# Low Energy antikaon-nucleon/nuclei interaction studies by AMADEUS



Investigation of **in-medium modification of the KN interaction** fundamental for the low-energy QCD in the non perturbative regime.

**Chiral perturbation theory (ChPT)**: effective field theory where mesons and baryons represent the effective degrees of freedom instead of the fundamental quark and gluon fields.

$$\mathcal{L}_{eff} = \mathcal{L}_{mesons}(\Phi) + \mathcal{L}_B(\Phi, \Psi_B)$$

- chiral symmetry is spontaneously broken → existence of massless and spinless Nambu-Goldstone bosons which are identified with the pions (SU(2)). Explicitly broken by quark masses.
- very successful in describing the πN, ππ and NN interactions in the low-energy regime.

## Problematic extension of the theory to the s sector, not directly applicable to the $\overline{K}N$ channel.

ChPT not applicable to the KN channel due to the emerging of the  $\Lambda(1405)$  and the  $\Sigma(1385)$  resonances just below the KN mass threshold (~1432 MeV)



•  $\Lambda(1405)$  I=0 J<sup>P</sup> =  $\frac{1}{2}^{-1}$ M =  $(1405.1^{+1.3}_{-1.0})$  MeV  $\Gamma$  =  $(50.5 \pm 2.0)$  MeV decay modes:  $\Sigma \pi$  (I=0) 100%

•  $\Sigma(1385)$  I=1 JP = 3/2<sup>+</sup> decay modes:  $\Lambda \pi$  (I=1) (87.0 ± 1.5) %  $\Sigma \pi$  (I=1) (11.7 ± 1.5) %

**Possible solutions:** 

Non-perturbative Coupled Channels approach: Chiral Unitary SU(3) Dynamics

 $\succ$  Phenomenological  $\overline{K}N$  and NN potentials

The parameters of the models are constrained by the existing scattering data  $\rightarrow$  above the threshold

Phen. [Y. Ikeda and T. Sato, Phys. Rev. C76, 035203 (2007)]
 Chiral [S. Ohnishi, Y. Ikeda, T. Hyodo, W. Weise, Phys.Rev. C93 (2016) no.2, 025207]
 → aslo see the talk



The parameters of the models are constrained by the existing scattering data  $\rightarrow$  above the threshold

Chiral J. Révai [Few Body Systems 59(2018)49] → aslo see the talk



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...but... large differences in the subthreshold extrapolations! Significantly weaker attraction in chiral SU(3) models than in phenomenological potential models.



**Chiral unitary models:**  $\Lambda(1405)$  is an I = 0 quasibound state emerging from the coupling between the  $\overline{KN}$  and the  $\Sigma\pi$  channels. Two poles in the neighborhood of the  $\Lambda(1405)$ :



Chiral dynamics predicts significantly weaker attraction than AY (local, energy independent) potential in far-subthreshold region

**Chiral unitary models:**  $\Lambda(1405)$  is an I = 0 quasibound state emerging from the coupling between the  $\overline{KN}$  and the  $\Sigma\pi$  channels. Two poles in the neighborhood of the  $\Lambda(1405)$ :

*two poles*: about 1420 ; about = 1380 MeV Phys. Lett. B 500 (2001), Phys. Rev. C 66 (2002), (Nucl. Phys. A 725(2003) 181) ... many others .. (Nucl. Phys. A881, 98 (2012)) ... others

mainly coupled to KN

mainly coupled to  $\Sigma \pi \rightarrow$  line-shape depends on production mechanism

• J. Révai [Few Body Systems 59(2018)49]: solving the LS equation with full WT potential  $\rightarrow$  low mass pole in the KN -  $\Sigma\pi$  system disappears

#### Pole positions (MeV):

	Ζ <sub>1</sub>	<b>Z</b> <sub>2</sub>
U	1428 – 35 <i>i</i>	1384 – 62 <i>i</i>
V	1425 – 21 <i>i</i>	_

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BUBBLE CHAMBER search of the  $\Lambda(1405)$ :

- O. Braun et al. Nucl. Phys. B129 (1977) 1

K- induced reactions on d  $\rightarrow \Sigma^{-}\pi^{+}n$  the resonance is found & 1420 MeV

