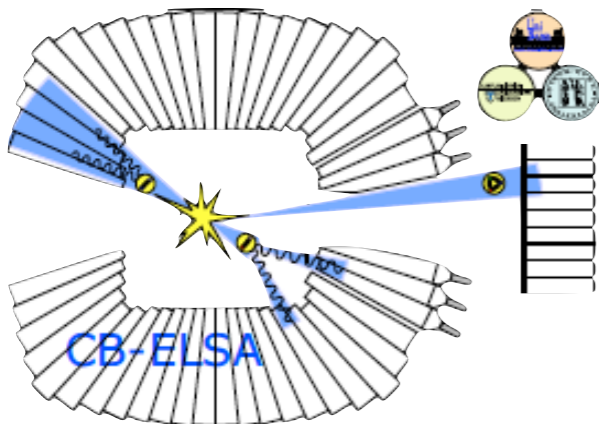
Determination of the ω - and η' -nucleus optical potential

Mariana Nanova

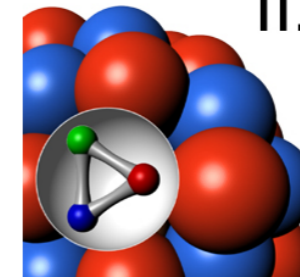
for the CBELSA/TAPS Collaboration

Outline:

- ◆ motivation
- ◆ exp. approaches to study the in-medium properties of mesons
- ◆ experimental results on the real and imaginary part of the ω - and η' -nucleus optical potential
- ◆ summary & outlook

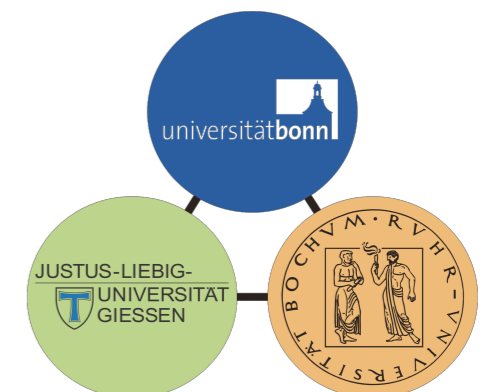


MESON2016
Cracow, 2th - 7th June 2016



II. Physikalisches
Institut

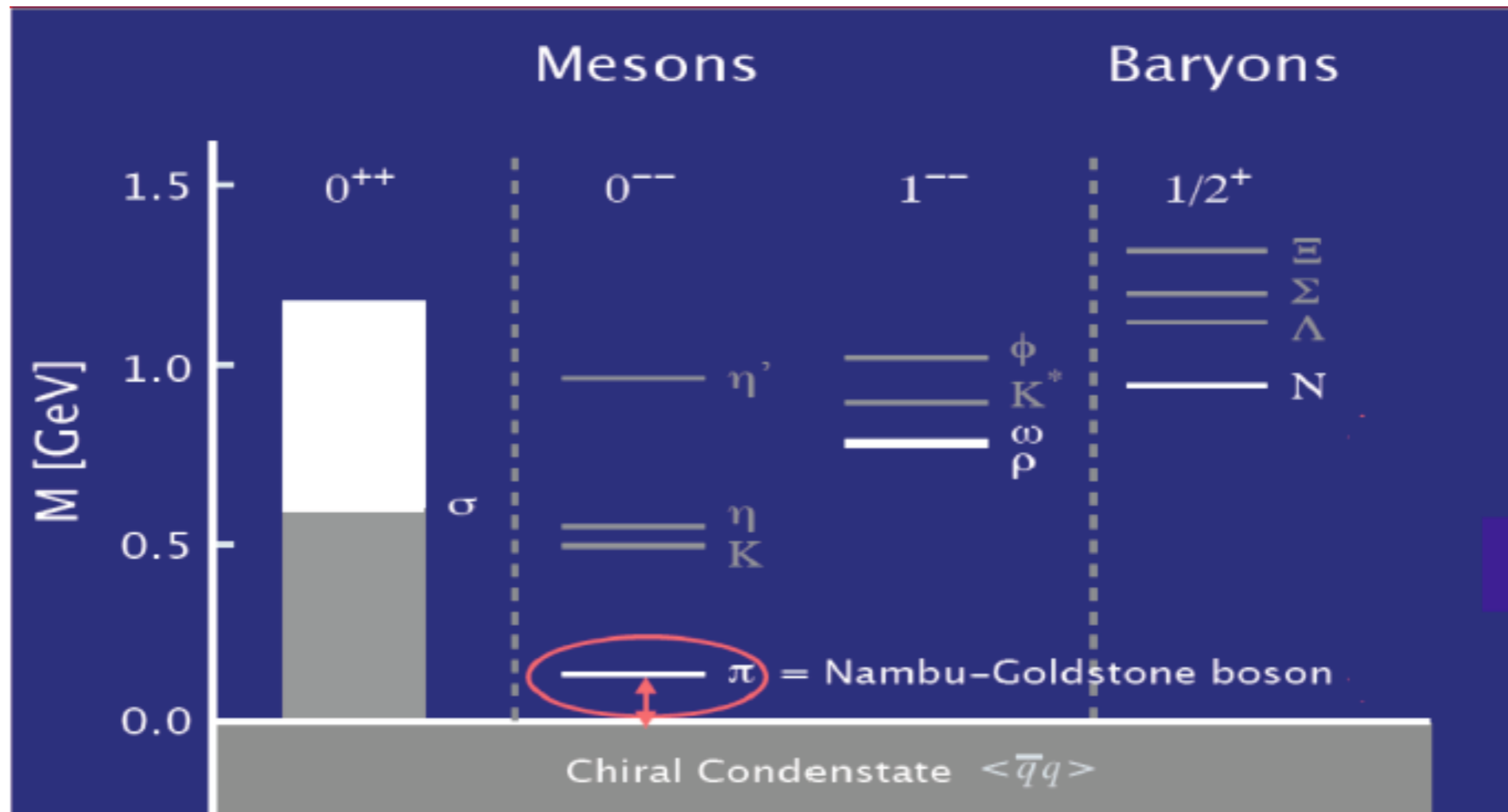
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GIESSEN



*funded by the DFG within SFB/TR16

baryons and mesons

- ◆ QCD vacuum as a Bose-Einstein condensate of $\bar{q}q$
- ◆ all states (particles) are created out of the vacuum state (“excitations of the QCD-vacuum”)
- ◆ the ground-state structure influences the particle properties



if the QCD ground state changes in a medium

\Rightarrow properties of hadrons (“excited states”) are also expected to change

hadrons in the medium

how do the hadron properties (mass, width) change in a dense nuclear medium ??

pioneering papers:

V. Bernard and U.-G. Meißner,
NPA 489 (1988) 647

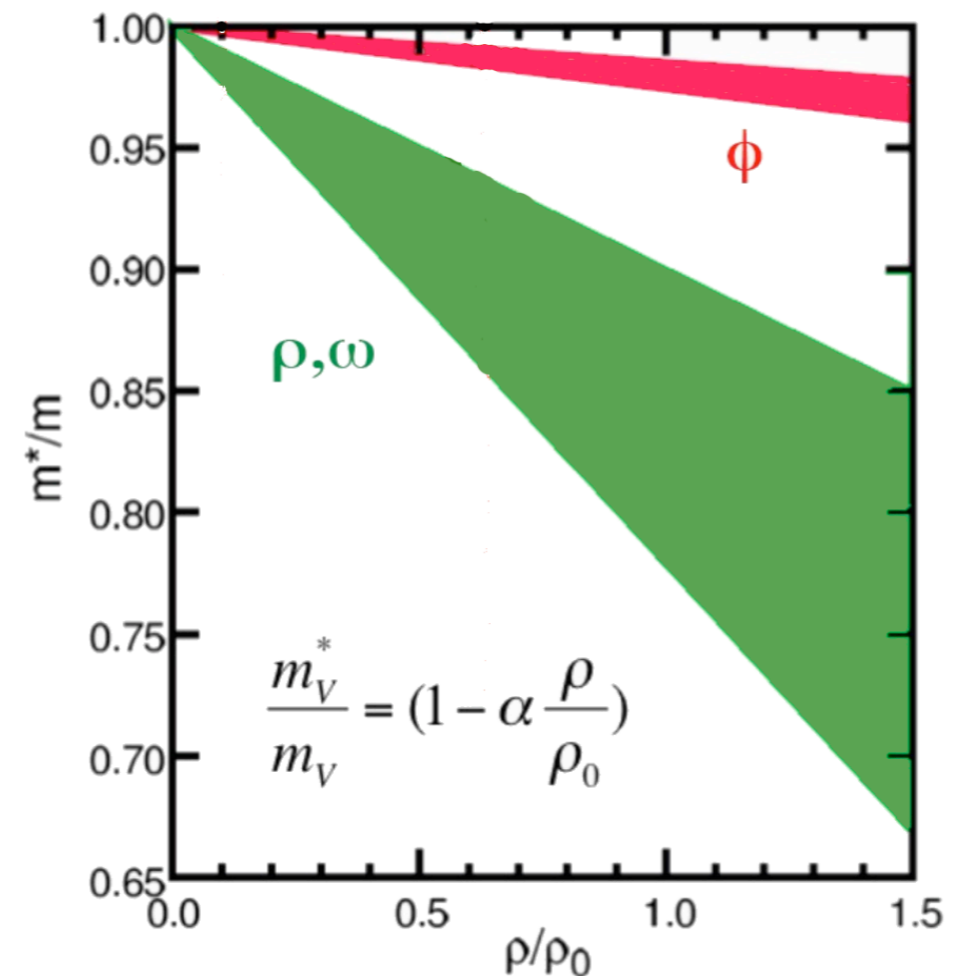
“Brown-Rho Scaling”

G.E.Brown and M. Rho, $\frac{m^*}{m} \approx \frac{\langle \bar{q}q \rangle^*}{\langle \bar{q}q \rangle_0} \approx 0.8(\rho \approx \rho_0)$
PRL 66 (1991) 2720

T.Hatsuda and S. Lee, $\frac{m_V^*}{m_V} = (1 - \alpha \frac{\rho}{\rho_0}); \alpha \approx 0.18$
PRC 46 (1992) R34

QCD sum rule approach:

drop of ρ, ω mass by
about 15% at $\rho = \rho_0$

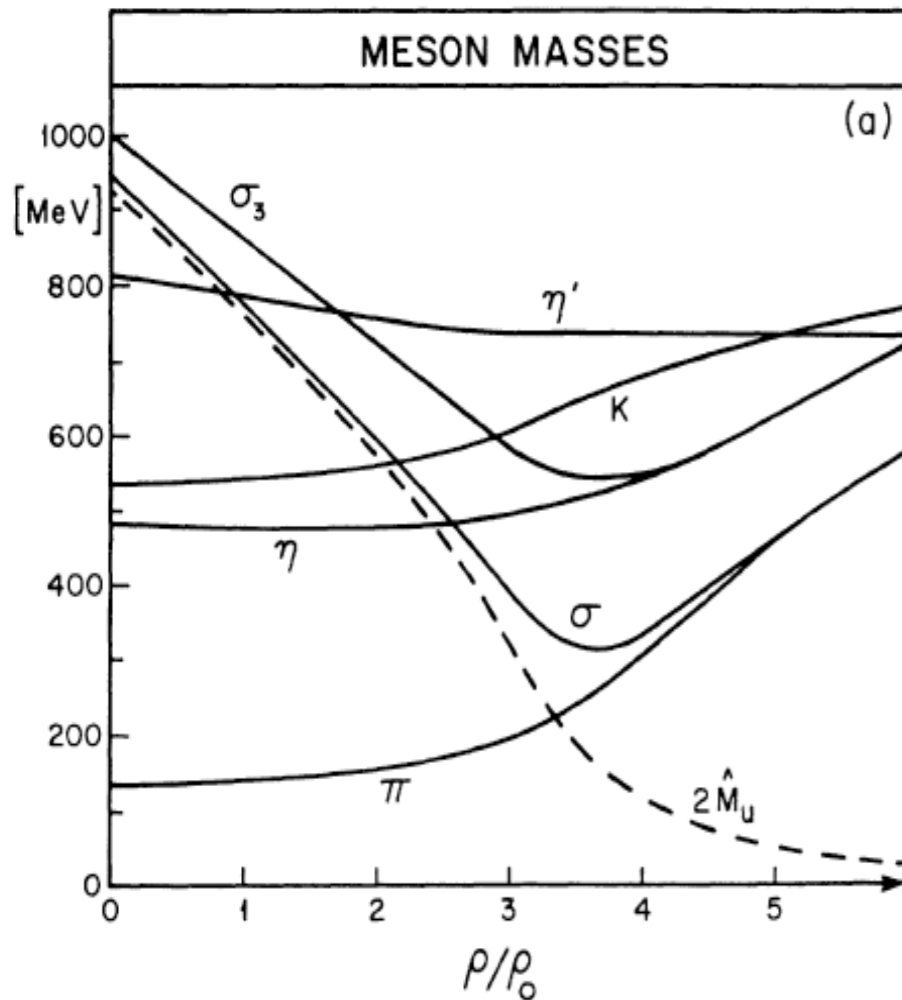


widespread theoretical and
experimental activities to search for
in-medium modifications of hadrons

hadronic models: predictions for η' in-medium mass

NJL-model

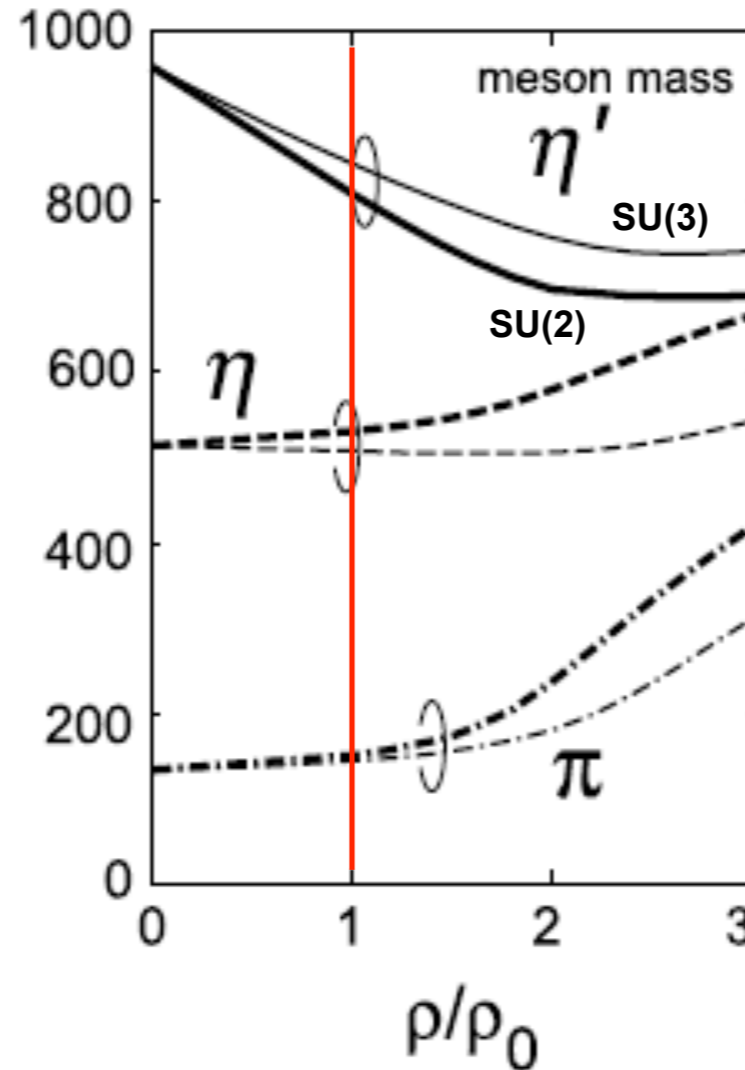
V. Bernard and U.-G. Meissner,
Phys. Rev.D 38 (1988) 1551



almost no dependence of
 η' mass on density

NJL-model

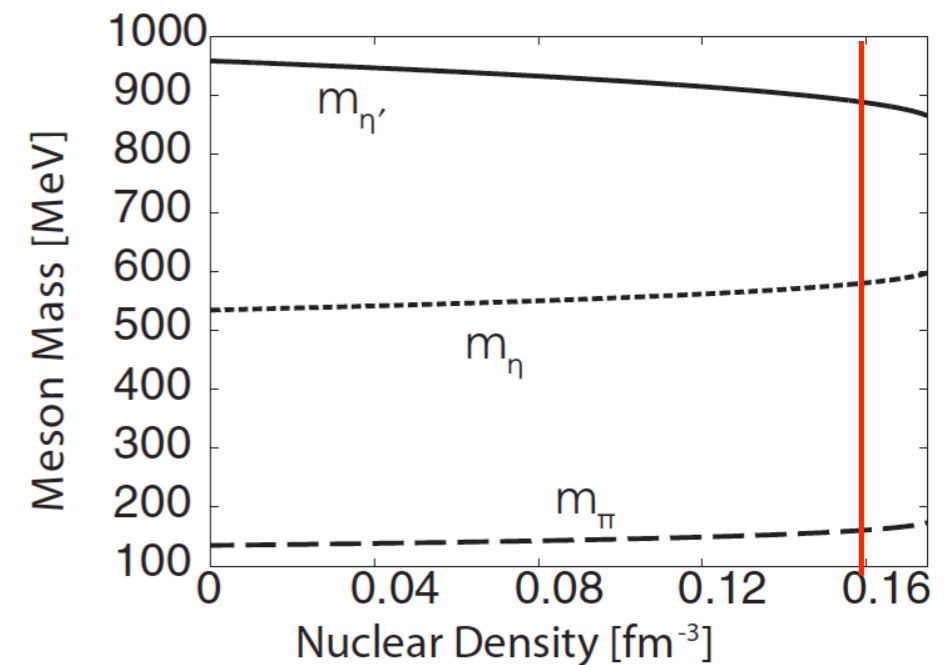
H. Nagahiro et. al,
Phys. Rev. C 74 (2006) 045203



