

Search for the Kaonic Bound State $K^{\text{bar}}NN$ at J-PARC

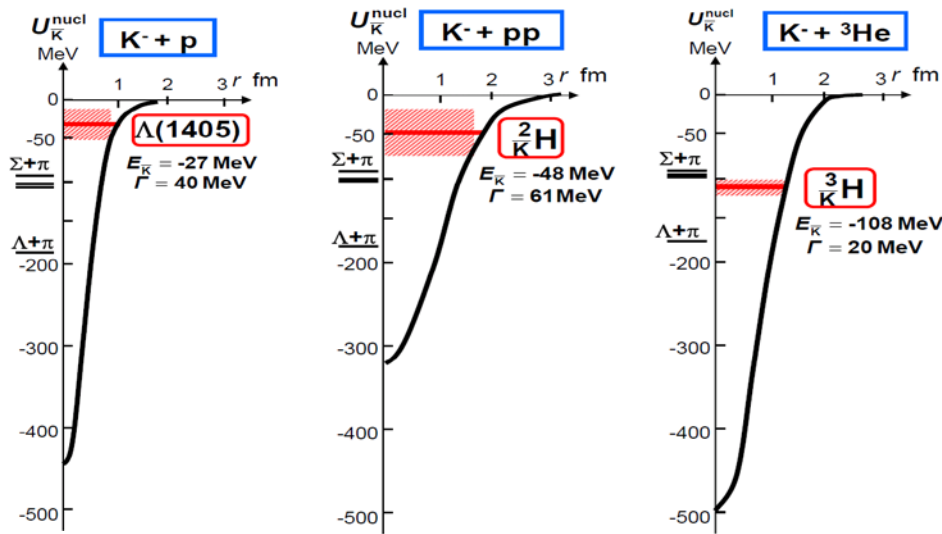
via Δp and $\pi\Sigma p$ decay channels

F. Sakuma, RIKEN 

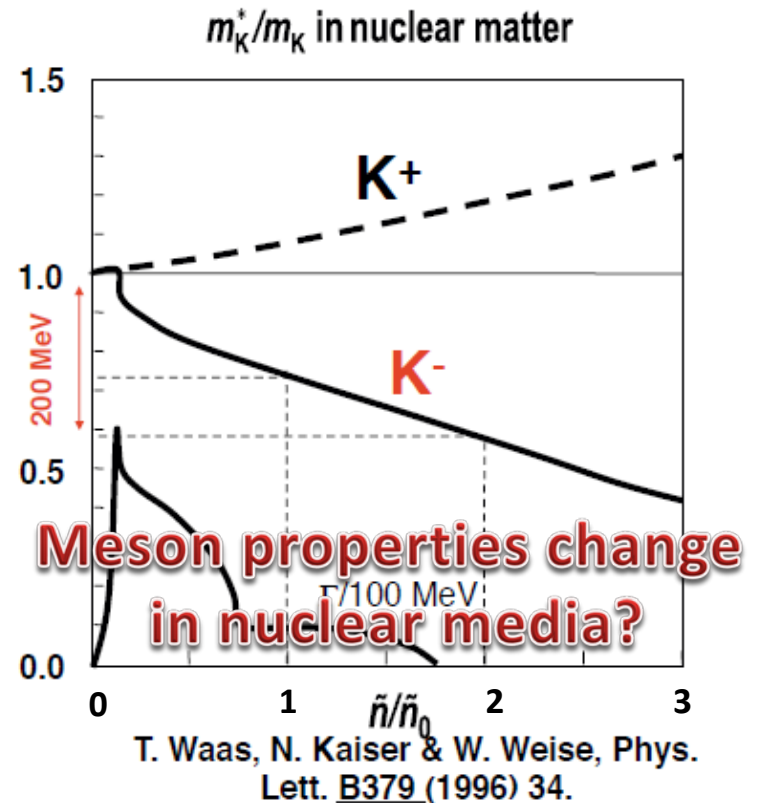
on behalf of the J-PARC E15
collaboration

Kaonic Nuclei

- Bound states of nucleus and anti-kaon
- Predicted as a consequence of **attractive $K^{\text{bar}}N$ interaction in $l=0$**