- D. W. Thomas et al., Nucl. Phys. B56 (1973) 15 pion induced reaction  $\pi$ - p  $\rightarrow$  K+ $\pi \Sigma$  the resonance is found & 1405 MeV

- R. J. Hemingway, Nucl. Phys. B253 (1985) 742  $K^-p \rightarrow \pi^-\Sigma^+(1660) \rightarrow \pi^-(\pi^+\Lambda(1405)) \rightarrow \pi^-\pi^+(\pi\Sigma) \& 4.2 \text{ GeV}$ analysed by Dalitz and Deloff  $M = 1406.5 \pm 4.0 \text{ MeV}, \ \Gamma = 50 \pm 2\text{MeV}$ 

- HADES coll. Phys. Rev. C 87, 025201 (2013)

 $pp \rightarrow p K^+ \pi \Sigma$  the resonance is found & 1390 MeV

THE "LINE-SHAPE" OF THE Λ(1405) DEPENDS ON THE OBSERVED CHANNEL !!

$$\frac{d\sigma(\Sigma^{-}\pi^{+})}{dM} \propto \frac{1}{3} |T^{0}|^{2} + \frac{1}{2} |T^{1}|^{2} + \frac{2}{\sqrt{6}} Re(T^{0}T^{1*})$$
$$\frac{d\sigma(\Sigma^{+}\pi^{-})}{dM} \propto \frac{1}{3} |T^{0}|^{2} + \frac{1}{2} |T^{1}|^{2} - \frac{2}{\sqrt{6}} Re(T^{0}T^{1*})$$
$$\frac{d\sigma(\Sigma^{0}\pi^{0})}{dM} \propto \frac{1}{3} |T^{0}|^{2}$$

THE "LINE-SHAPE" OF THE Λ(1405) DEPENDS ON THE OBSERVED CHANNEL !!

$$\begin{aligned} \frac{d\sigma(\Sigma^{-}\pi^{+})}{dM} &\propto \frac{1}{3} \left| T^{0} \right|^{2} + \frac{1}{2} \left| T^{1} \right|^{2} + \frac{2}{\sqrt{6}} Re(T^{0}T^{1*}) \\ \frac{d\sigma(\Sigma^{+}\pi^{-})}{dM} &\propto \frac{1}{3} \left| T^{0} \right|^{2} + \frac{1}{2} \left| T^{1} \right|^{2} - \frac{2}{\sqrt{6}} Re(T^{0}T^{1*}) \\ \frac{d\sigma(\Sigma^{0}\pi^{0})}{dM} &\propto \frac{1}{3} \left| T^{0} \right|^{2} \end{aligned}$$

IS DIFFERENT IN  $\Sigma^+\pi^-$  VS  $\Sigma^-\pi^+$ 

DUE TO ISOSPIN INTERFERENCE

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IS DIFFERENT IN  $\Sigma^+\pi^-$  VS  $\Sigma^-\pi^+$ 

DUE TO ISOSPIN INTERFERENCE

THE CLEANEST SIGNATURE OF THE Λ(1405) IS GIVEN BY THE NEUTRAL CHANNEL:

- is free from isospin interference
- is purely I = 0, no  $\Sigma(1385)$  contamination.

### $\Lambda(1405)$ .. the golden channel

Crystall Ball: K-p  $\rightarrow \Sigma^0 \pi^0 \pi^0$  for kaon momentum in the range (514-750 MeV/c). S. Prakhov et al. Phys Rev. C70 (2004) 03465 (interpreted by Magas et al. PRL 95, 052301 (2005))



COSY julich:  $pp \rightarrow pK^+ \Sigma^0 \pi^0$ 

(I. Zychor et al., Phys. Lett. B 660 (2008) 167)



CLAS:  $\gamma p \rightarrow K^+ \Sigma \pi$ 

AIP Conf.Proc. 1441 (2012) 296-298



Fig. 4. a) Missing–mass  $MM(p_{Fd}K^+)$  distribution for the  $pp \to pK^+p\pi^-X^0$  reaction for events with  $M(p_{Sd}\pi^-) \approx m(\Lambda)$  and  $MM(pK^+p\pi^-) > 190 \,\mathrm{MeV/c^2}$ . Exper-

Two main biases:

- the kinematical energy threshold 1412 MeV •  $(M_{K} + M_{p} - |BE_{p}|)$  the high pole energy region is closed,
- The shape and the amplitude of the NON-RESONANT •  $\Sigma \pi$  production below KbarN threshold is unknown.