$\Delta m_{\eta'}(\rho_0) \approx -150 \text{ MeV}$
 $\Delta m_{\eta}(\rho_0) \approx +20 \text{ MeV}$

linear σ model

S. Sakai and D. Jido
PRC 88 (2013) 064906



$\Delta m_{\eta'}(\rho_0) \approx -80 \text{ MeV}$

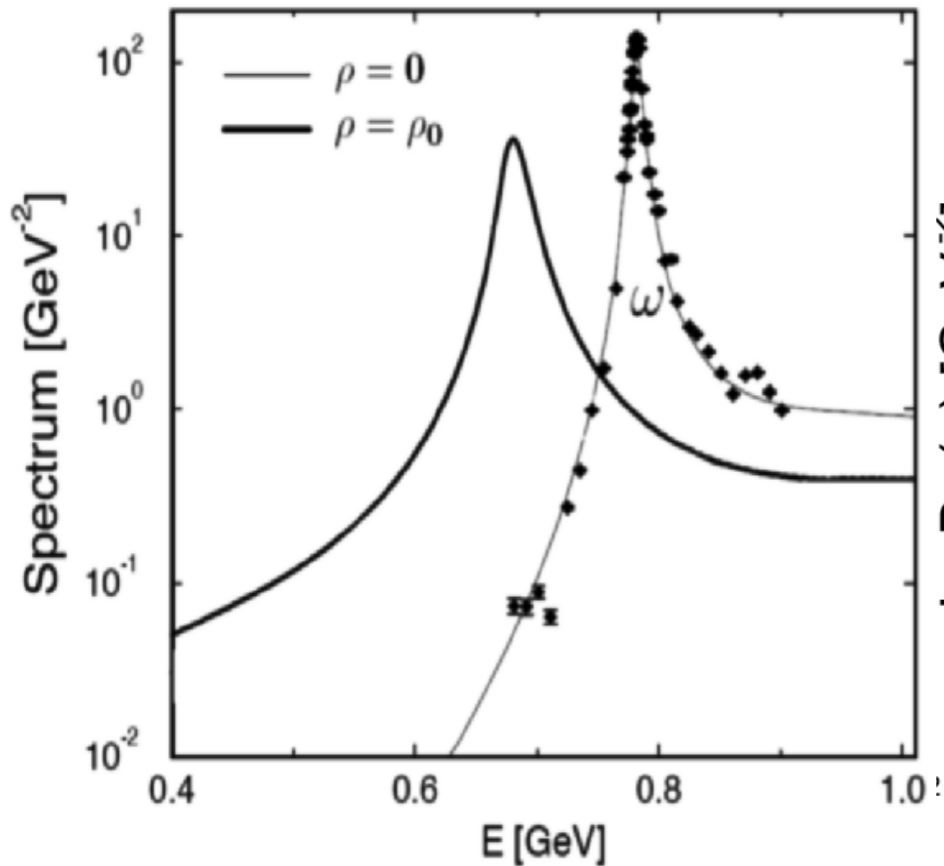
QMC-model

S. Bass and A. Thomas,
PLB 634 (2006) 368

$\Delta m_{\eta'}(\rho_0) \approx -40 \text{ MeV}$
for $\theta_{\eta\eta'} = -20^\circ$

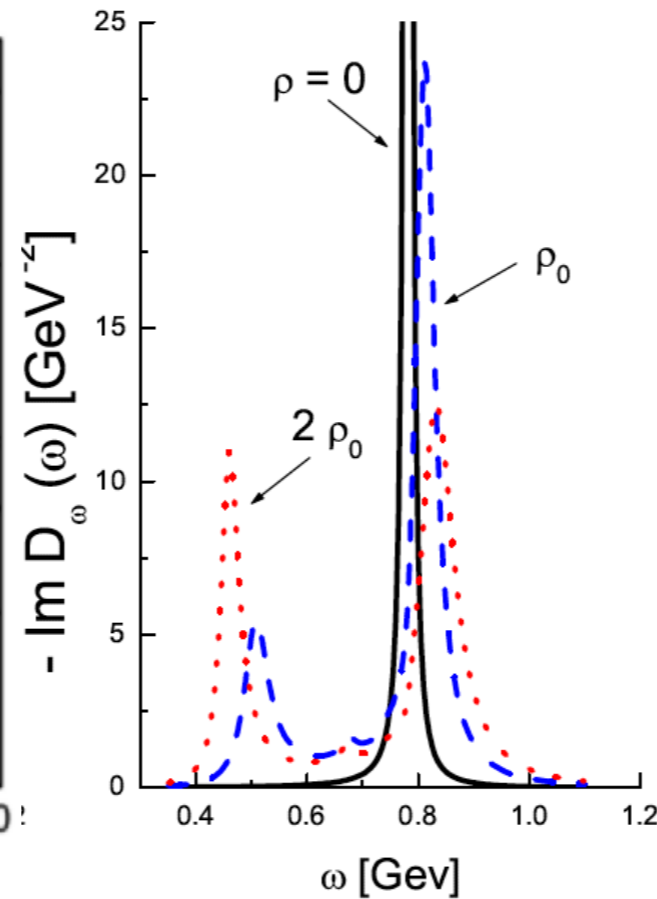
hadronic models: predictions for ω -spectral functions

F. Klingl et al.,
NPA 610 (1997) 297;
NPA 650 (1999) 299



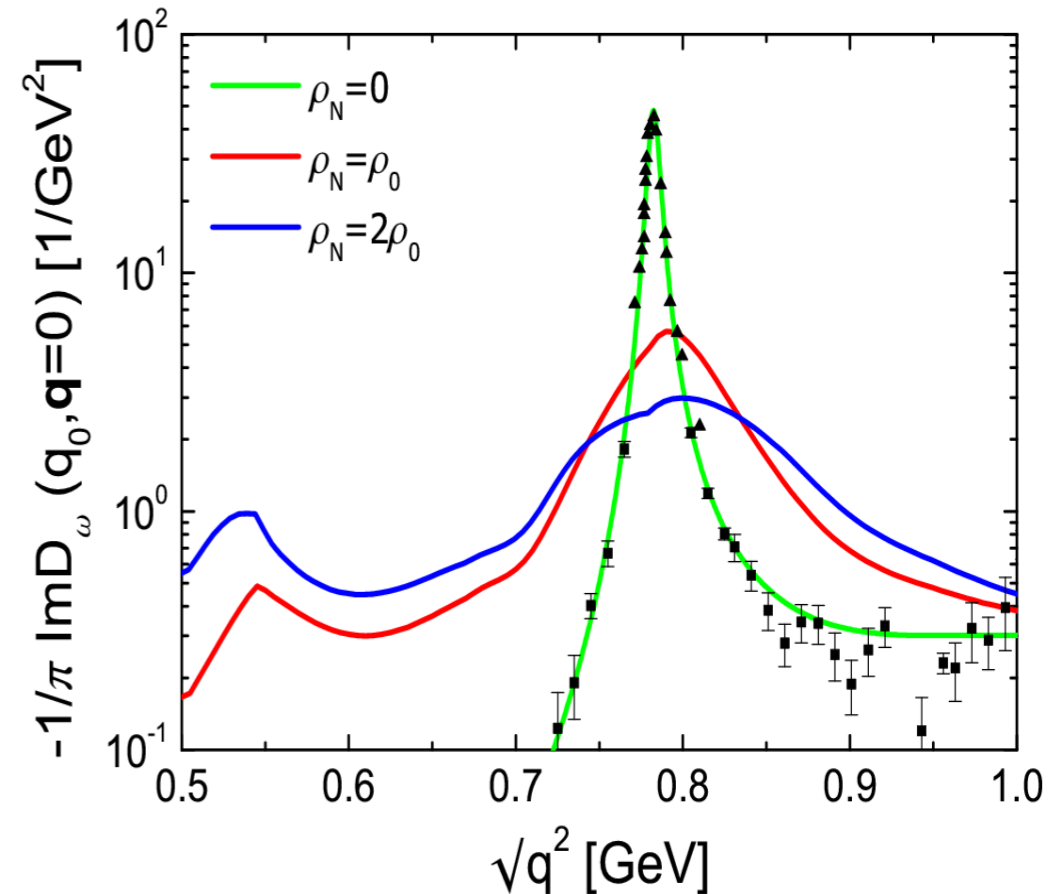
- ◆ lowering of in-medium mass
 - ◆ broadening of resonance with increasing nuclear density
- $\text{Re}(U) \neq 0; \text{Im}(U) \neq 0$

M. Lutz et al.,
NPA 706 (2002) 437



splitting into ω -like and N^*N^{-1} mode due to coupling to nucleon resonances

P. Mühlich et al., NPA 780 (2006) 187



spectral function for ω meson at rest:

almost no mass shift;
strong in-medium broadening
 $\text{Re}(U) \approx 0; \text{Im}(U)$ large

experimental task: search for $\left\{ \begin{array}{l} \text{mass shift?} \\ \text{broadening?} \\ \text{structures?} \end{array} \right\}$ of hadronic spectral functions

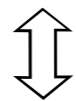
meson-nucleus optical potential

H. Nagahiro and S. Hirenzaki,
PRL 94 (2005) 232503

$$U(r) = V(r) + iW(r)$$

$$V(r) = \Delta m(\rho_0) \cdot \frac{\rho(r)}{\rho_0}$$

real part



in-medium mass modification

$$W(r) = -\Gamma_0/2 \cdot \frac{\rho(r)}{\rho_0} \\ = -\frac{1}{2} \cdot \hbar c \cdot \rho(r) \cdot \sigma_{inel} \cdot \beta$$

imaginary part



lifetime shortened
in-medium width, absorption
inelastic cross section

mass and lifetime (width) may be changed in the medium

experimental approaches to determine the meson-nucleus optical potential

$$U(r) = V(r) + iW(r)$$

real part

$$V(r) = \Delta m(\rho_0) \cdot \frac{\rho(r)}{\rho_0}$$

- ◆ line shape analysis
- ◆ excitation function
- ◆ momentum distribution
- ◆ meson-nucleus bound states

imaginary part

$$W(r) = -\Gamma_0/2 \cdot \frac{\rho(r)}{\rho_0} \\ = -\frac{1}{2} \cdot \hbar c \cdot \rho(r) \cdot \sigma_{inel} \cdot \beta$$

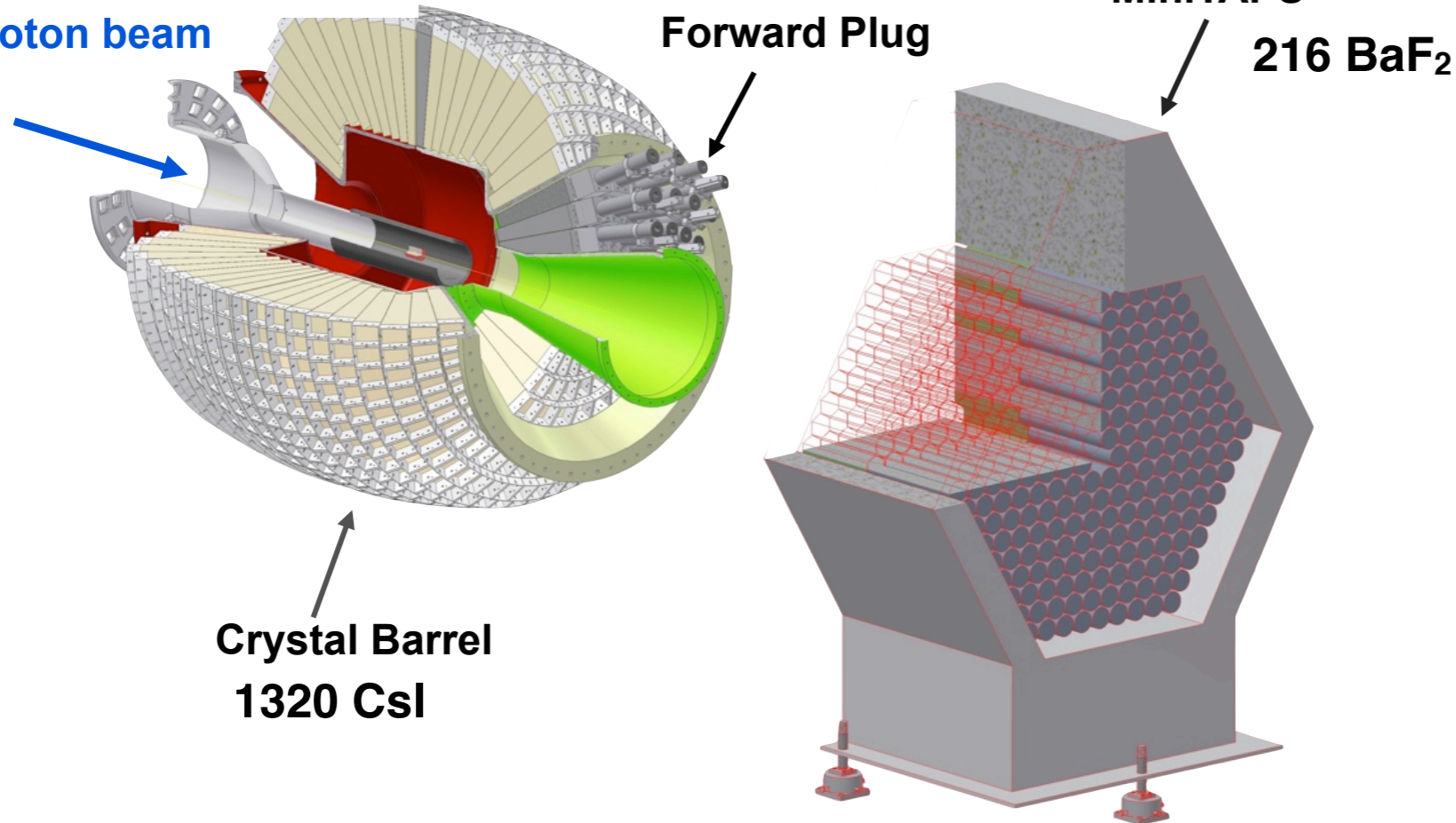
- ◆ transparency ratio measurement

$$T_A = \frac{\sigma_{\gamma A \rightarrow \eta' X}}{A \cdot \sigma_{\gamma N \rightarrow \eta' X}}$$

D. Cabrera et al., NPA 733 (2004)130

CBELSA/TAPS experiment

$E_\gamma = 0.7 - 3.1$ GeV
photon beam



Crystal Barrel
1320 CsI

MiniTAPS
216 BaF₂

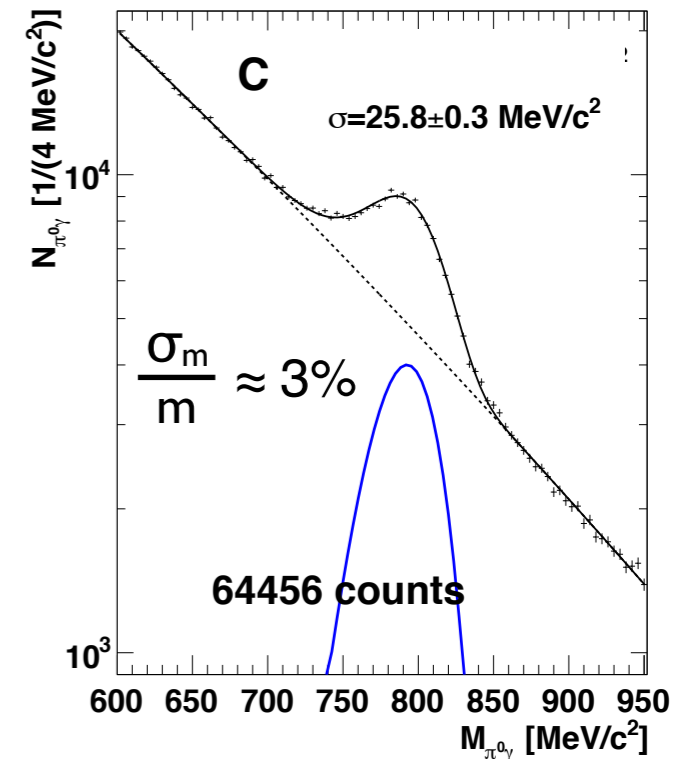
solid target: ¹²C and ⁹³Nb

4π photon detector: ideally suited for identification of multi-photon final states

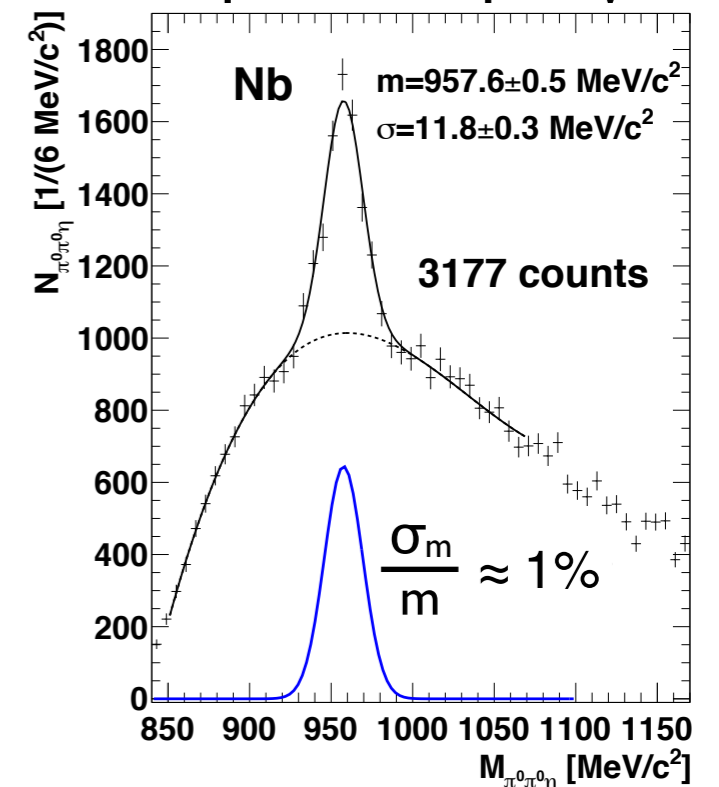
$\omega \rightarrow \pi^0 \gamma \rightarrow 3\gamma$ BR 8.2%

$\eta' \rightarrow \pi^0 \pi^0 \eta \rightarrow 6\gamma$ BR 8.5%

$\omega \rightarrow \pi^0 \gamma \rightarrow 3\gamma$



$\eta' \rightarrow \pi^0 \pi^0 \eta \rightarrow 6\gamma$



The real part of the meson-nucleus
optical potential

the real part of the ω -nucleus potential

J.Weil, U. Mosel and V. Metag, PLB 723 (2013) 120 $\omega \rightarrow \pi^0 \gamma$

sensitive to nuclear density at **production** point and not at **decay** point

◆ measurement of the excitation function of the meson

in case of dropping mass -
higher meson yield for given \sqrt{s}
because of increased phase space
due to lowering of the production threshold