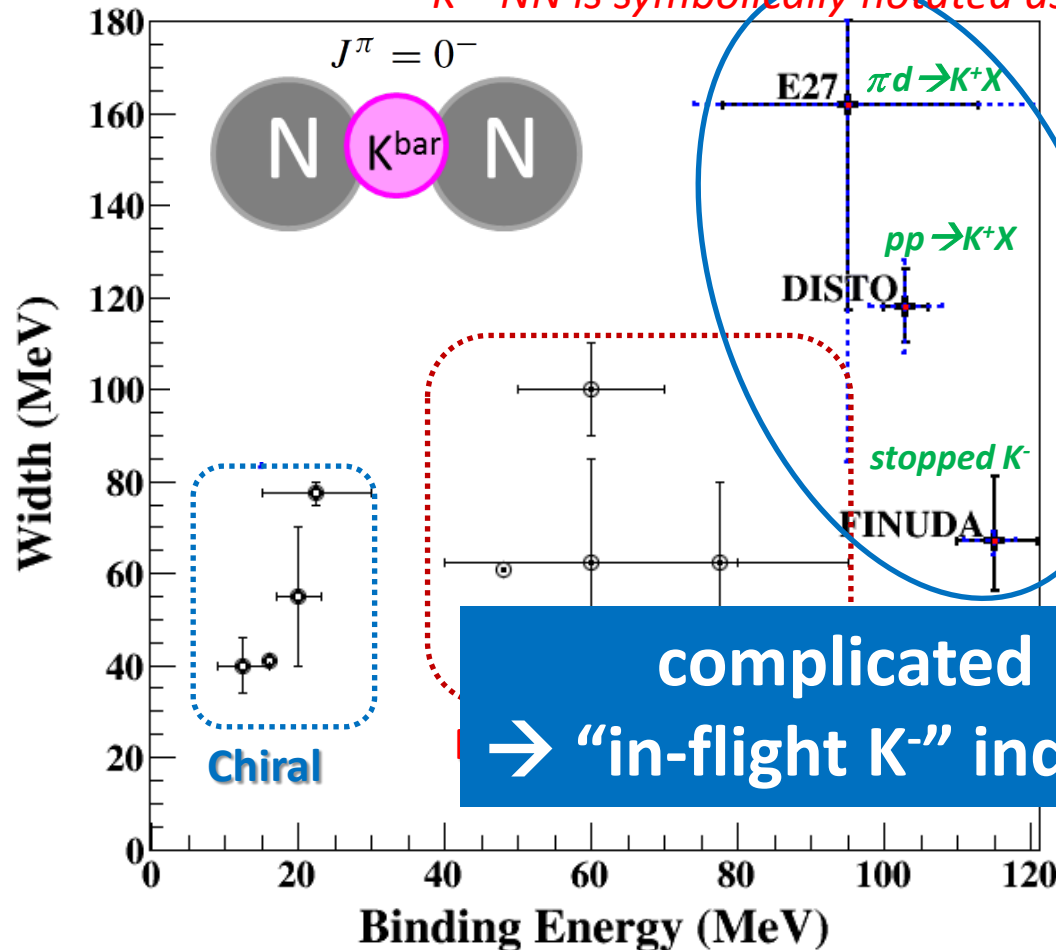
Y.Akaishi & T.Yamazaki, PLB535, 70(2002).