Fig. 6. Detailed differences in  $M_{\Sigma\pi}$  spectra among the Hyodo–Weise prediction and the present model predictions

#### An ideal experiment:

- $\Lambda(1405)$  is produced in K- p absorption  $\rightarrow$  mainly coupled to the high mass pole, •
- $\Lambda(1405)$  is observed in the  $\Sigma^0\pi^0$  decay channel (pure isospin 0), ٠
- K- is absorbed in-flight on a bound proton with  $p_{K} \sim 100$  MeV,  $\Sigma \pi$  invariant mass • gain of ~ 10 MeV to open an energy window to the high mass pole.
- Knowledge of the  $\Sigma\pi$  NON-RESONANT production amplitude ... a choice of the • resonant amplitude necessary to model simulations

## How deep can an antikaon be bound in a nucleus?



#### **Possible Bound States:**

$(K^{-}pp) \rightarrow \Lambda p$	$(K^{-}ppn) \rightarrow \Lambda d$
$\rightarrow \Sigma^0 p$	$\rightarrow \Sigma^0 d$

predicted <u>if</u> strong KN interaction in the I=0 channel. [Wycech (1986) - Akaishi & Yamazaki (2002)]

#### K<sup>-</sup>pp bound state

....at the end of 2015

	BE (MeV)	$\Gamma (MeV)$	Reference
Dote, Hyodo, Weise	17-23	40-70	Phys.Rev.C79 (2009) 014003
Akaishi, Yamazaki	48	61	Phys.Rev.C65 (2002) 044005
Barnea, Gal, Liverts	16	41	Phys.Lett.B712 (2012) 132-137
Ikeda, Sato	60 - 95	45-80	Phys.Rev.C76 (2007) 035203
Ikeda, Kamano, Sato	9-16	34-46	Prog.Theor.Phys. (2010) 124(3): 533
Shevchenko, Gal, Mares	55 - 70	90-110	Phys.Rev.Lett.98 (2007) 082301
Revai, Shevchenko	32	49	Phys.Rev.C90 (2014) no.3, $034004$
Maeda, Akaishi, Yamazaki	51.5	61	Proc.Jpn.Acad.B 89, (2013) 418
Bicudo	14.2-53	13.8 - 28.3	Phys.Rev.D76 (2007) 031502
Bayar, Oset	15 - 30	75-80	Nucl.Phys.A914 (2013) 349
Wycech, Green	40-80	40-85	Phys.Rev.C79 (2009) 014001

Experiments reporting DBKNS				
KEK-PS E549	T. Suzuki at al. MPLA23, 2520-2523 (2008)			
FINUDA	M. Agnello et al. PRL94, 212303 (2005)	Extraction of a signal		
DISTO	T. Yamazaki et al. PRL104 (2010)	Extraction of a signal		
OBELIX	G. Bendiscioli et al. NPA789, 222 (2007)	Extraction of a signal		
HADES	G. Agakishiev et al. PLB742, 242-248 (2015)	Upper limit		
LEPS/SPring-8	A.O. Tokiyasu et al. PLB728, 616-621 (2014)	Upper limit		
J-PARC E15	T. Hashimoto et al. PTEP, 061D01 (2015)	Upper limit		
J-PARC E27	Y. Ichikawa et al. PTEP, 021D01 (2015)	Extraction of a signal		

## How deep can an antikaon be bound in a nucleus?

#### interpreted in

T. Sekihara, E. Oset, A. Ramos, Prog. Theor. Exp. Phys (2016) (12): 123D03



[from the talk of T. Nagae at HYP2015, Sep. 10, 2015]

## **Bound state search in K- induced reactions**

E549 at KEK:  $K_{stop}^{-} + {}^{4}He \rightarrow \Lambda + p + X'$ 

#### detected particles



Measurement of yields and shapes of the K- multinucleon yields is mandatory to solve the puzzle!