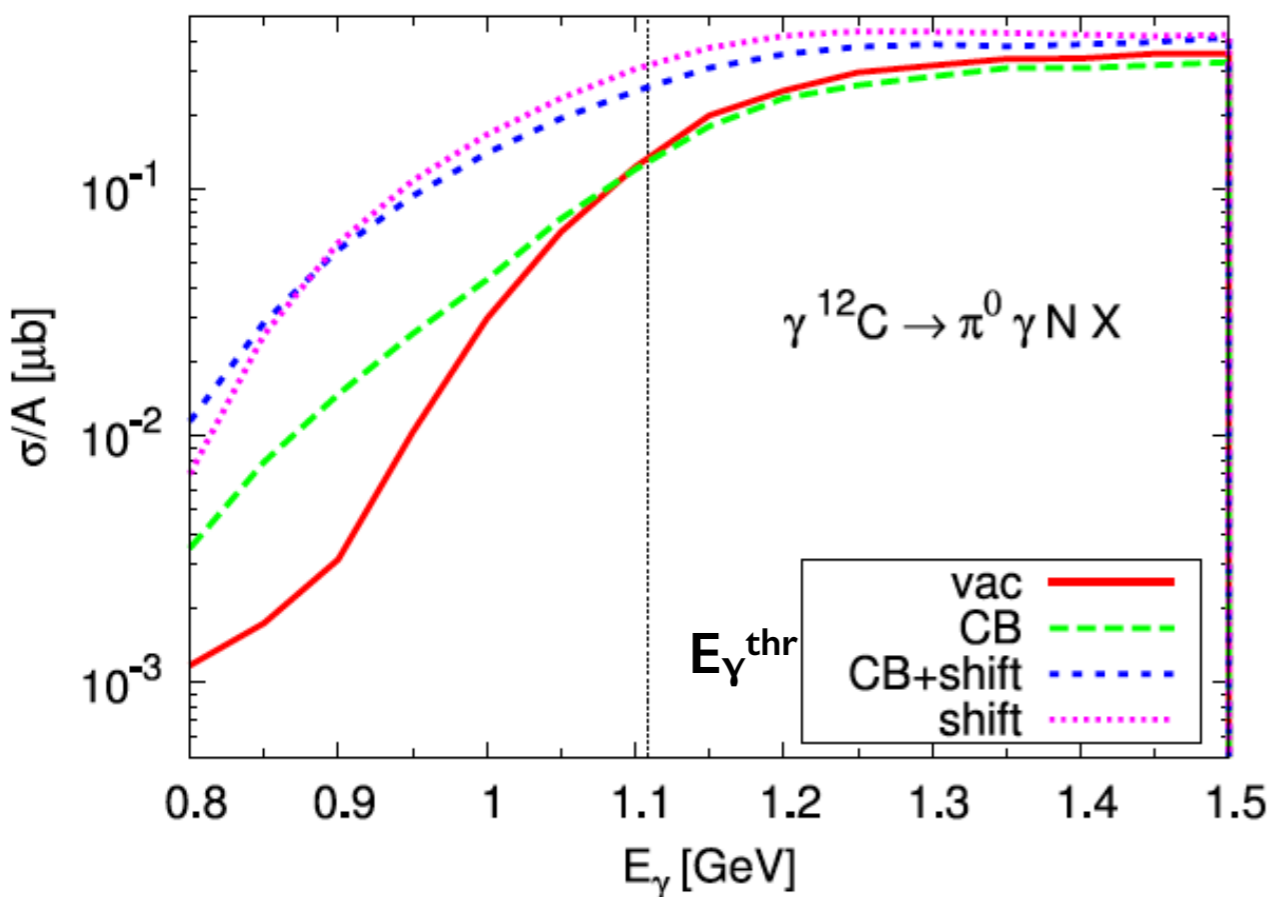
⇒ cross section enhancement

◆ momentum distribution of the meson:

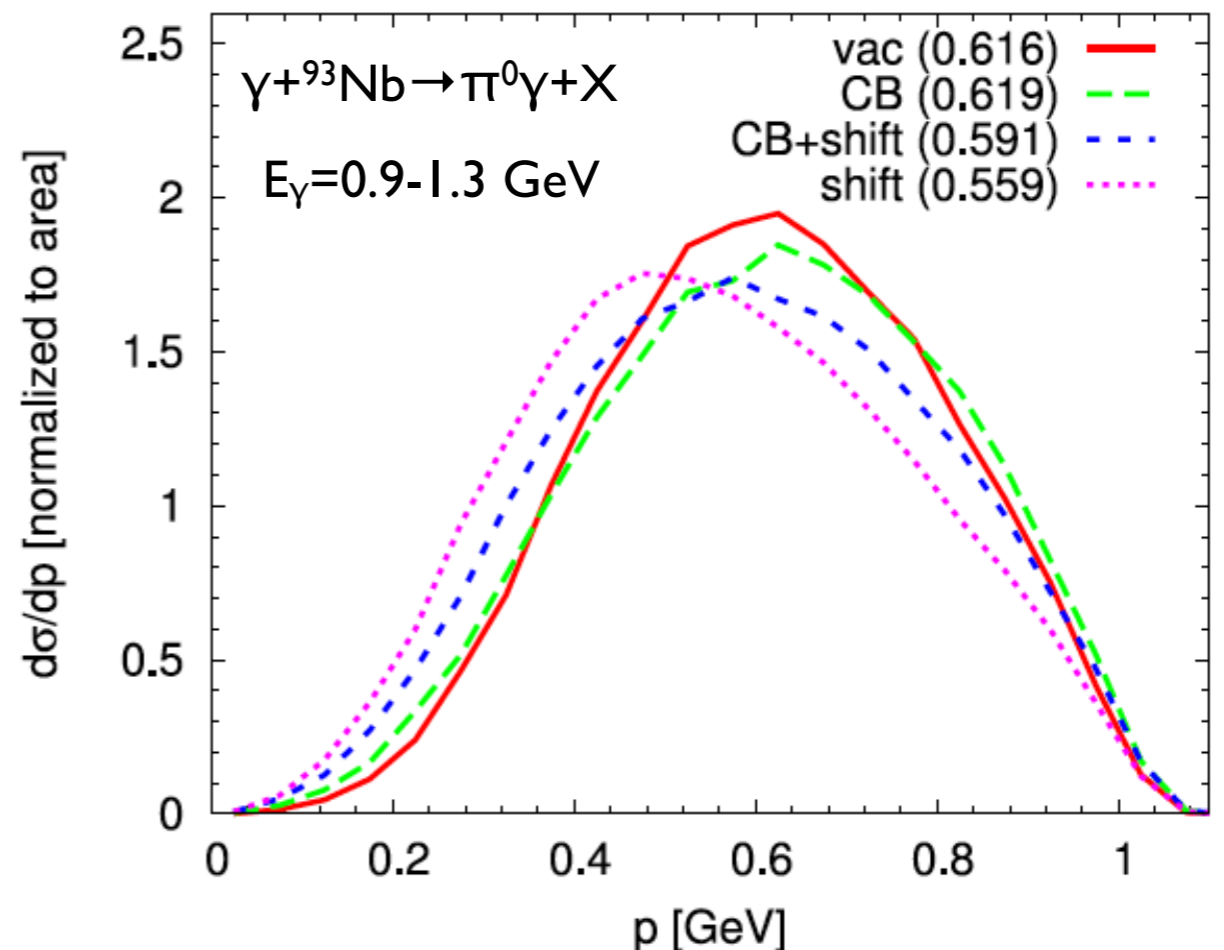
in case of dropping mass - when leaving the nucleus hadron has to become on-shell;
mass generated at the expense of kinetic energy

⇒ downward shift of momentum distribution

$\pi^0 \gamma$ excitation function



$\pi^0 \gamma$ momentum distribution



excitation function for ω photoproduction off C comparison with GiBUU calculation

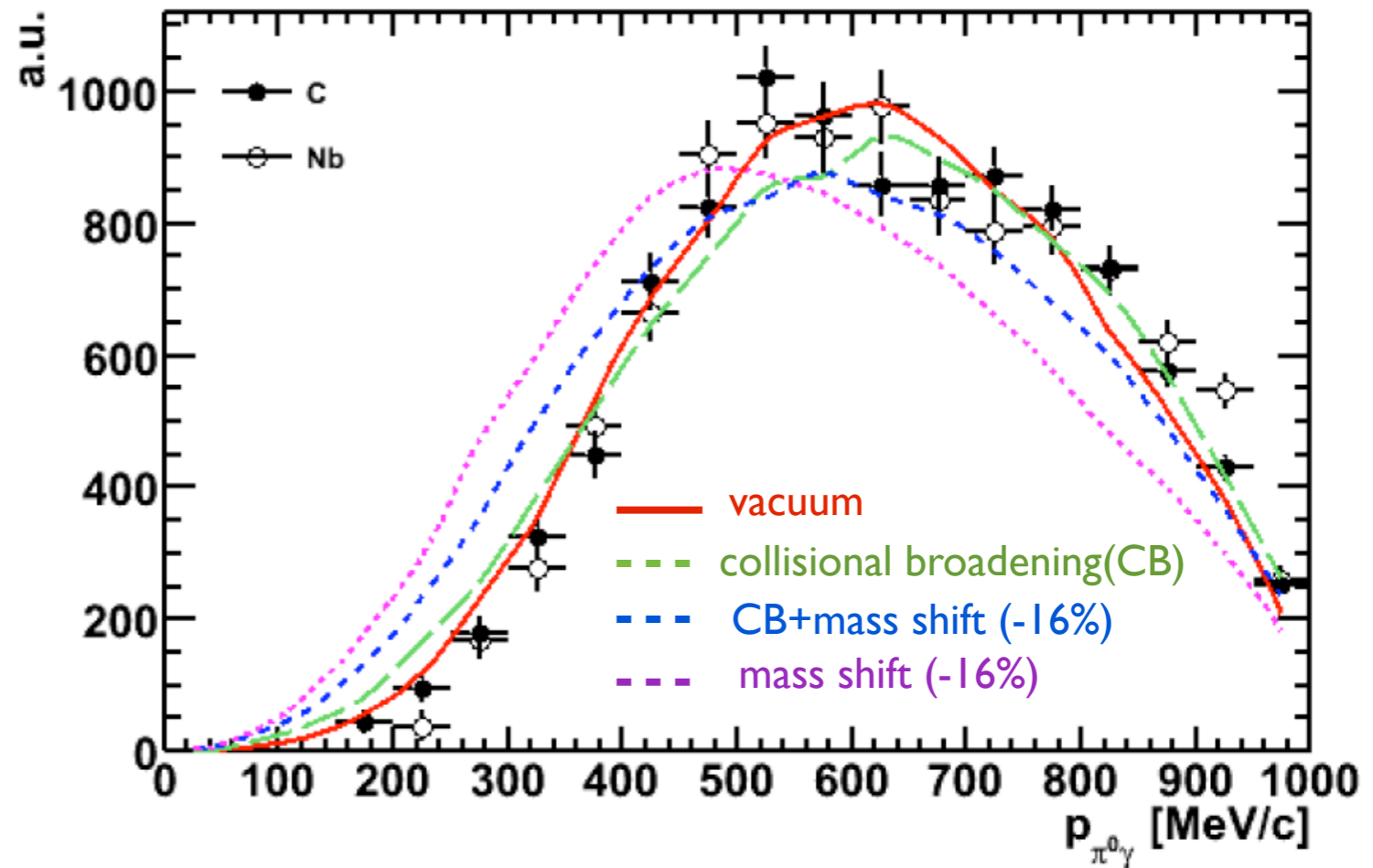
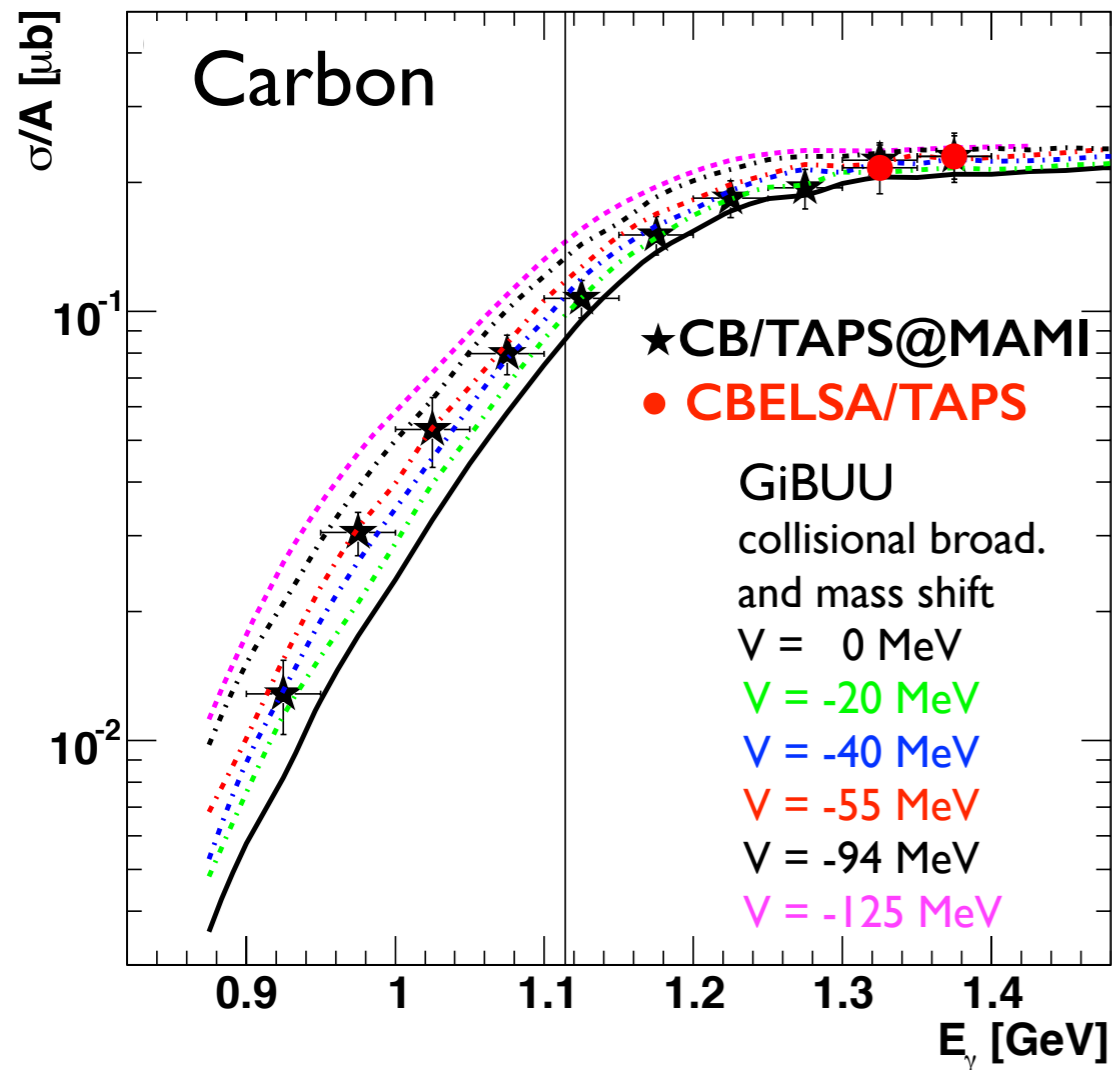
CB/TAPS @ MAMI

V. Metag et al., PPNP, 67 (2012) 530

M.Thiel et al., EPJA 49 (2013) 132

excitation function

momentum distribution



$$V(\rho=\rho_0) = -(42 \pm 17(\text{stat}) \pm 20(\text{syst})) \text{ MeV}$$

data not consistent with strong mass shift scenario ($\Delta m/m \approx -16\%$)

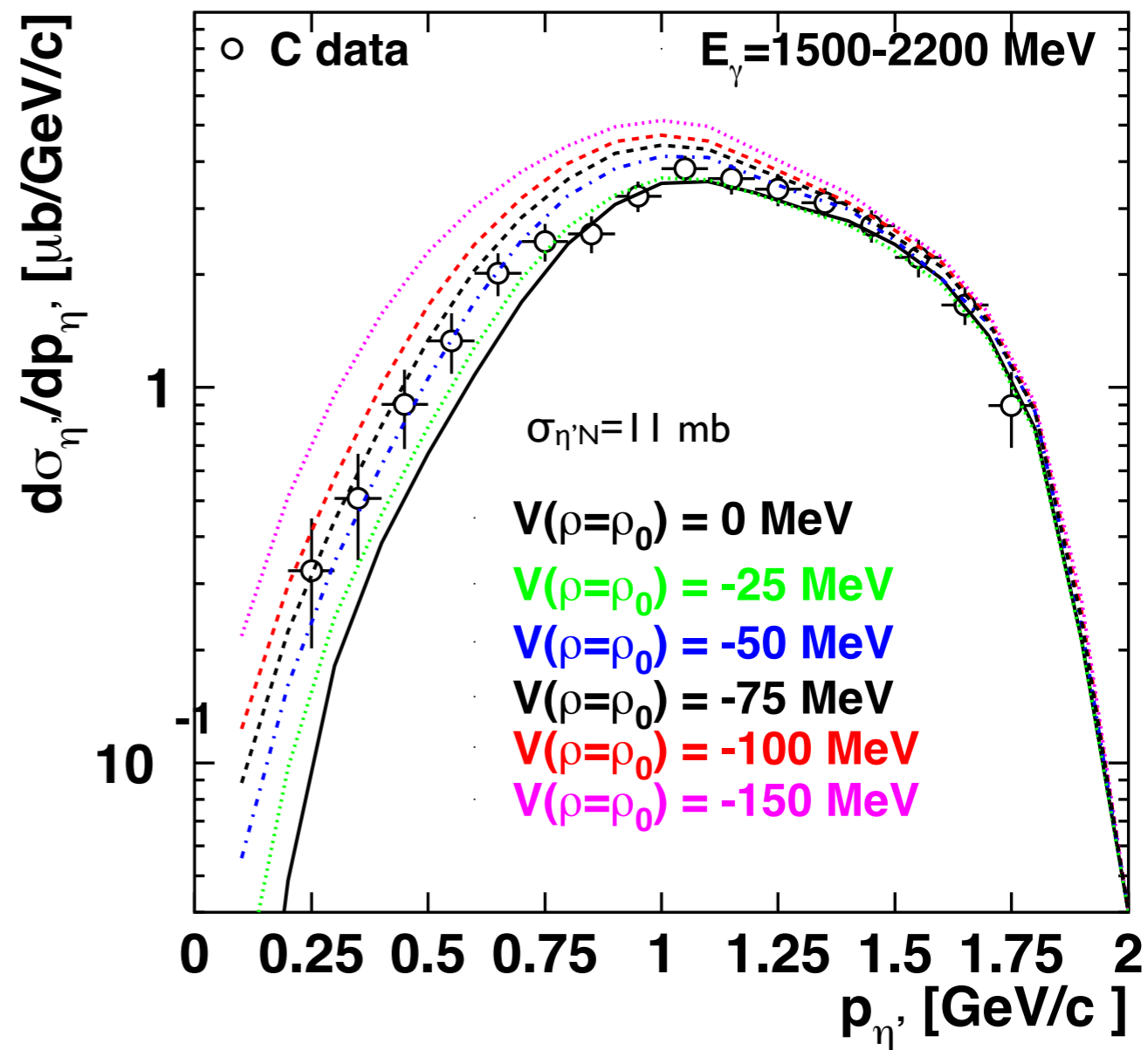
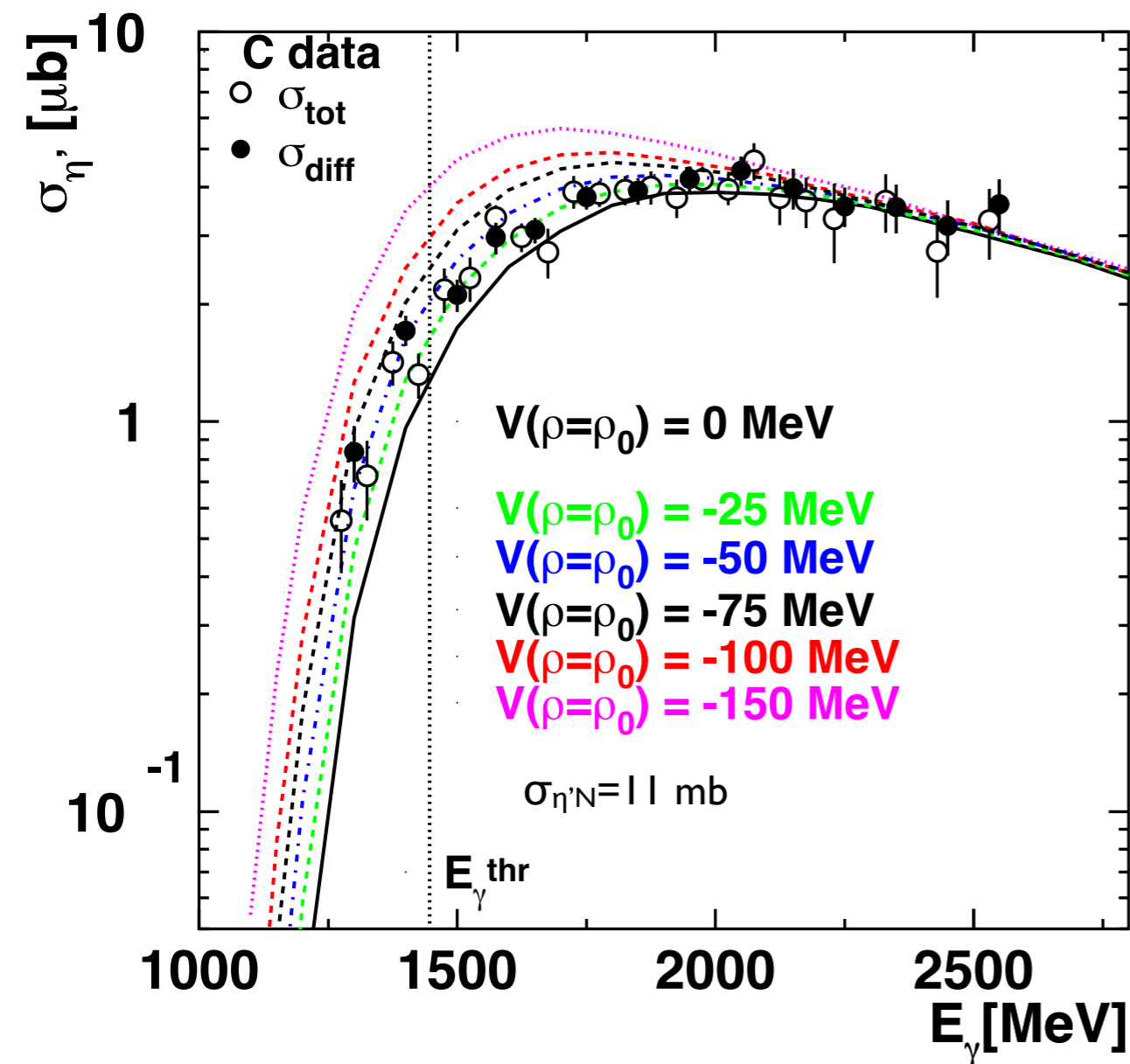
excitation function and momentum distribution for η' photoproduction off C