- Will provide new insight on **$K^{\text{bar}}N$ interaction in media**

Present Status of $K^{\text{bar}}\text{NN} = \text{“K-pp”}$

$K^{\text{bar}}\text{NN}$ is symbolically notated as “K-pp” in this talk



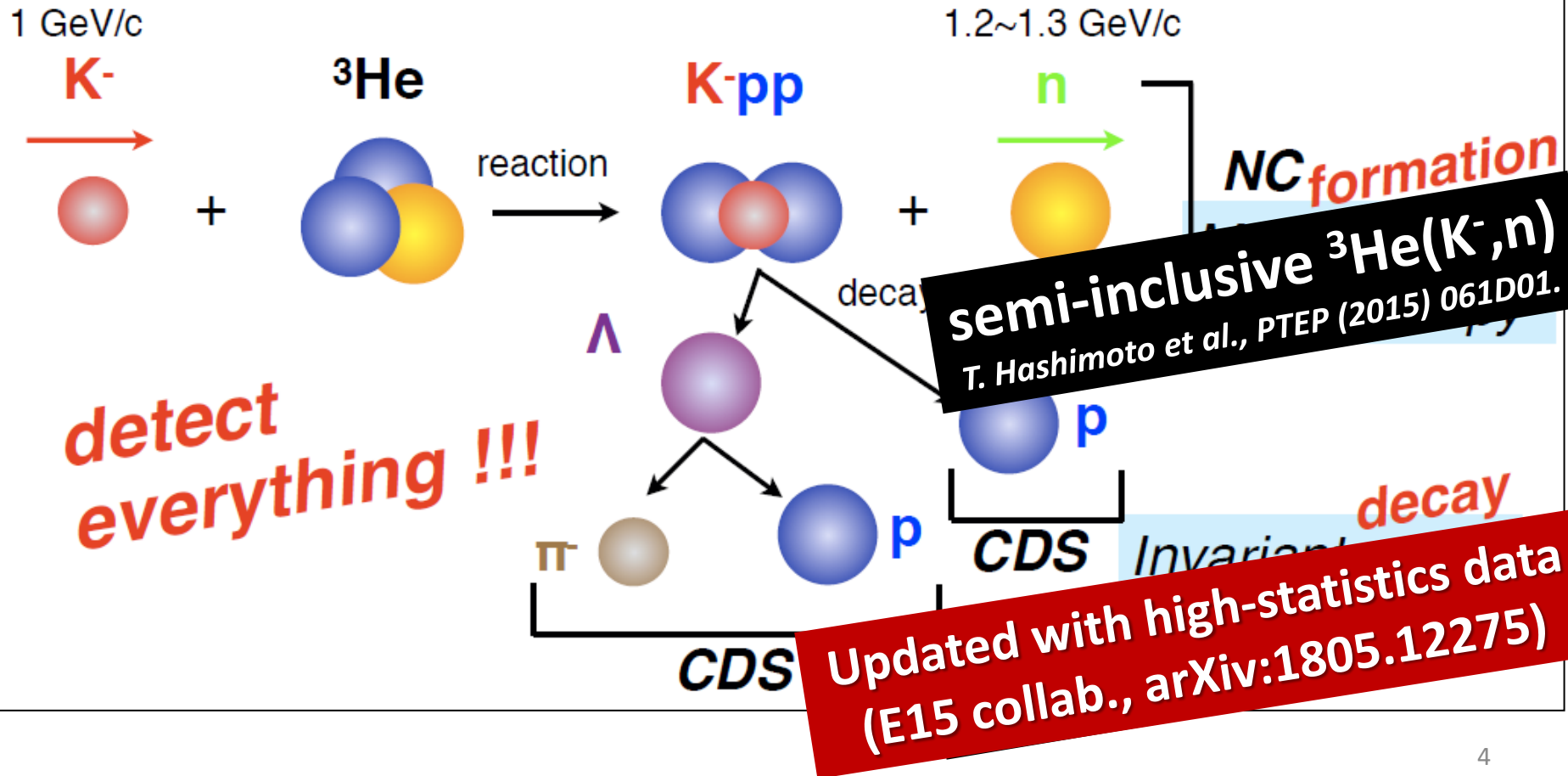
Upper limits were also obtained:

- LEPS@SPring8 [Inclusive $d(\gamma, K^+\pi^-)X$]
- HADES@GSI [Exclusive $pp \rightarrow p\Lambda K^+$]

J-PARC E15 Experiment

- ${}^3\text{He}(\text{in-flight } K^-, n)$ reaction @ 1.0 GeV/c

😊 2NA processes and Λ decays can be discriminated kinematically



“ K^-pp ”, a \bar{K} -Meson Nuclear Bound State, Observed in ${}^3\text{He}(K^-, \Lambda p)n$ Reactions

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¹⁸ Tohoku University, Sendai, 982-0826, Japan and

¹⁹ Lund University, Lund, 221 00, Sweden

(J-PARC E15 Collaboration)