They are the counterpart of the non-resonant single nucleon capture

- 1NA: K<sup>-</sup> single nucleon absorption
- 2NA: K<sup>-</sup> two nucleon absorption
- 2NA + conversion, multi-nucleon, or Bound State?

## and K- multi-nucleon cross section?



Transport models and collision calculations need the measurement of the Kmulti-nucleon cross sections at low energy

In medium K properties investigated in heavy-ion & proton nuclei collisions, K<sup>-</sup> mass modification extrapolated from the K<sup>-</sup> production yield

still missing!

. . .

## **AMADEUS** scientific case

- Nature of the  $\Lambda(1405)$  & K<sup>-</sup>N amplitude below threshold  $\rightarrow$  **Y** $\pi$  **CORRELATION STUDIES**
- K<sup>-</sup> multi-nucleons absorptions cross sections
- kaonic nuclear clusters
  - $\rightarrow$  **YN CORRELATION STUDIES** (i.e.  $\Lambda p$ ,  $\Sigma^0 p$ , and  $\Lambda t$  final states)
  - Low-energy charged kaon cross sections for low momenta (100 MeV/c)
  - YN scattering → extremely poor experimental information from scattering data (strong impact on the EoS of Neutron Stars Related to NS merging radiation + GW emission)

## **AMADEUS & DAΦNE**



• **back to back** K<sup>+</sup>K<sup>-</sup> topology



**AMADEUS step 0**  $\rightarrow$  KLOE 2004-2005 dataset analysis ( $\mathscr{L} = 1.74 \text{ pb}^{-1}$ )



#### KLOE → see talk by D. KISIELEWSKA

• Cilindrical drift chamber with a  $4\pi$  geometry and electromagnetic calorimeter

#### 96% acceptance

- optimized in the energy range of all **charged particles** involved
- good performance in detecting photons and neutrons checked by kloNe group [M. Anelli et al., Nucl Inst. Meth. A 581, 368 (2007)]<sub>20</sub>

## K<sup>-</sup> absorption on light nuclei

from the materials of the KLOE detector DC gas (90% He, 10%  $C_4H_{10}$ ) & DC wall (C + H) **AT-REST** (K<sup>-</sup> absorbed from atomic orbit) or **IN-FLIGHT**  $(p_{v} \sim 100 MeV)$ YOKE S.C. COIL Cryostat **Barrel** calorimeter Advantage: DRIFT CHAMBER excellent resolution ..  $\sigma_{pA} = 0.49 \pm 0.01$  MeV/c in DC gas  $\sigma_{mvv} = 18.3 \pm 0.6 \text{ MeV/c}^2$ 7 m **Disadvantage:** Not dedicated target → different nuclei **contamination** → **complex interpretation** .. but  $\rightarrow$  new features ... K<sup>-</sup> in flight absorption. 21

6 m

## At-rest VS in-flight K<sup>-</sup> captures

AT-REST K<sup>-</sup> absorbed from atomic orbit (p<sub>K</sub>~ 0 MeV)



### <u>IN-FLIGHT</u> (p<sub>к</sub>~100MeV)



## Pure graphite Carbon target

#### Advantages:

- gain in statistics
- pure K<sup>-</sup> Carbon absorptions
  - pure absorptions at-rest.



- MC simulation: 26% of K<sup>-</sup> stopped in <sup>12</sup>C
- Thickness optimized to maximize the number of stopping K<sup>-</sup> in the targed (minimizing energy loss)

(~90 pb<sup>-1</sup>; analyzed 37 pb<sup>-1</sup>, x 1.5 statistics)

## K<sup>-</sup> - N single nucleon absorption

## resonant and non-resonant amplitudes

#### $\Lambda(1405)$ case





FIG. 4: Theoretical  $(\pi^0 \Sigma^0)$  invariant mass distribution for an finitial kaon lab momenta of 687 MeV. The non-symmetrized text distribution also contains the factor 1/2 in the cross section.

FIG. 5: Two experimental shapes of  $\Lambda(1405)$  resonance. See text for more details.