CBELSA/TAPS @ ELSA

data: M. Nanova et al., PLB 727 (2013) 417

$\gamma C \rightarrow \eta' X$

calc.: E. Paryev, J. Phys. G 40 (2013) 025201



$$V_{\eta'}(\rho=\rho_0) = -(40 \pm 6) \text{ MeV}$$

$$V_{\eta'}(p_{\eta'} \approx 1.1 \text{ GeV}/c; \rho=\rho_0) = -(32 \pm 11) \text{ MeV}$$

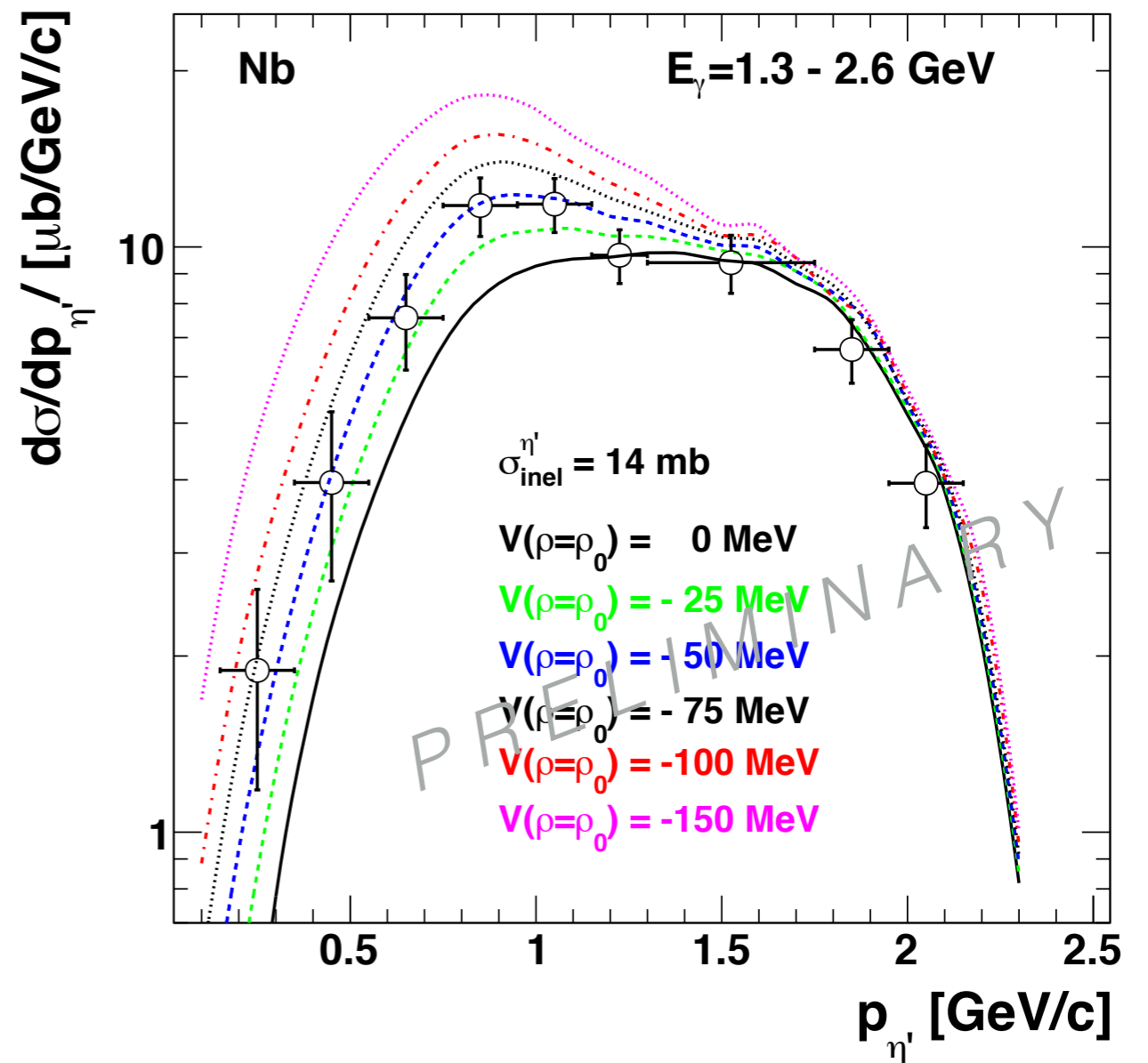
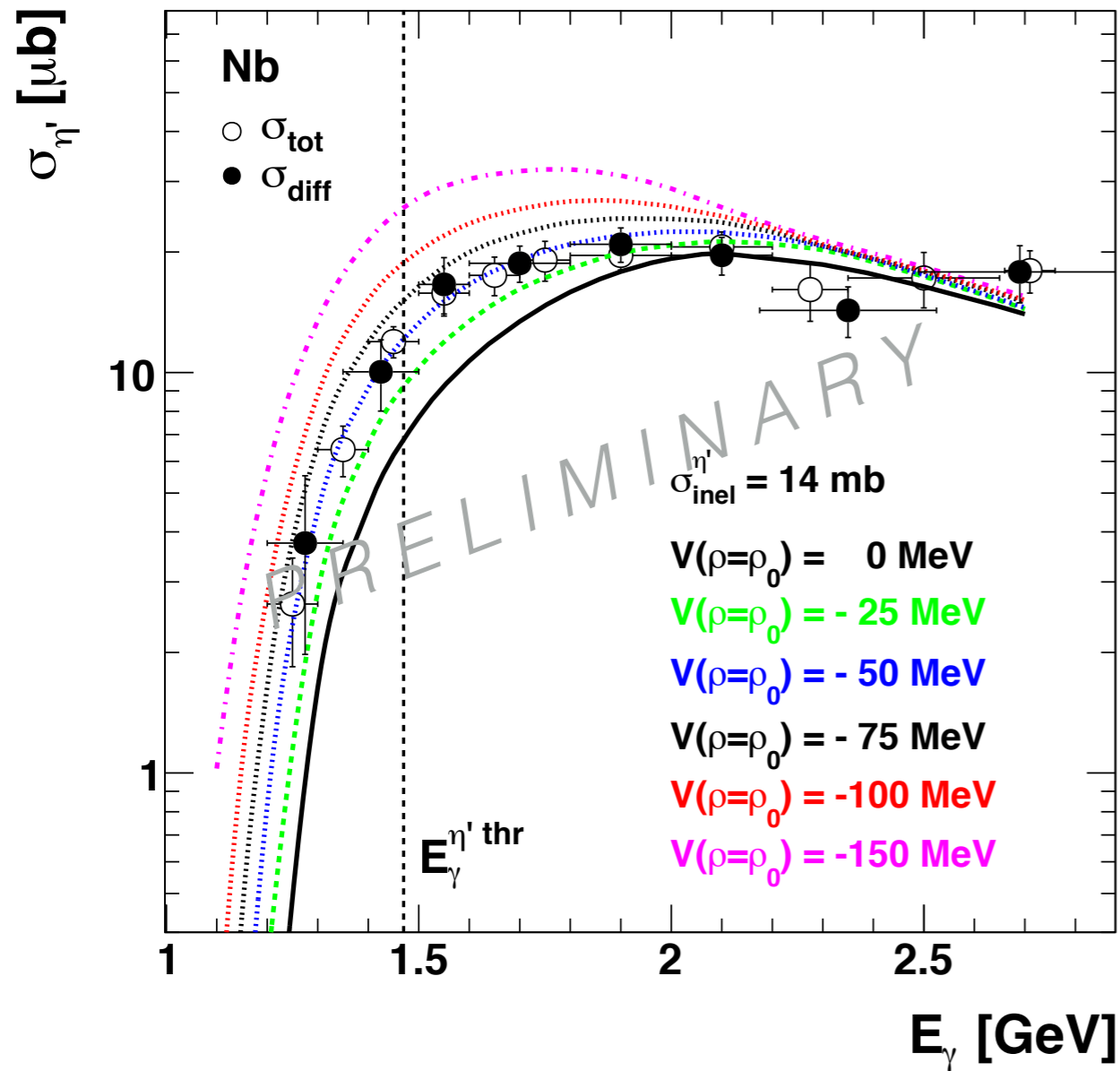
data disfavour strong mass shifts

excitation function and momentum distribution for η' photoproduction off Nb

CBELSA/TAPS @ ELSA

γ Nb \rightarrow η' X

M. Nanova et al., submitted to PRC for publication



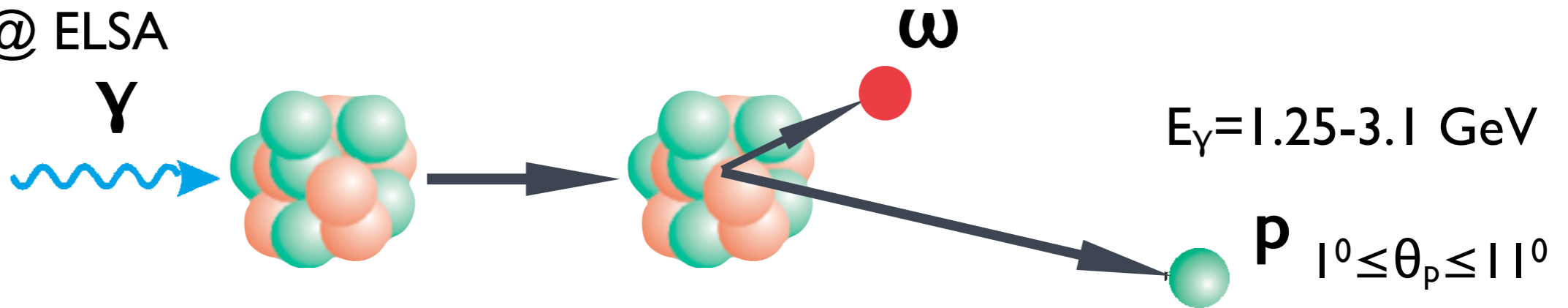
$$V_{\eta'}(\rho=\rho_0) = -(46 \pm 15) \text{ MeV}$$

$$V_{\eta'}(p_{\eta'} \approx 1.14 \text{ GeV}/c; \rho=\rho_0) = -(41 \pm 22) \text{ MeV}$$

data disfavour strong mass shifts

real part of ω -nucleus potential from ω kinetic energy

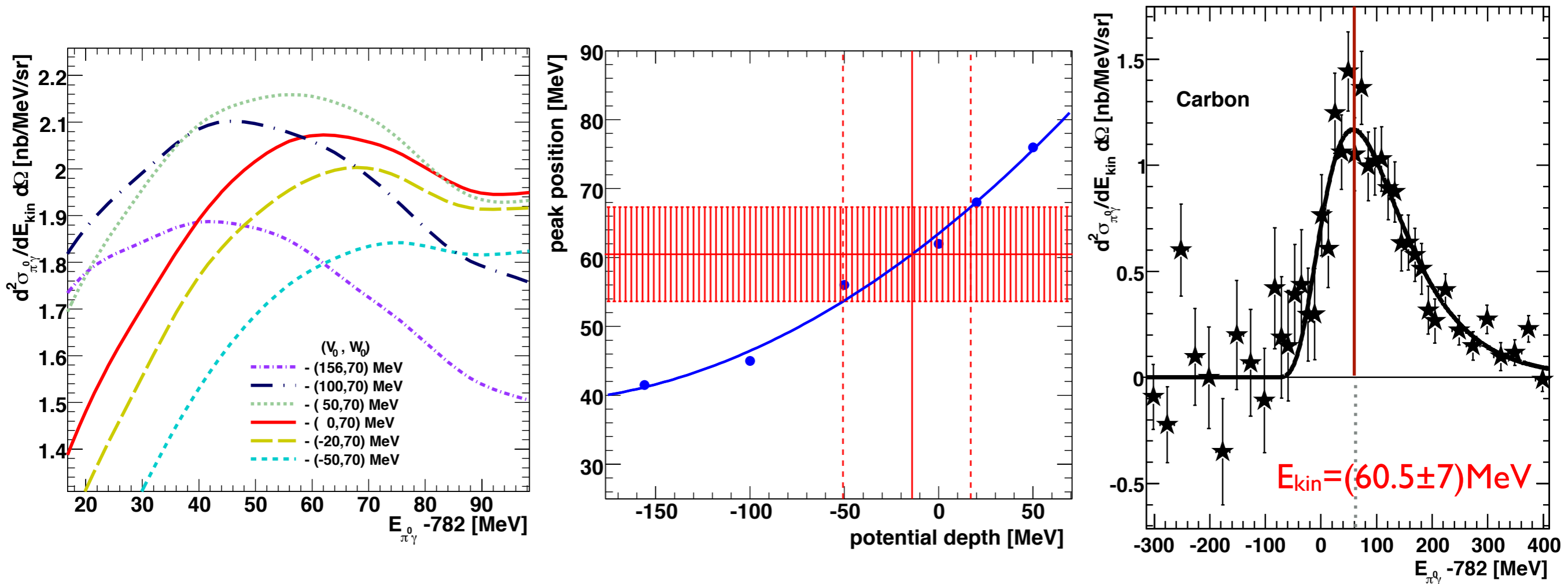
CBELSA/TAPS @ ELSA



the higher the attraction the lower the kinetic energy of the ω meson

H. Nagahiro, priv. com.

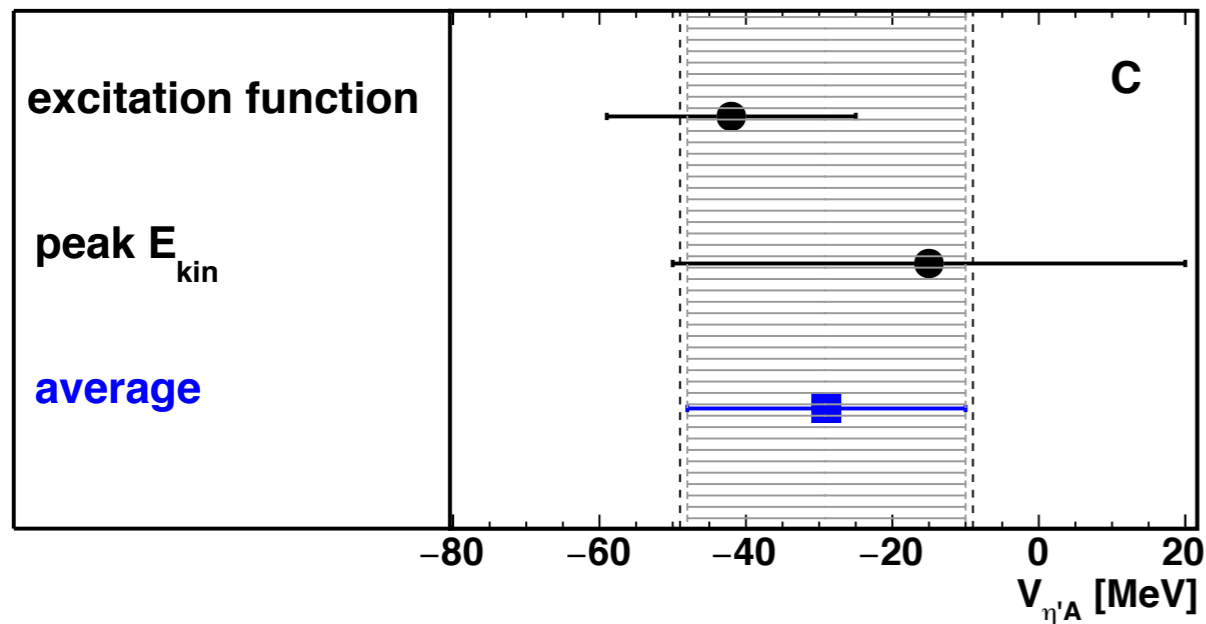
S. Friedrich et al., PLB 736 (2014) 26



$$V_\omega(p_\omega \approx 300 \text{ MeV}/c; \rho = \rho_0) = -(15 \pm 35) \text{ MeV}$$