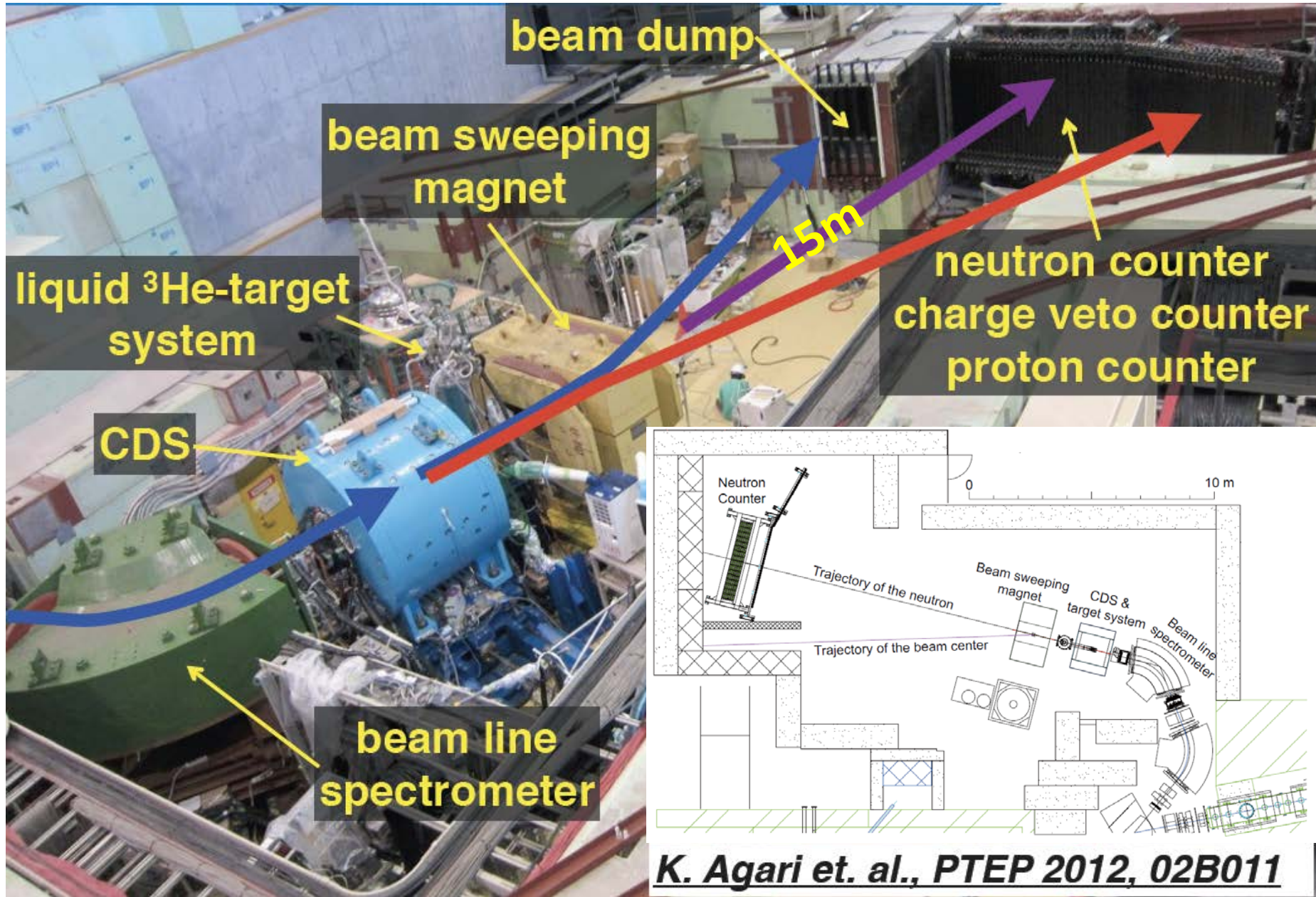
“ K^-pp ” in ${}^3\text{He}(K^-, \Lambda p)n$

We observed a distinct resonance peak in the Λp invariant-mass spectrum of ${}^3\text{He}(K^-, \Lambda p)n$, well below the mass threshold of $M(K^-pp)$. By selecting a relatively large momentum-transfer region $q = 350 \sim 650$ MeV/c, one can clearly separate the resonance peak from the quasi-free process, $\bar{K}N \rightarrow \bar{K}N$ followed by the non-resonant absorption by the two spectator-nucleons $\bar{K}NN \rightarrow \Lambda N$. We found that the simplest fit to the observed peak gives us a Breit-Wigner pole at $B_{Kpp} = 47 \pm 3$ (stat.) $_{-6}^{+3}$ (sys.) MeV having a width $\Gamma_{Kpp} = 115 \pm 7$ (stat.) $_{-9}^{+10}$ (sys.) MeV, and the S -wave Gaussian reaction form-factor parameter $Q_{Kpp} = 381 \pm 14$ (stat.) $_{-0}^{+57}$ (sys.) MeV/c, as a new form of the nuclear bound system with strangeness – “ K^-pp ”.