#### $\Lambda(1405)$ : extracting the resonant I = 0 contribution

PID optimised, data fit is ongoing

necessary the input of the  $\Lambda\pi^-$  measurement

- K. Piscicchia et al., APP B48 (2017) 10, 1875
- C. Curceanu, K. Piscicchia et al., APP B46 (2015) 1, 203



## **Resonant VS non-resonant**

 $K^{-}N \rightarrow (Y^{*}?) \rightarrow Y\pi$ how much comes from resonance?

Non resonant transition amplitude never measured before below threshold can be obtained exploiting K<sup>-</sup> N in-medium absorption,

chosen target <sup>4</sup>He

- K<sup>-</sup> angular momentum at the capture
   absorbing nucleon wave function

known quantities

## **Resonant VS non-resonant**

Investigated using:  $\mathbf{K}^{-} \mathbf{n}^{-} \rightarrow \mathbf{\Lambda} \pi^{-}$  direct formation in <sup>4</sup>He

> the goal is to measure  $|f^{N-R}|_{\Lambda\pi}$  (I=1) | to get information on  $|f^{N-R}|_{\Sigma\pi}$  (I=0) |



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## Resonant VS non-resonant

## Investigated using: $K^{-}''n'' \rightarrow \Lambda \pi^{-}$ to extract $|f^{N-R}_{\Lambda \pi}(I=1)|$ below threshold

J. Hrtankova, and J. Mares, Phys. Rev. C96, (2017) 015205 A. Cieply et al., Nucl. Phys. A954, (2016) 17



K<sup>- 4</sup>He → Λp<sup>- 3</sup>He resonant and non-resonant processes K. P., S. Wycech and C. Curceanu, Nucl. Phys. A954 (2016) 75-93 R. Del Grande, K. P., S. Wycech, Acta Phys. Pol. B 48 (2017) 1881



**Theoretical shapes for :** 

total  $\Lambda\pi^{-}$  momentum spectra for the resonant ( $\Sigma^{*-}$ ) and non-resonant (I = 1) processes were calculated, for both S-state and P-state K<sup>-</sup> capture at-rest and in-flight. Corrections to the amplitudes due to  $\Lambda/\pi$  final state interactions were estimated.



## Simultaneous fit : $(p_{\Lambda\pi-} - m_{\Lambda\pi-} - \cos(\theta_{\Lambda\pi-}))$



## Comparison



 $m_{\Lambda\pi}$  fit Light band sys err. Dark band stat. Err.

### **Outcome of the measurement**

From the well known  $\Sigma^*$  transition probability:

$$\frac{\mathrm{NR} - \mathrm{ar}}{\mathrm{RES} - \mathrm{ar}} = \frac{\int_{0}^{pmax} P_{ar}^{nr}(p_{\Lambda\pi}) dp_{\Lambda\pi}}{\int_{0}^{pmax} P_{ar}^{res}(p_{\Lambda\pi}) dp_{\Lambda\pi}} =$$

 $= |f_{ar}^s|^2 \cdot 8,94 \cdot 10^5 \text{MeV}^2.$ 

 $|f_{ar}^{nr}| = |A_{K-n \to \Lambda \pi^-}| = (0.334 \pm 0.018 \operatorname{stat}_{-0.058}^{+0.034} \operatorname{syst}) \operatorname{fm}$ 

compatible with K<sup>-</sup>  $p \rightarrow \Lambda \pi^0$  scattering above threshold

J. K. Kim, Columbia University Report, Nevis 149 (1966),

J. K. Kim, Phys Rev Lett, 19 (1977) 1074:

E = -33  MeV	$p_{lab} = 120 \text{ MeV}$	$160 { m MeV}$	$200 {\rm ~MeV}$	$245~{\rm MeV}$
$0.334 \pm 0.018  \mathrm{stat}^{+0.034}_{-0.058} \mathrm{syst}$	0.33(11)	0.29(10)	0.24(6)	0.28(2)

## **Outcome of the measurement**



To compare with theoretical calculations:

- 1) extract the amplitude for each model ..  $A_{K-n} = (\text{ReF}_{K-n}^2 + \text{ImF}_{K-n}^2)^{1/2}$
- 2) scale the amplitudes for the K<sup>-</sup>n couplings to the  $\Sigma^{-}\pi^{0}$  and  $\Sigma^{0}\pi^{-}$  channels:

$$\frac{Prob_{K^-n\to\Lambda\pi^-}}{Prob_{K^-n\to\Sigma^-\pi^0}} = \frac{Ph_{K^-n\to\Lambda\pi^-}}{c_1Ph_{K^-n\to\Sigma^-\pi^0}}$$

$$\underset{=(1,-1) \text{ component}}{\text{Isospin (I, I_2) = }}$$

$$\frac{Prob_{K^-n\to\Lambda\pi^-}}{Prob_{K^-n\to\Sigma^0\pi^-}} = \frac{Ph_{K^-n\to\Lambda\pi^-}}{c_2Ph_{K^-n\to\Sigma^0\pi^-}}$$
Phase spaces ratios

## **Outcome of the measurement**

 $|f_{ar}^{s}| = (0.334 \pm 0.018 \operatorname{stat}_{-0.058}^{+0.034} \operatorname{syst}) \operatorname{fm}.$ 



A<sub>K-n →  $\Lambda\pi$ </sub> (s<sup>1/2</sup> ~ 1400 MeV)<sup>1/2</sup>  $E_{Kn} = -|B_n| - \frac{p_3^2}{2\mu_{\pi,\Lambda,3He}}$ Nucl. Phys. A954 (2016) 75-93 Phys. Rev. C 96 (2017) 045204

Phys. Lett. B 702 (2011) 402–407

Nucl.Phys. A968 (2017) 35-47

K. P., S. Wycech, L. Fabbietti et al. Phys.Lett. B782 (2018) 339-345

## K<sup>-</sup> - multiN absorption and search for bound states

## AMADEUS contribution from low energy K<sup>-12</sup>C absorption $\Sigma^0 p / \Lambda p$ final states

O. Vazquez Doce et al, Phys Lett B 758, (2016) 134



From the contributions to the fit, the yields are extracted for K- stop

	yield / $K_{stop}^{-} \cdot 10^{-2}$	$\sigma_{stat} \cdot 10^{-2}$	$\sigma_{syst} \cdot 10^{-2}$
2NA-QF	0.127	$\pm 0.019$	$+0.004 \\ -0.008$
2NA-FSI	0.272	$\pm 0.028$	$^{+0.022}_{-0.023}$
Tot 2NA	0.376	$\pm 0.033$	$^{+0.023}_{-0.032}$
3NA	0.274	$\pm 0.069$	$^{+0.044}_{-0.021}$
Tot 3body	0.546	$\pm 0.074$	$+0.048 \\ -0.033$
4NA + bkg.	0.773	$\pm 0.053$	$^{+0.025}_{-0.076}$

- disappearance of the bound state in K<sup>-12</sup>C induced reaction explained

рк ~ 100 MeV/с

for

## K<sup>-</sup> - 4NA cross section & BR

### At available data

Available data:

• in Helium :

bubble chamber experiment
 [M.Roosen, J.H. Wickens, II Nuovo Cimento 66, (1981), 101]
 K<sup>-</sup> stopped in liquid helium, Λ dn/t search. 3 events compatible with the Λt kinematics were found

 $BR(K^{-4}He \rightarrow \Lambda t) = (3 \pm 2) \times 10^{-4}/K_{stop}$ 

global, no 4NA

Solid targets

- FINUDA [Phys.Lett. B669 (2008) 229] (40 events in different solid targets)

#### **∧t** available data

FINUDA presented [Phys.Lett.B (2008) 229]:

- a study of Λ vs t momentum correlation and an opening angle distribution
- 40 events collected and added together coming from different targets (<sup>6,7</sup>Li, <sup>9</sup>Be)



# $K^- {}^{4}He \rightarrow \Lambda t \ cross \ section, \ DC \ gas \ sample \ contributing \ processes:$

#### single nucleon absorption (1NA)

#### **Spectator tritons have low momentum:**





4NA processes – K<sup>-</sup> absorbed on FREE α:

- $K^{-4}He \rightarrow \Lambda t$
- $K^{\text{-}\,4}He \ \rightarrow \ \Sigma^0 t$  ,  $\Sigma^0 \ \rightarrow \ \Lambda \gamma$



# K<sup>-</sup> <sup>4</sup>He → At cross section, DC gas sample contributing processes:

Main background: K<sup>-</sup> absorption on <sup>12</sup>C (isobutane contamination)

4NA:  $K^{-12}C \rightarrow K^{-}(\alpha)^{\text{bound 8}}Be \rightarrow \Lambda/\Sigma^{0} t^{-8}Be$ 

7 MeV/c<sup>2</sup> lower invariant mass threshold respect to:

4NA: 
$$K^{-4}He \rightarrow K^{-}(\alpha)^{\text{free}} \rightarrow \Lambda/\Sigma^{0} t$$

+

all possible elastic/inelastic FSI processes with primary  $\Lambda/\Sigma$  formation



Measured K<sup>- 12</sup>C sample from K- captures in wall:



### K- <sup>4</sup>He $\rightarrow$ At 4NA fit







## $K^{-12}C \rightarrow \Lambda/\Sigma^0 t^8Be 4NA$ without FSI

**BR**(K<sup>-4</sup>He(4NA)  $\rightarrow$   $\wedge$ t) = **1.5 ± 0.5 × 10<sup>-4</sup> (stat)** /K<sub>stop</sub>

σ(K<sup>12</sup>C (4NA) → Λt<sup>8</sup>Be) = 0.58 ± 0.11 (stat) mb

σ(K<sup>12</sup>C (4NA) → Σ<sup>0</sup>t <sup>8</sup>Be) = **1.88 ± 0.35 (stat)** mb



## **Perspective:**

### **Measurement of the**

 $K^-H \rightarrow \Sigma^0 \pi^0$  cross section for  $p_{\kappa} = 97 \pm 10$  MeV/c

#### Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012)



### Low momentum $p_{\Sigma^+}$ structure in $\Sigma^+\pi^-$ formation



Fig. 5. Momentum distributions of sigmas from the  ${}^{6}Li(K_{stop}^{-}, \pi^{\pm}\Sigma^{\mp})A'$  reactions. The grey-filled histograms are the measured distributions. The distributions of Monte-Carlo generated sigmas are depicted by full dots, and with open diagrams are represented the M-C generated sigmas being reconstructed by FINUDA.

FINUDA coll. M. Agnello et al., Phys. Lett. B704 (2011) 474.  $K^{-6}Li \rightarrow \Sigma^{+}\pi^{-}A'$ 

### Low momentum $p_{\Sigma^+}$ structure in $\Sigma^+\pi^-$ formation

K. Piscicchia et al., EPJ Web Conf. 137 (2017) 09005.

 $K^{-9}Be \rightarrow \Sigma^{+}\pi^{-} + + \alpha + n + t$ 

no structure at low momentum

 $K^{-12}C \rightarrow \Sigma^+\pi^- A'$ 

structure at low momentum

amounts some % of the total yield

also in thiner targets

(not explained by energy loss)

Hypothesis: Σ<sup>+</sup> trapped in a Gamov state, interplay of the attractive nuclear potential & repulsive Coulomb barrier

S. Wycech, K. P., EPJ Web. Conf. 130 (2016) 02011

R. Del Grande, K. P. and S. Wycech, Acta Phys.Polon. B48 (2017) 1881

S. Wycech, K. P., On Gamov states of  $\Sigma^+$  hyperons, Acta Phys.Polon. B48 (2017) 1861

## Gamov state formation of a $\Sigma^+$ in light nuclei?

### ... work in progress

Gamov peak following in-flight capture

 $K^{-12}C \rightarrow \Sigma^+\pi^{-11}Be$ 

about 3% of the large peak Breit – Wigner - (Ε, Γ) = (1405,40) ; (1410,40) ; (1420,40)

Position p<sub>Σ+</sub> = 15 MeV/c peculiar structure due to the limitation of the phase space



![](_page_48_Figure_7.jpeg)

# ThanK<sup>-</sup> you