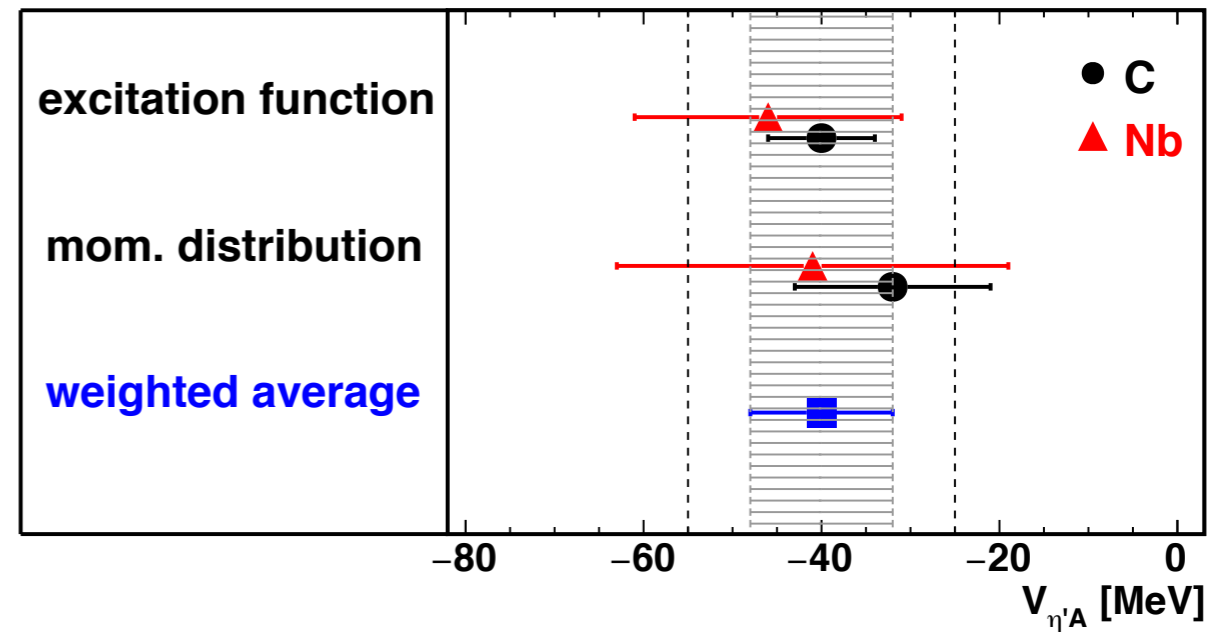
compilation of results for the real part of the ω - and η' -nucleus optical potential

ω



$$V_{\omega A}(\rho=\rho_0) = -(29 \pm 19(\text{stat}) \pm 20(\text{syst})) \text{ MeV}$$

η'



$$V_{\eta' A}(\rho=\rho_0) = -(40 \pm 8(\text{stat}) \pm 15(\text{syst})) \text{ MeV}$$

The imaginary part of the meson-nucleus optical potential: momentum dependence

momentum differential cross section for ω, η' produced off C, Nb

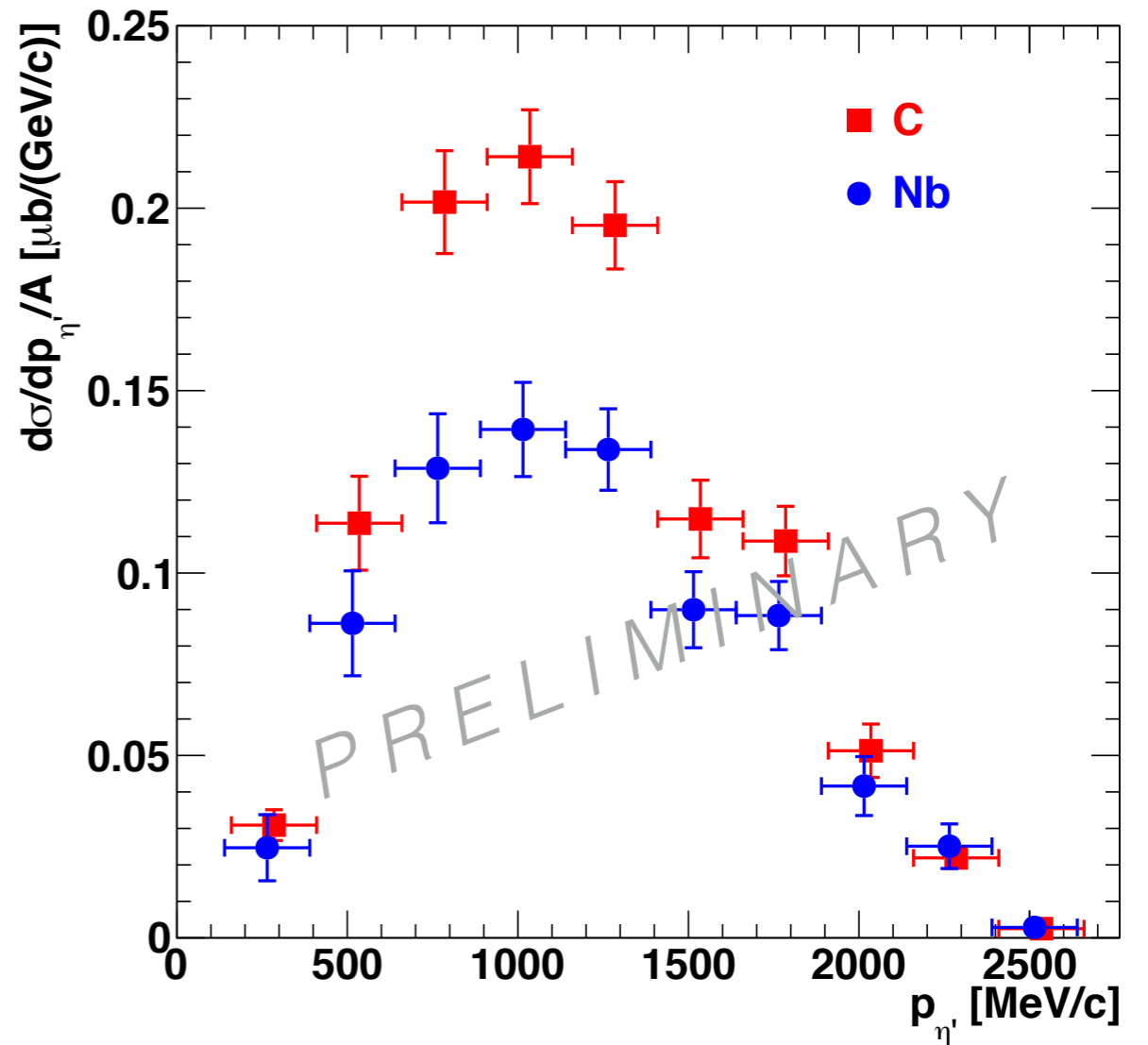
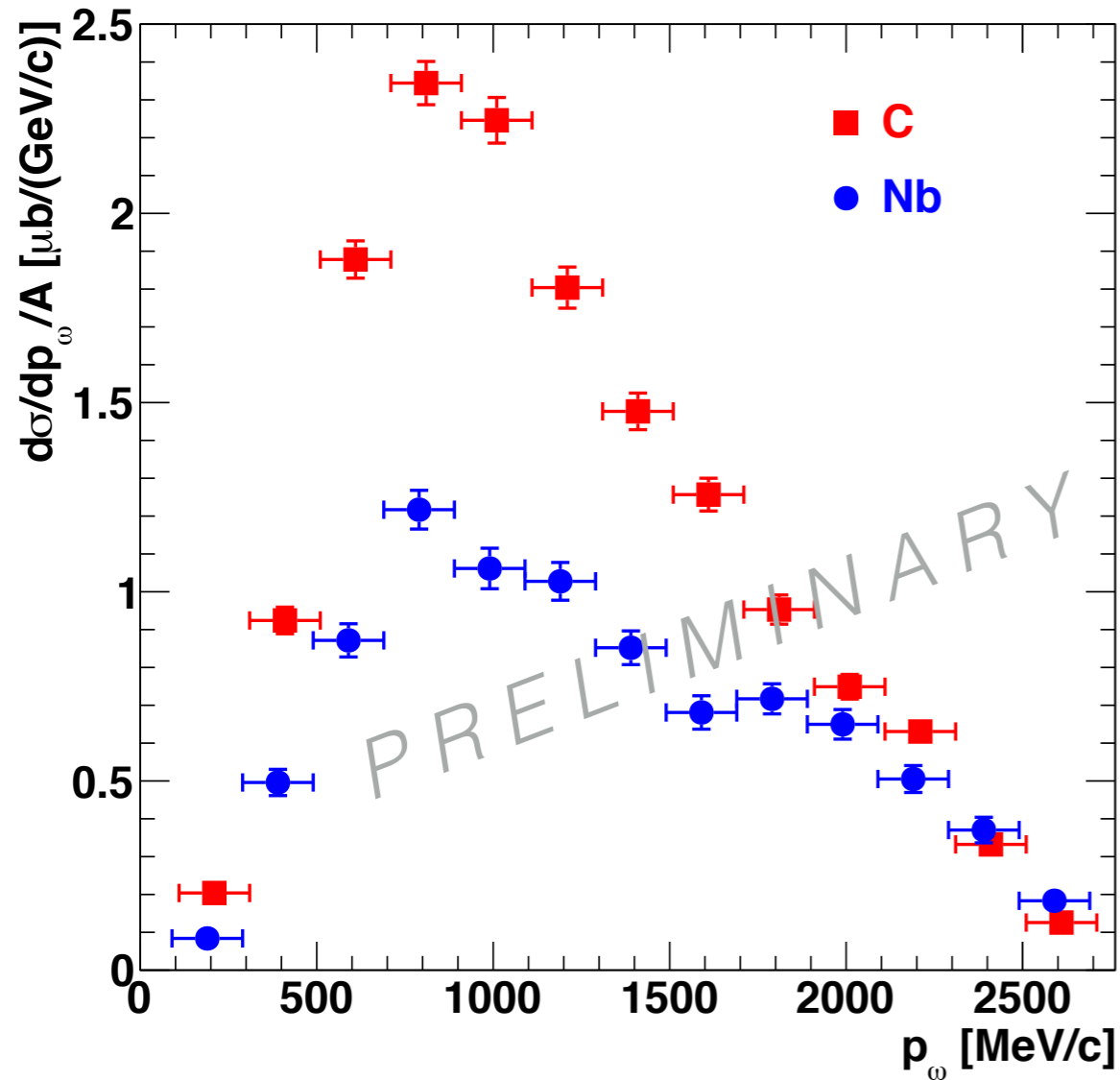
ω

$E_\gamma = 1.2\text{-}2.9\text{ GeV}$

η'

$\gamma\text{ C,Nb} \rightarrow \omega X$

$\gamma\text{ C,Nb} \rightarrow \eta' X$



momentum differential cross sections $\Rightarrow T_{\text{Nb/C}}^m(p_m) = \frac{12 \cdot \sigma_{\gamma\text{Nb} \rightarrow mX}(p_m)}{93 \cdot \sigma_{\gamma\text{C} \rightarrow mX}(p_m)}$

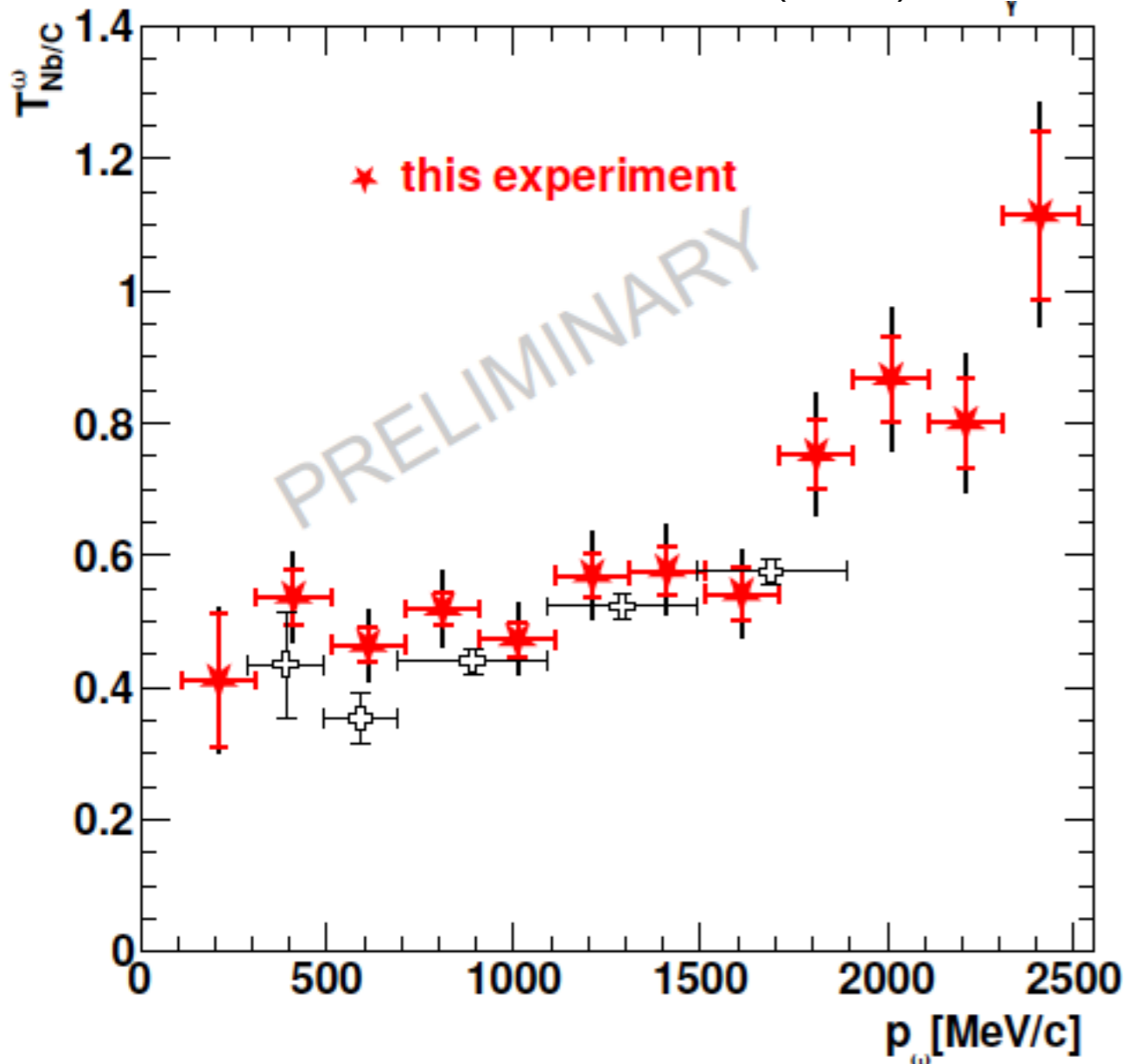
momentum dependence of transparency ratio for ω , η'

ω

$$T_{\text{Nb/C}}^m(p_m) = \frac{12 \cdot \sigma_{\gamma\text{Nb} \rightarrow mX}(p_m)}{93 \cdot \sigma_{\gamma\text{C} \rightarrow mX}(p_m)}$$

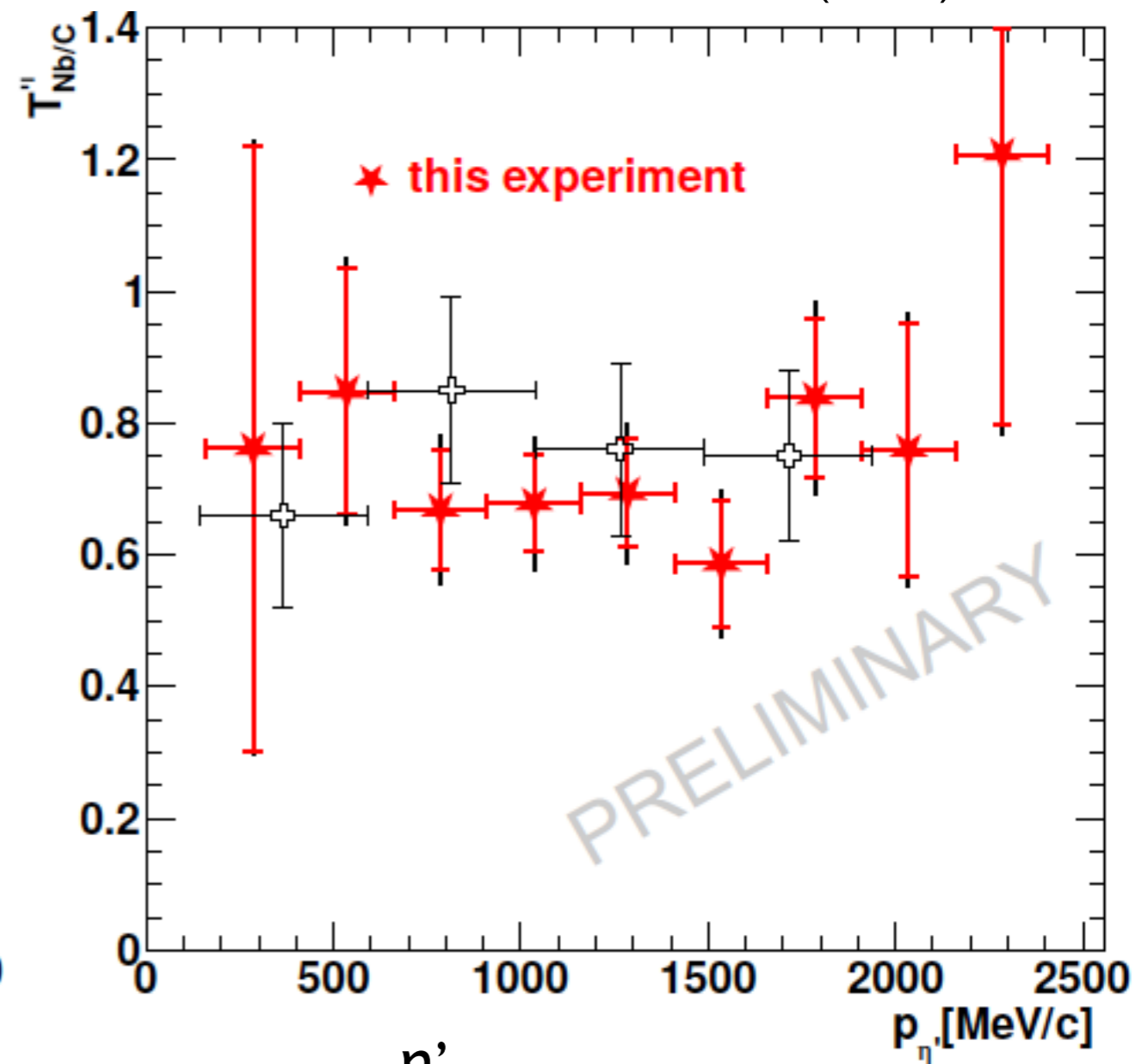
η'