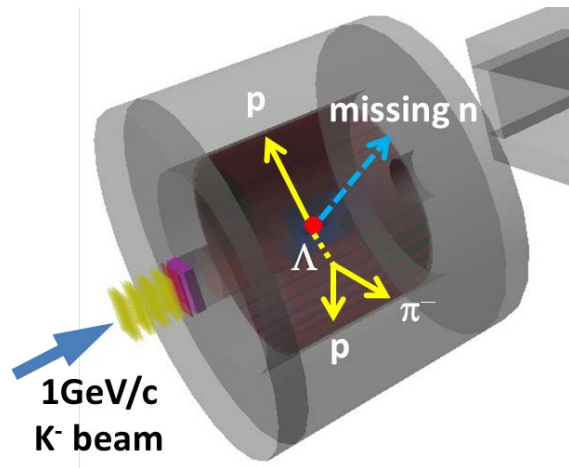
Since the prediction of the π -meson by Yukawa [1], there has been a long-standing question as to whether a mesonic nuclear bound state exists. Mesons are introduced as mediators between nucleons to confine them

in vacuum one needs energy m to produce them. If a mesonic nuclear bound state exists, it will form a quantum state at an energy E_M below m whose binding energy $B_M = m - E_M$. Many mesons have been examined

Experimental Setup @ K1.8BR

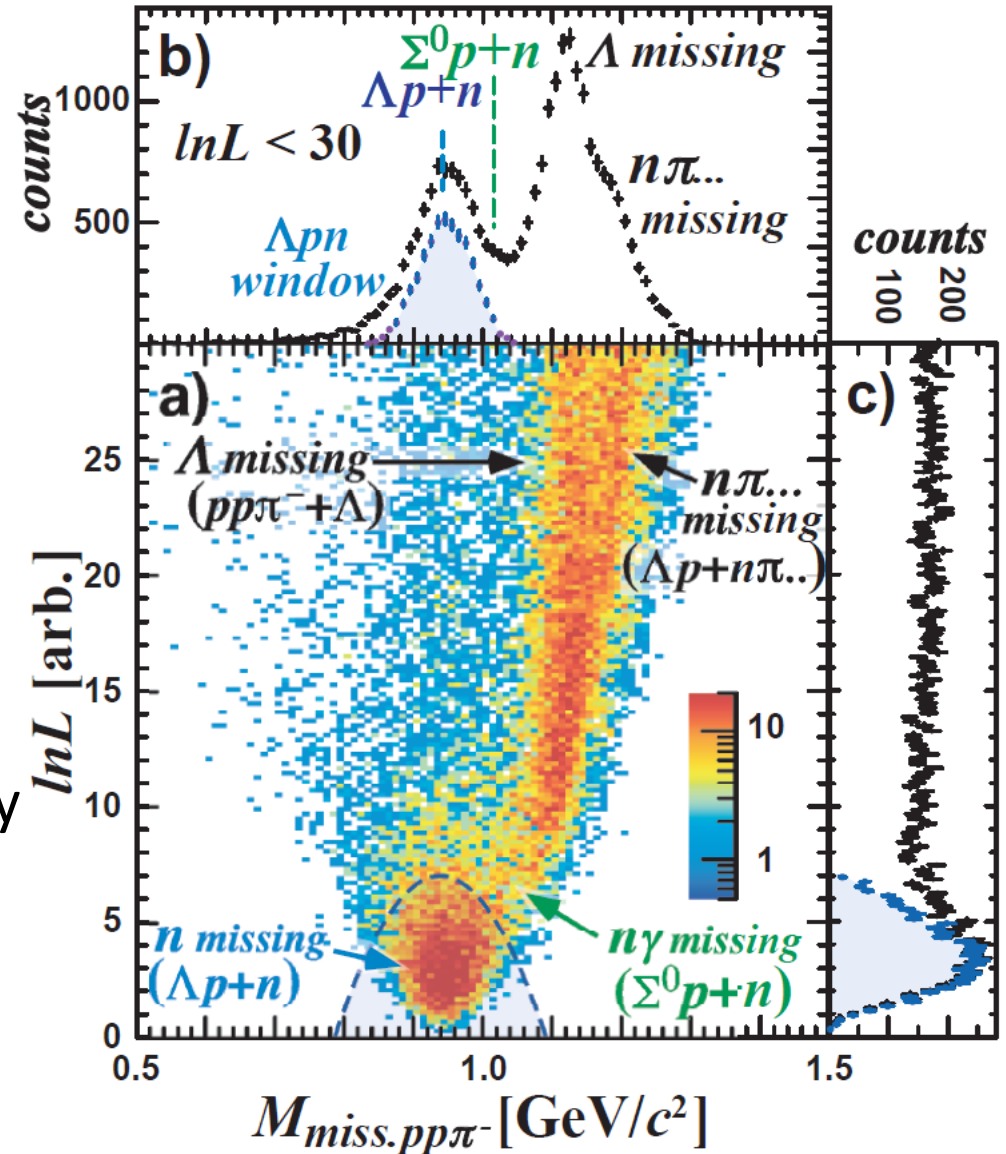


${}^3\text{He} + \text{K}^- \rightarrow \Lambda p n$ Selection



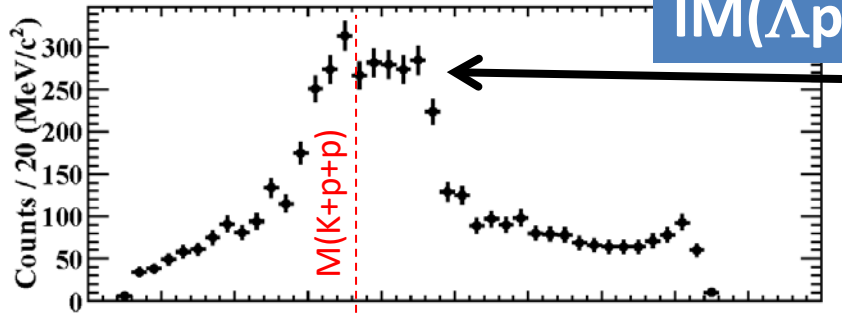
- $\Lambda \rightarrow \pi^- p$ and p are detected with CDS
 - A missing neutron is identified by missing-mass of ${}^3\text{He}(\text{K}^-, \Lambda p)n$
- $\Lambda p n_{\text{miss}}$ events are selected by log-likelihood method ($\ln L$)
 - distance-of-closest-approach for each vertex
 - kinematical constraint