⊕ M. Kotulla et al., PRL 100 (2008) 192302



$$T_{\text{Nb/C}}^\omega \approx 0.4-0.6$$

⊕ M. Nanova et al., PLB 710 (2012) 600



$$T_{\text{Nb/C}}^{\eta'} \approx 0.7-0.8$$

absorption of η' mesons much weaker than for ω mesons !!

imaginary part of the potential for ω, η'

Glauber model: high energy Eikonal approximation

$$T_{\text{Nb/C}}^m(p_m) \Rightarrow \Gamma_0^m(\rho=\rho_0)(p_m) = -2 \text{Im } U_0^m(p_m)$$

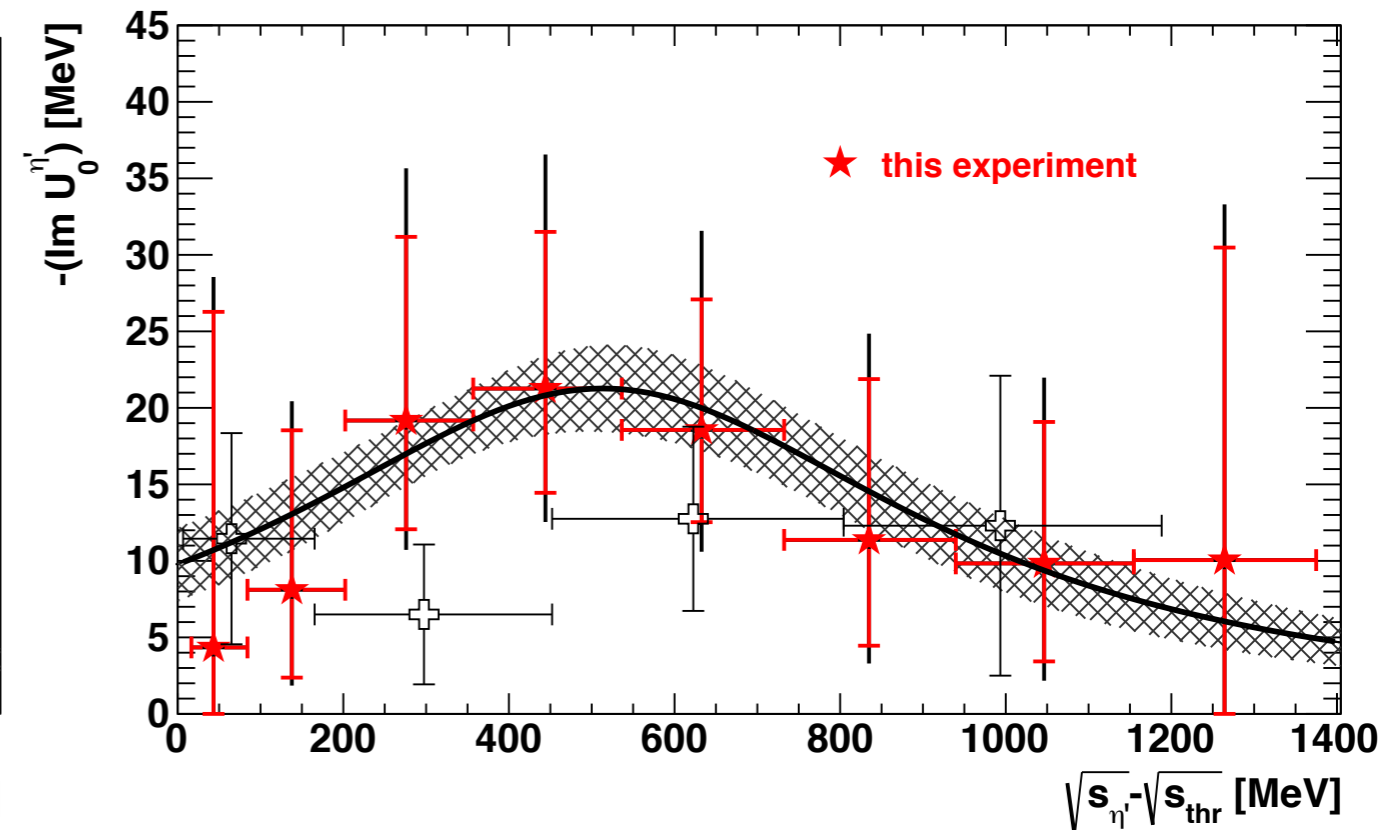
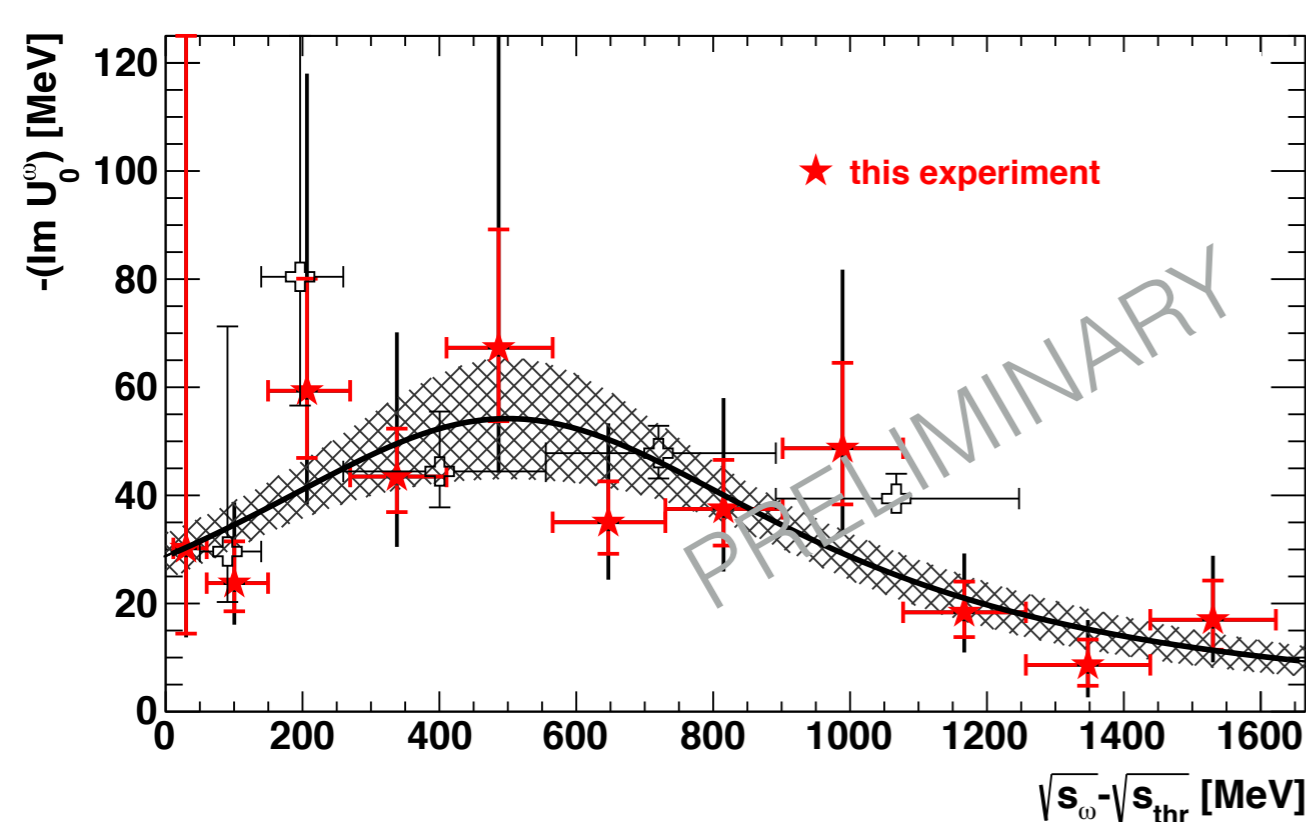
S. Friedrich et al., submitted to EPJA for publication

ω

η'

⊕ M. Kotulla et al., PRL 114 (2015) 199903

⊕ M. Nanova et al., PLB 710 (2012) 600



◆ extrapolation to production threshold:

$$\text{Im } U_0^{\omega}(\rho=\rho_0, p_{\omega}=0) = -(30 \pm 10) \text{ MeV}$$

$$\text{Im } U_0^{\eta'}(\rho=\rho_0, p_{\eta'}=0) = -(10 \pm 3) \text{ MeV}$$

◆ extension to higher energies allows for dispersion relation analysis, providing link between real and imaginary part of potential

compilation of results for real and imaginary part of the ω, η' -nucleus optical potential

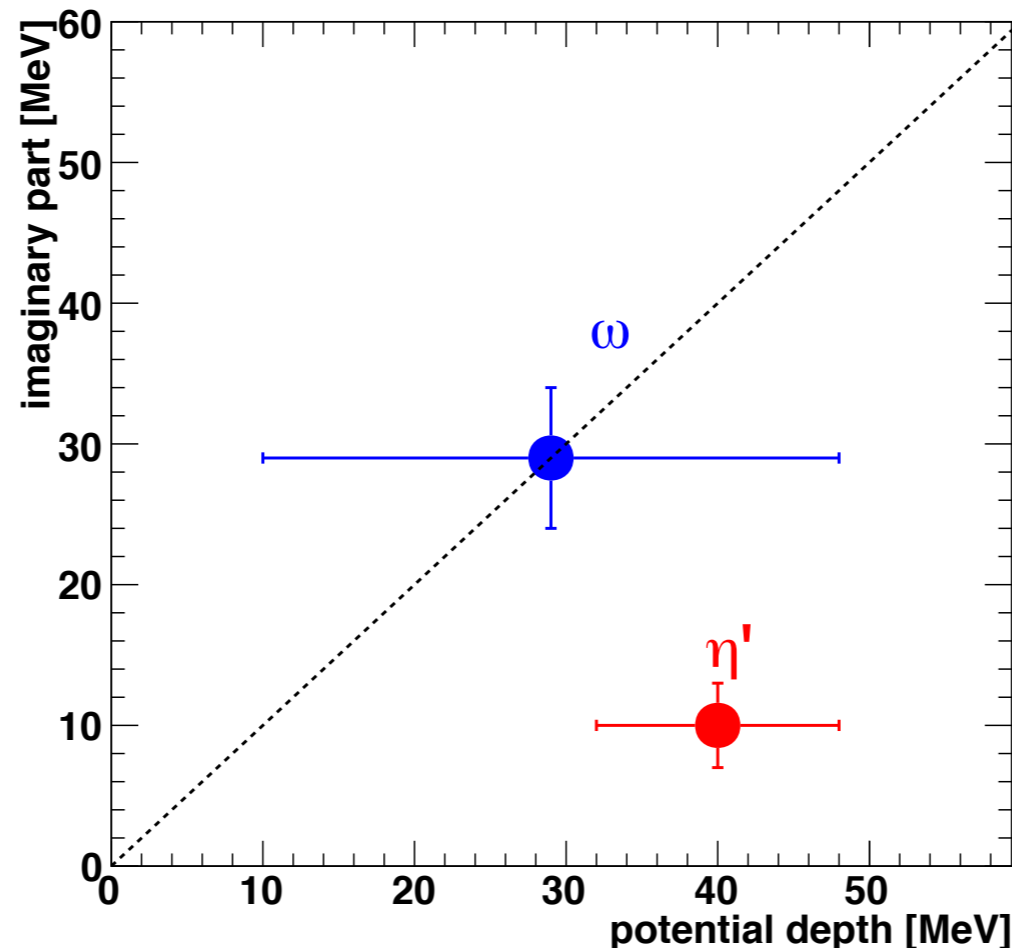
ω

$$U_{\omega A}(\rho=\rho_0)=$$

$$U_{\eta' A}(\rho=\rho_0)=$$

η'

$$-((29 \pm 19(\text{stat}) \pm 20(\text{syst}) + i(30 \pm 10)) \text{ MeV} \quad -((40 \pm 8(\text{stat}) \pm 15(\text{syst}) + i(10 \pm 3)) \text{ MeV}$$



$|\text{Im } U| \approx |\text{Re } U|; \Rightarrow \omega$ not a good candidate
to search for meson-nucleus bound states!

$|\text{Re } U| \gg |\text{Im } U|; \Rightarrow \eta'$ promising candidate
to search for mesic states

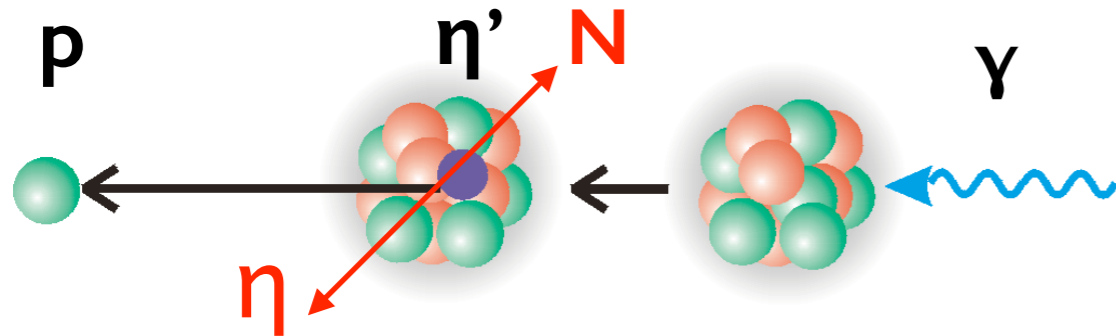
first (indirect) observation of in-medium mass shift of η' at $\rho=\rho_0$ and $T=0$

in good agreement with QMC model predictions (S. Bass et al., PLB 634 (2006) 368) 20

outlook: search for η' -mesic states in photo-nuclear reactions

B1: BGO-OD@ELSA

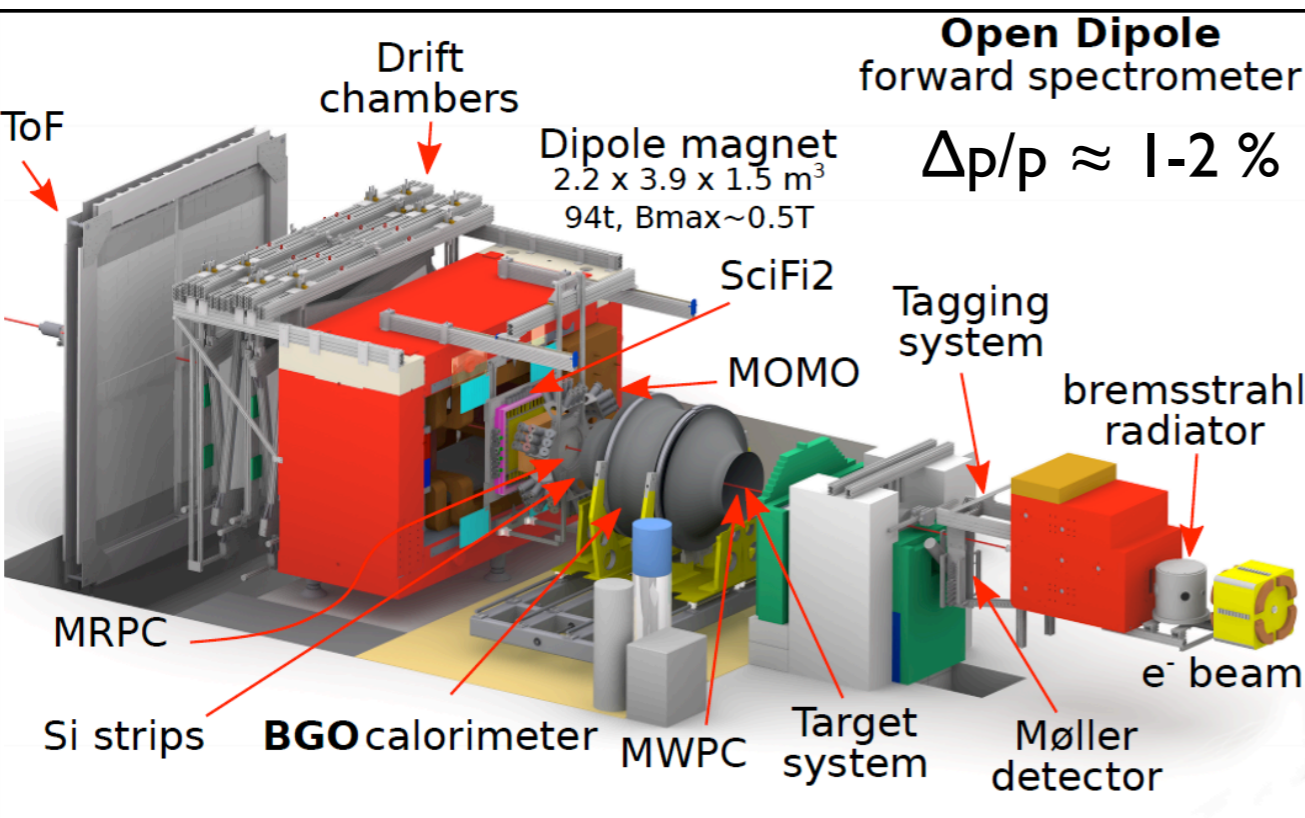
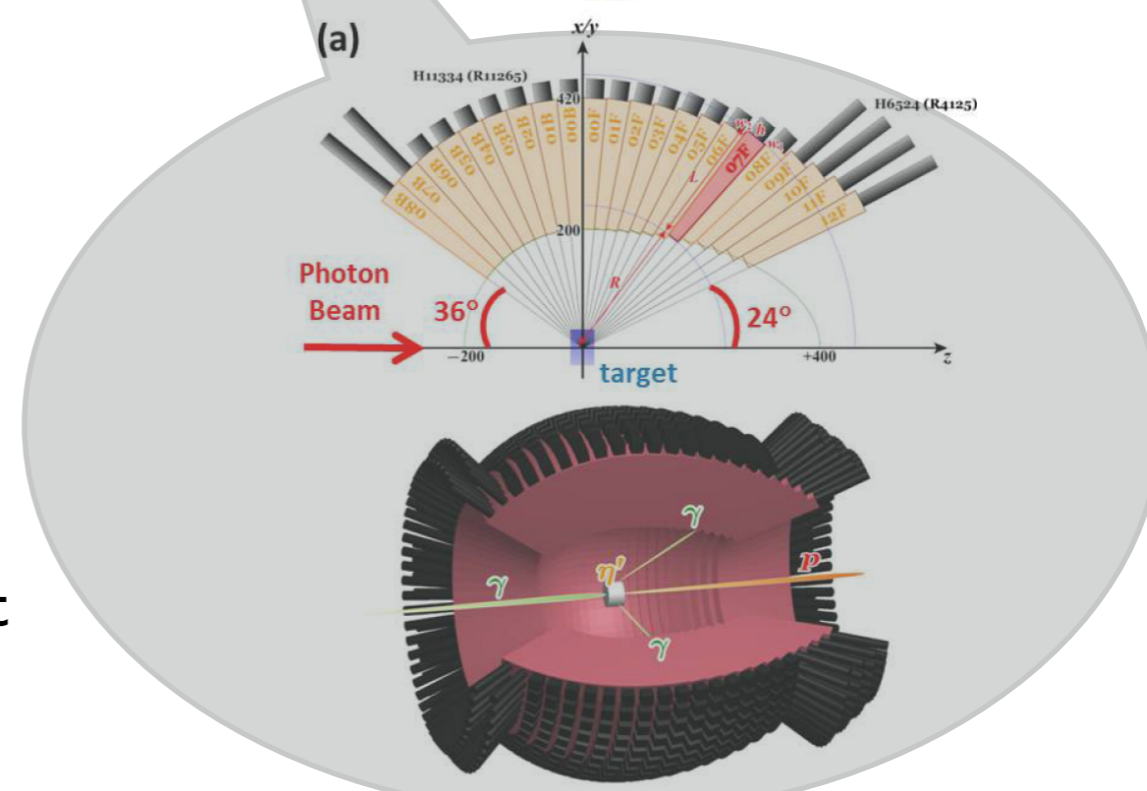
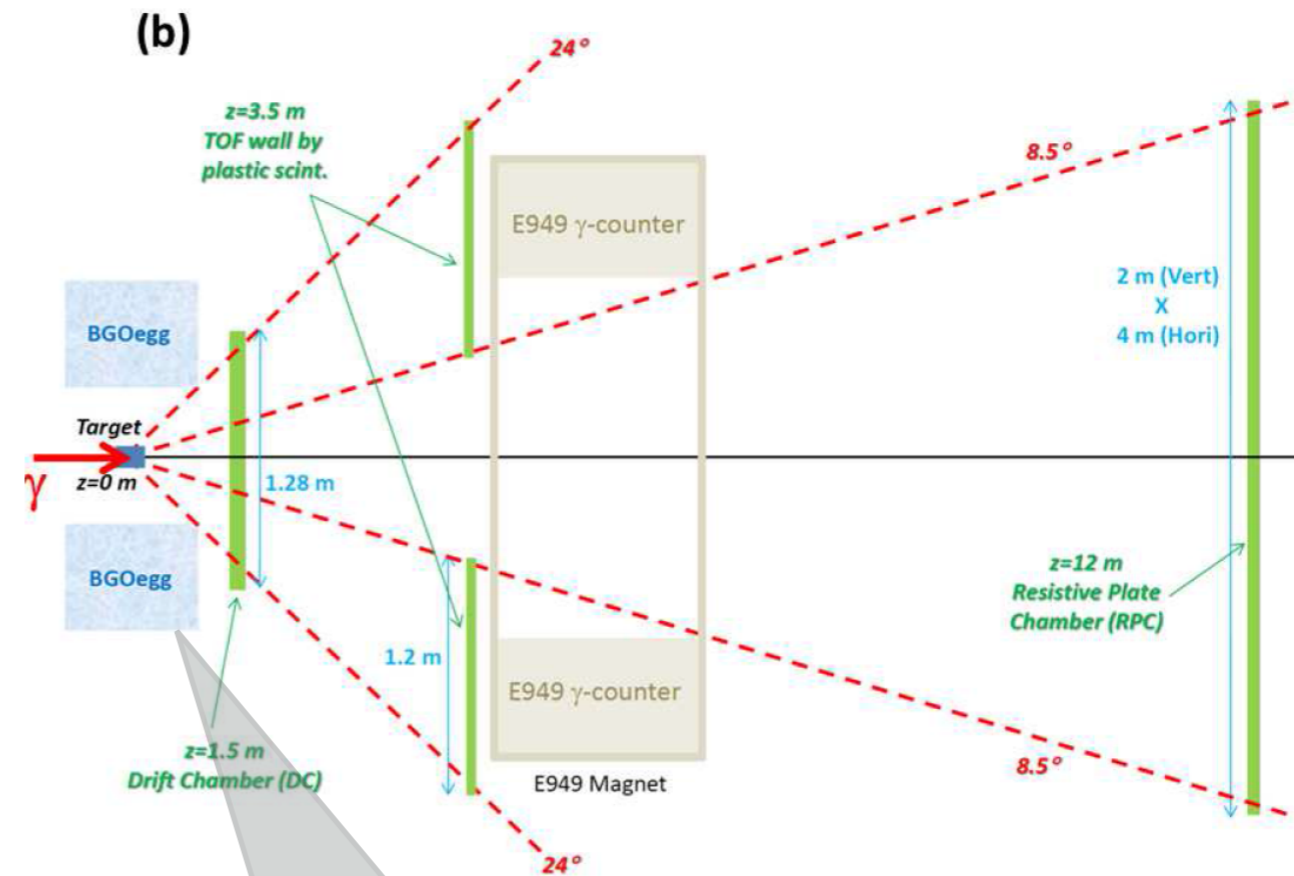
$^{12}\text{C}(\gamma, p) \eta' X @ 1.5\text{-}2.8 \text{ GeV}$



formation and decay of η' -mesic state

LEPS2@SPRing-8

$^{12}\text{C}(\gamma, p) \eta' X @ 1.5\text{-}2.4 \text{ GeV}$

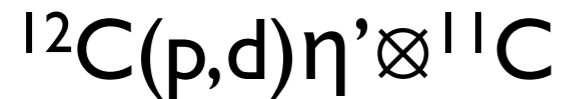


BGO-OD ideally suited for exclusive measurement

approved proposal: ELSA/3-2012-BGO

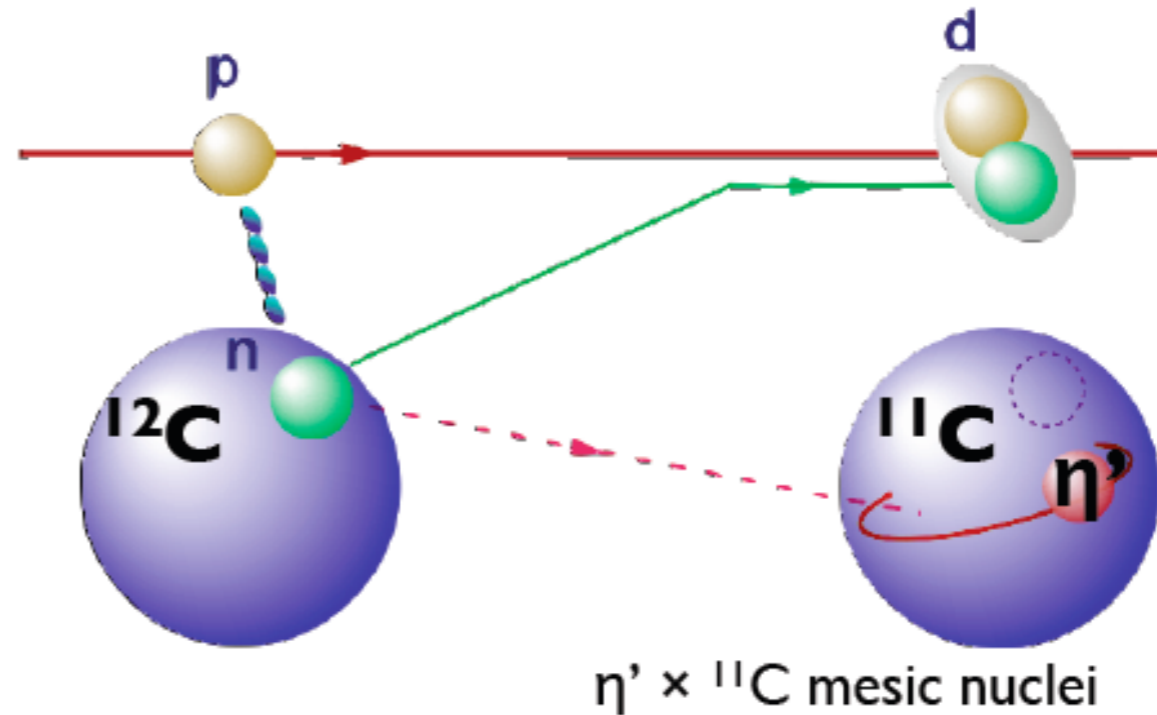
search for η' -mesic states in hadronic reactions

FRS@GSI: PRIME

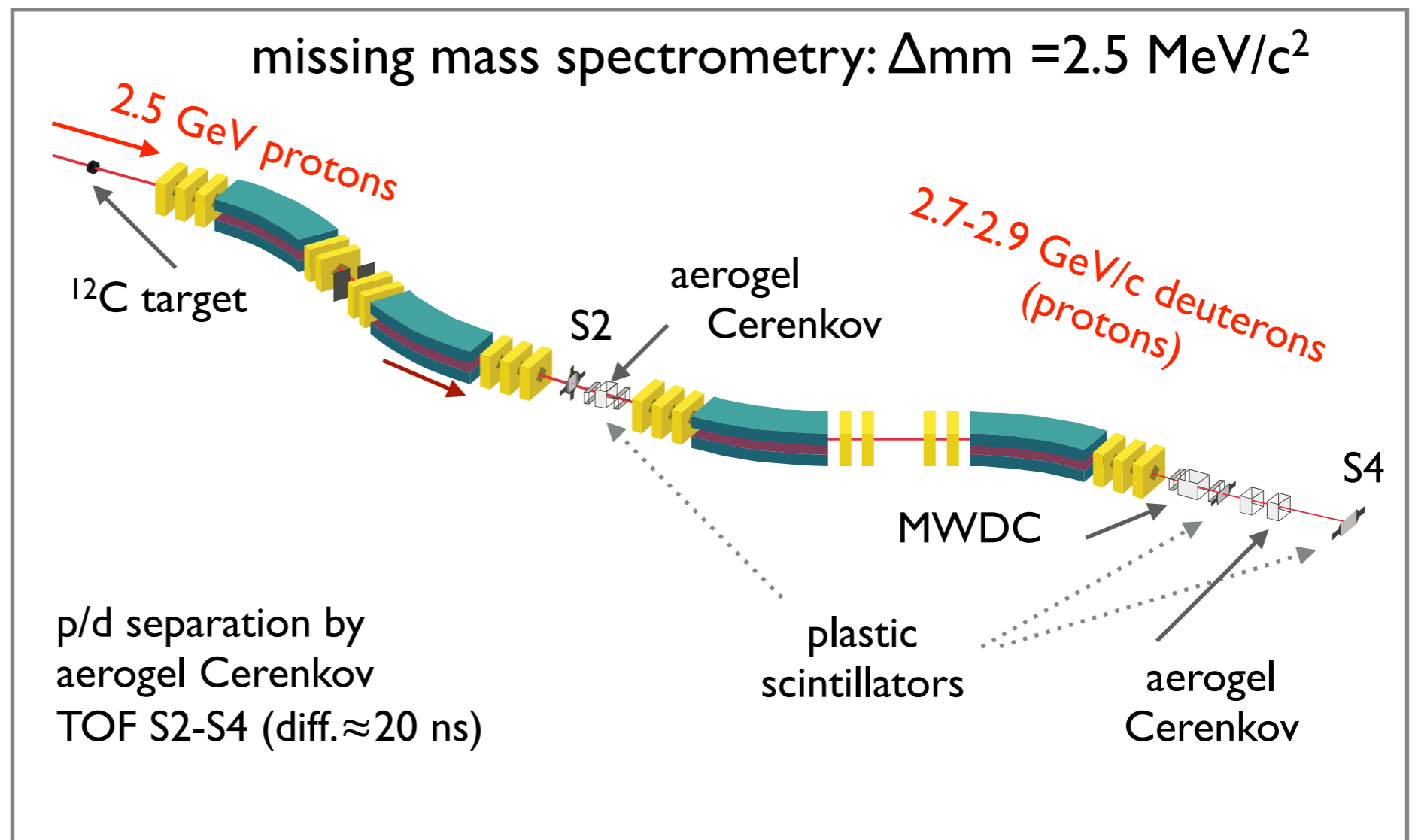


K. Itahashi et al., PETP 128 (2012) 601

H. Nagahiro et al., PRC 87 (2013) 045201



particle identification
by time-of-flight



Summary & Outlook

how do the hadron properties (mass, width) change
in a dense nuclear medium ??

meson properties do change in a strongly interacting medium !!

- ◆ all mesons are broadened; their lifetime is shortened through inelastic collisions
 $\Gamma_{\omega}(\rho=\rho_0; p=0) \approx 60 \text{ MeV}$; $\Gamma_{\eta'}(\rho=\rho_0; p=0) \approx 15 \text{ MeV}$;
- ◆ large mass modifications $|\Delta m| > 100 \text{ MeV}$ (as predicted by some calculations) have not been observed
- ◆ for the η' meson an in-medium mass drop of $\Delta m(\rho=\rho_0) \approx -40 \text{ MeV}$ has been determined
 - ◆ in-medium effects described within meson-nucleus optical
- ◆ the η' meson is a good candidate for forming meson-nucleus bound states since $|\text{Im } U| \ll |\text{Re } U|$
- ◆ search for η' mesic states ongoing

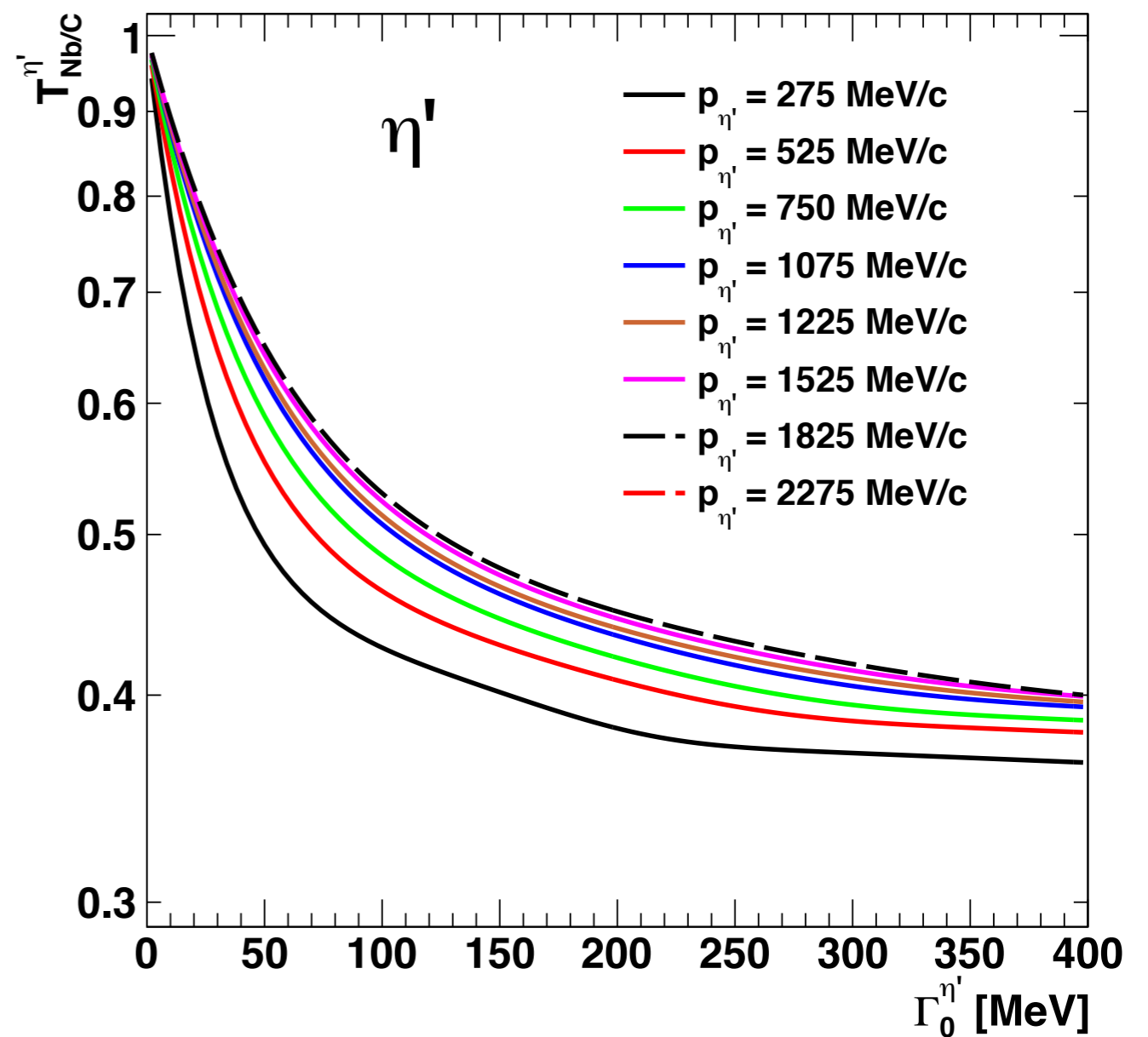
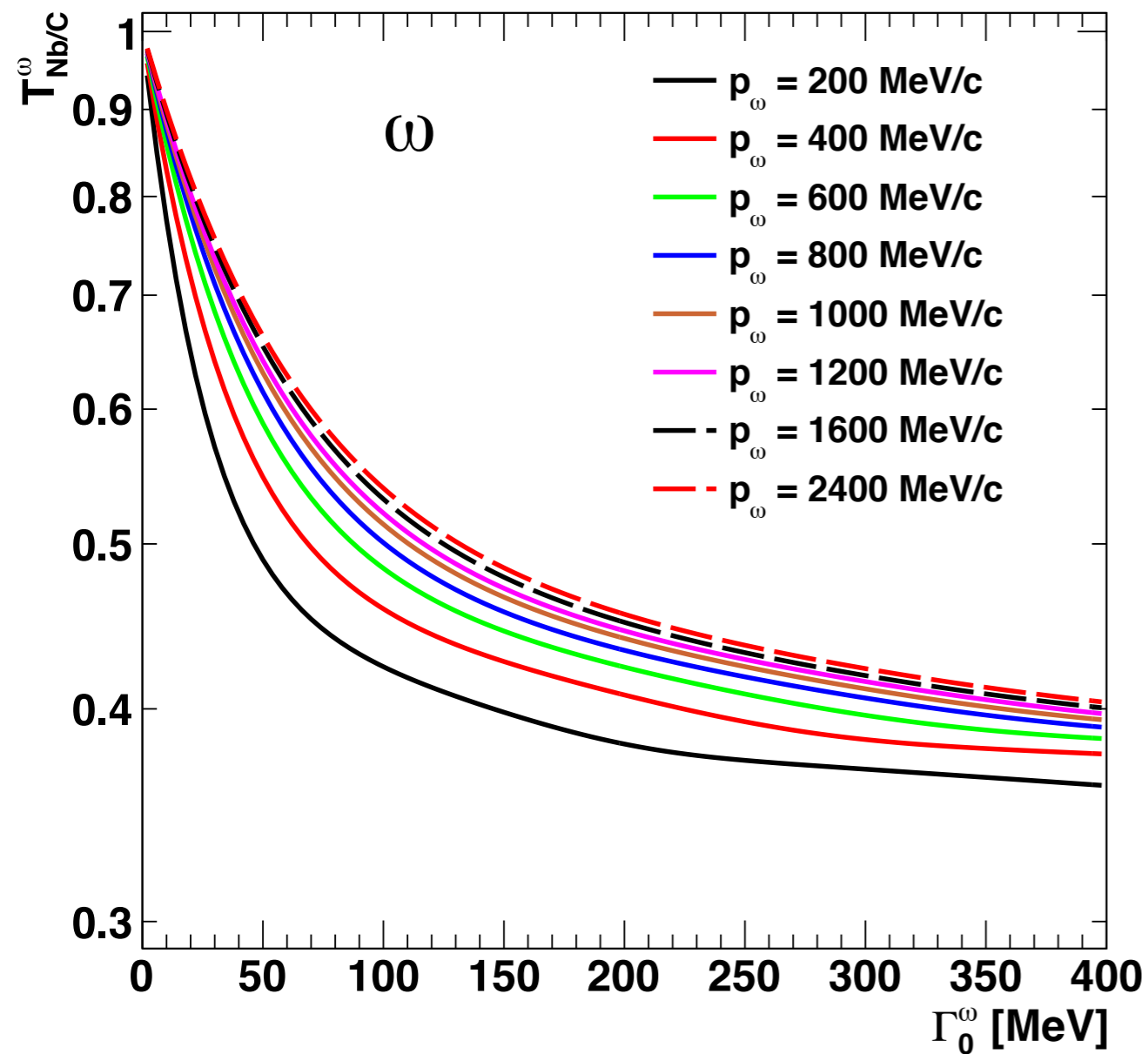
BACKUP

in-medium width from transparency ratio

Glauber model

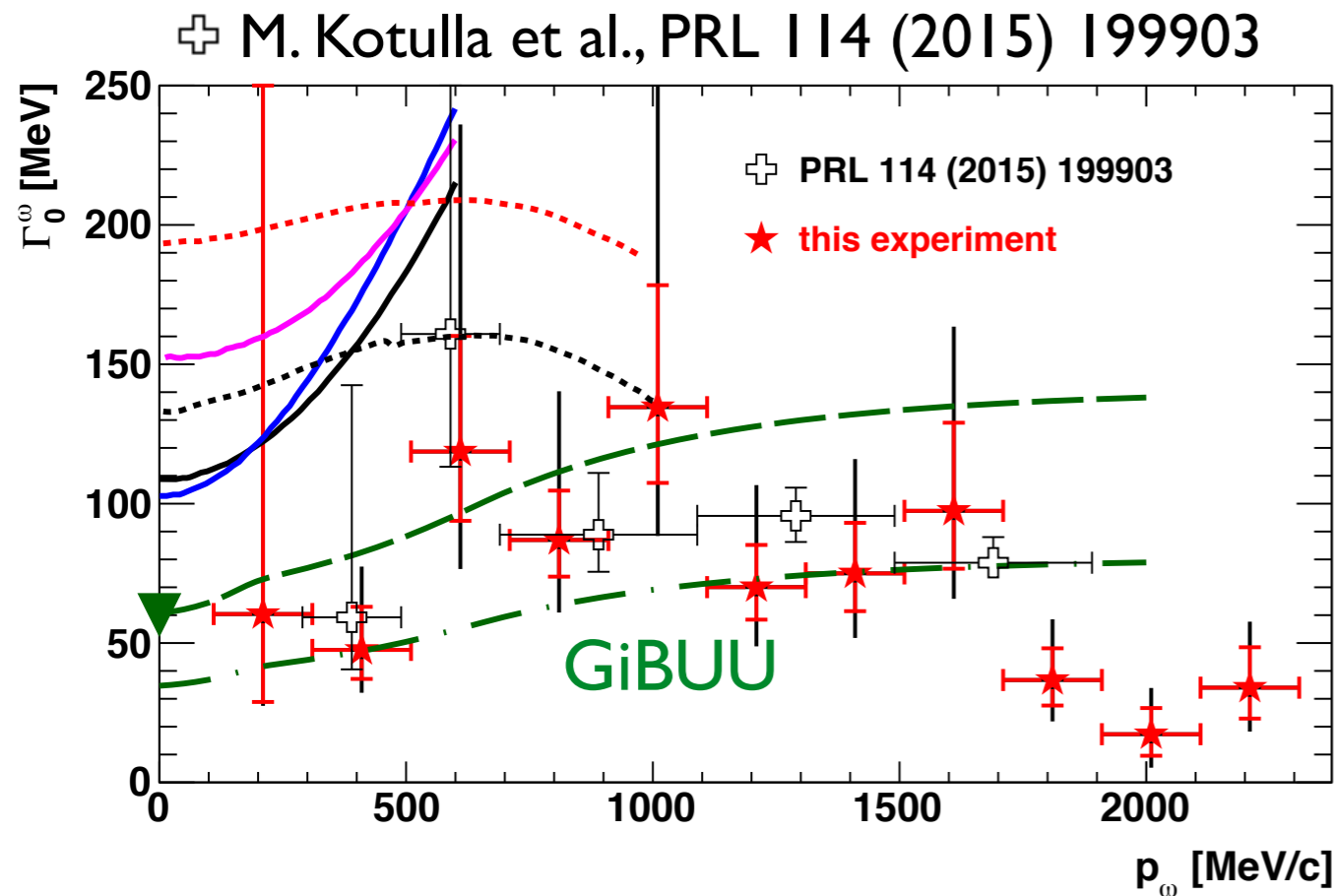
in high energy eikonal approximation

$$T_{\text{Nb/C}}(p) \longleftrightarrow \Gamma_0(p)$$



momentum dependence of ω, η' in-medium width

S. Friedrich

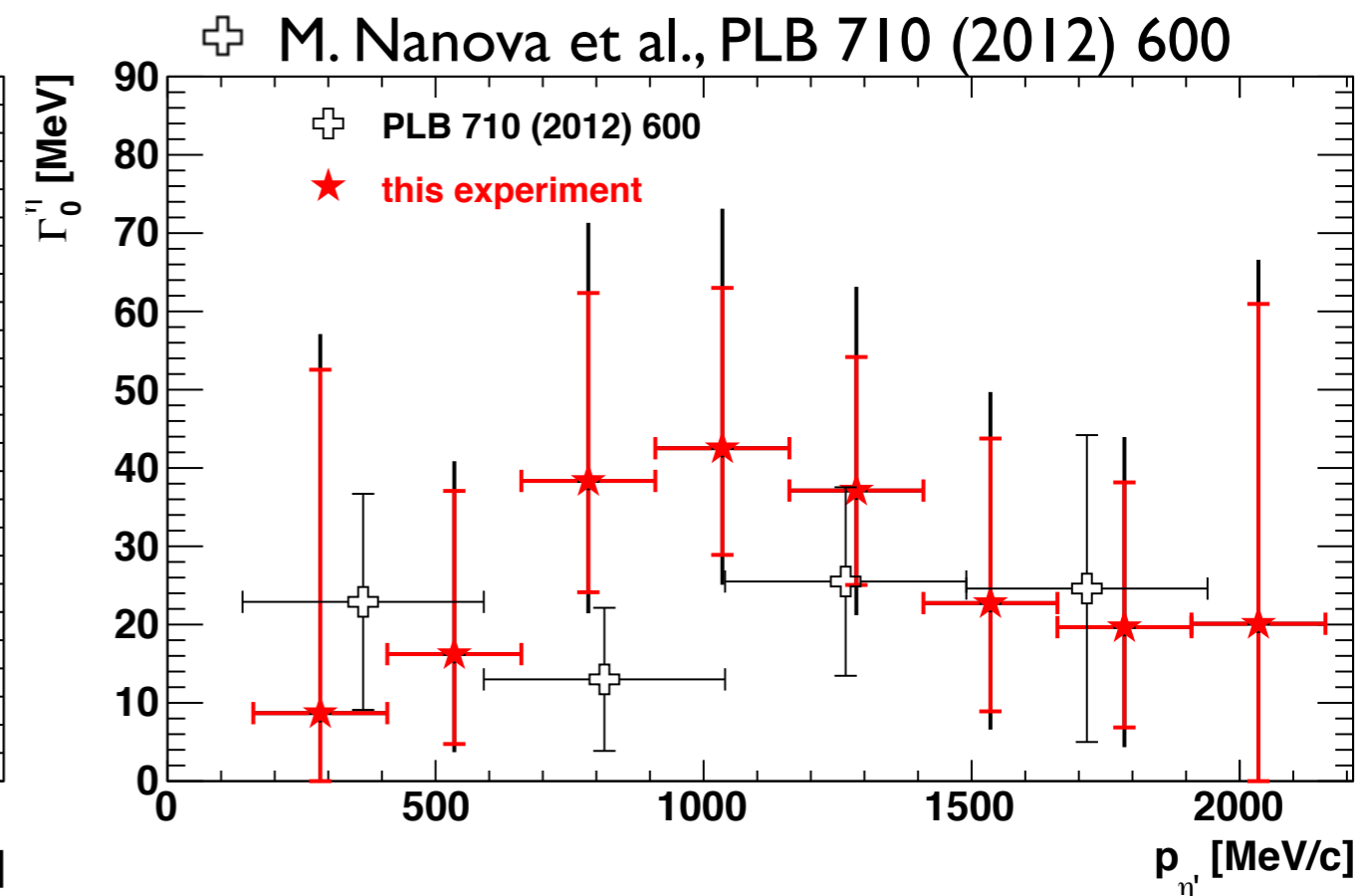


P. Mühlich et al., NPA 780 (2006) 187

O. Buss et al., Phys. Rep. 512 (2012) 1

A. Ramos et al., EPJA 49 (2013) 148

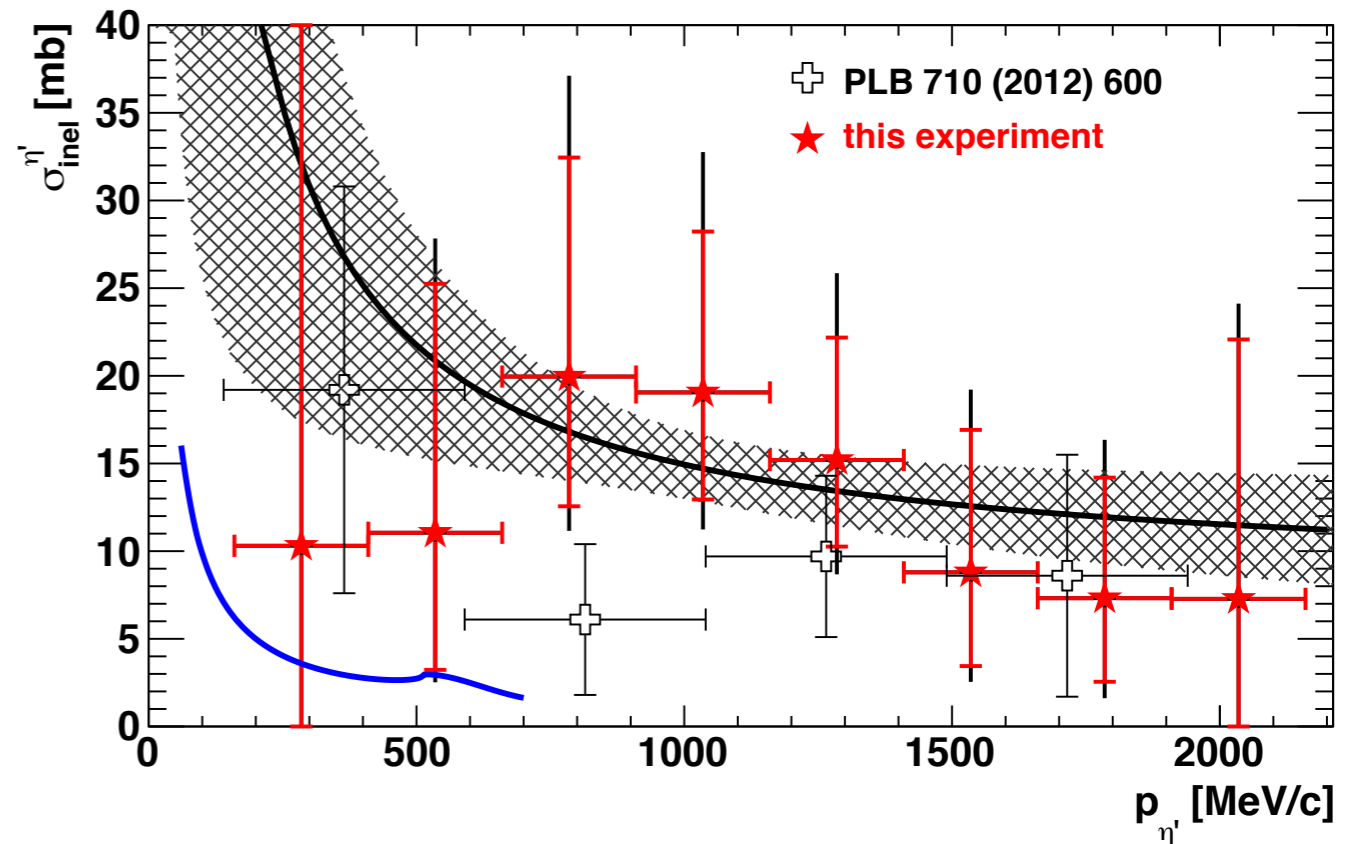
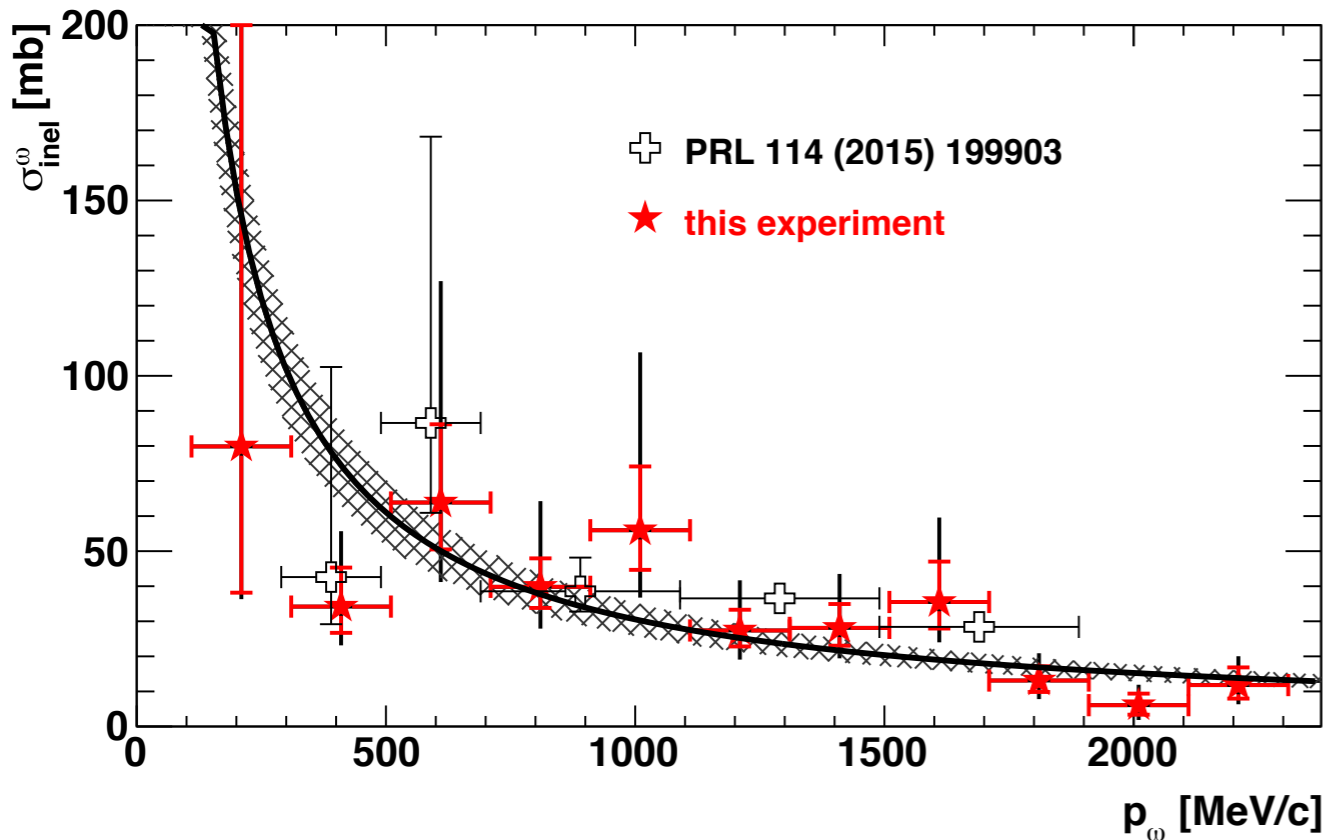
D. Cabrera and R. Rapp, PLB 729 (2014) 67



inelastic absorption cross section σ_{inel}

low density approximation

$$\Gamma(p) = \hbar c \cdot \beta \cdot \rho_0 \cdot \sigma_{inel}(p) \rightarrow \sigma_{inel}(p) = \frac{\Gamma(p)}{\hbar c \beta \rho_0}$$



E. Oset and A. Ramos, PLB 704 (2012) 334

$$\sigma_{inel}(p) [mb] = a + \frac{b}{p [GeV/c]}$$

$$a = 0.0 \pm 6.2$$

$$b = 31 \pm 4$$

$$a = 8.1 \pm 9.5$$

$$b = 6.8 \pm 9.8$$

$$\langle \sigma_{inel}(p) \rangle = (14 \pm 3) \text{ mb}$$

Dispersion relation analysis

work in progress (Horst Lenske (B7))

if self-energy Σ of the meson is an analytic function then imaginary and real part

related up to a constant by:

$$\text{Re } \Sigma(s) = -\frac{1}{\pi} P \int_0^{\infty} ds' \frac{\text{Im } \Sigma(s')}{s - s'} + \text{const}$$