E15 collab., arXiv:1805.12275

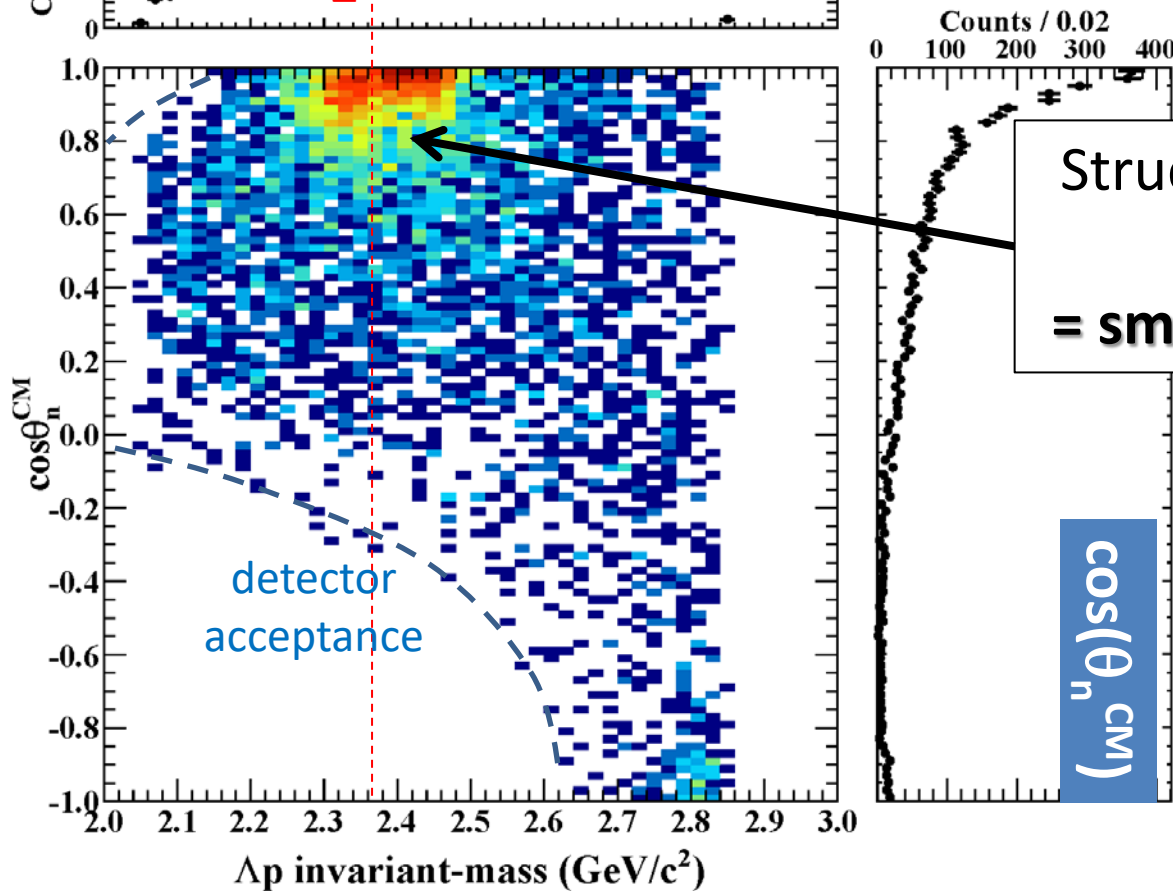


$IM(\Lambda p)$ vs. $\cos(\theta_n^{CM})$

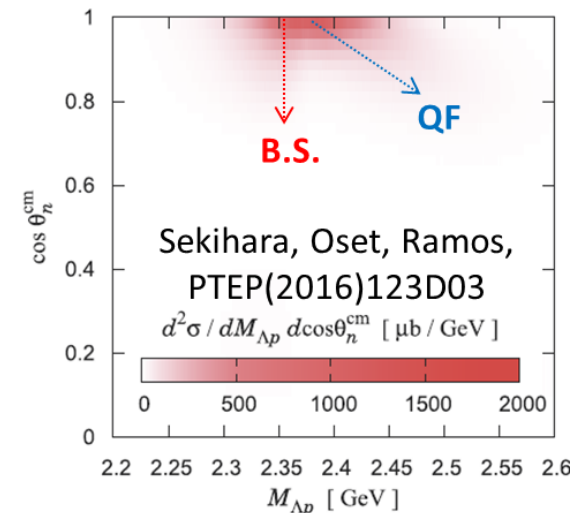
E15 collab., arXiv:1805.12275



Structures around the K^-pp threshold can be seen
= **bound-state + quasi-elastic**

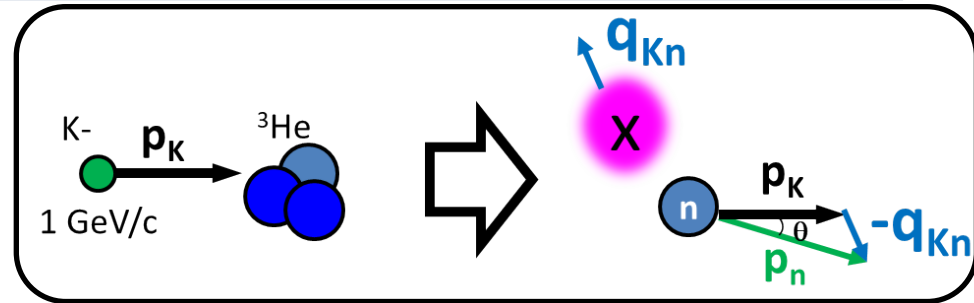
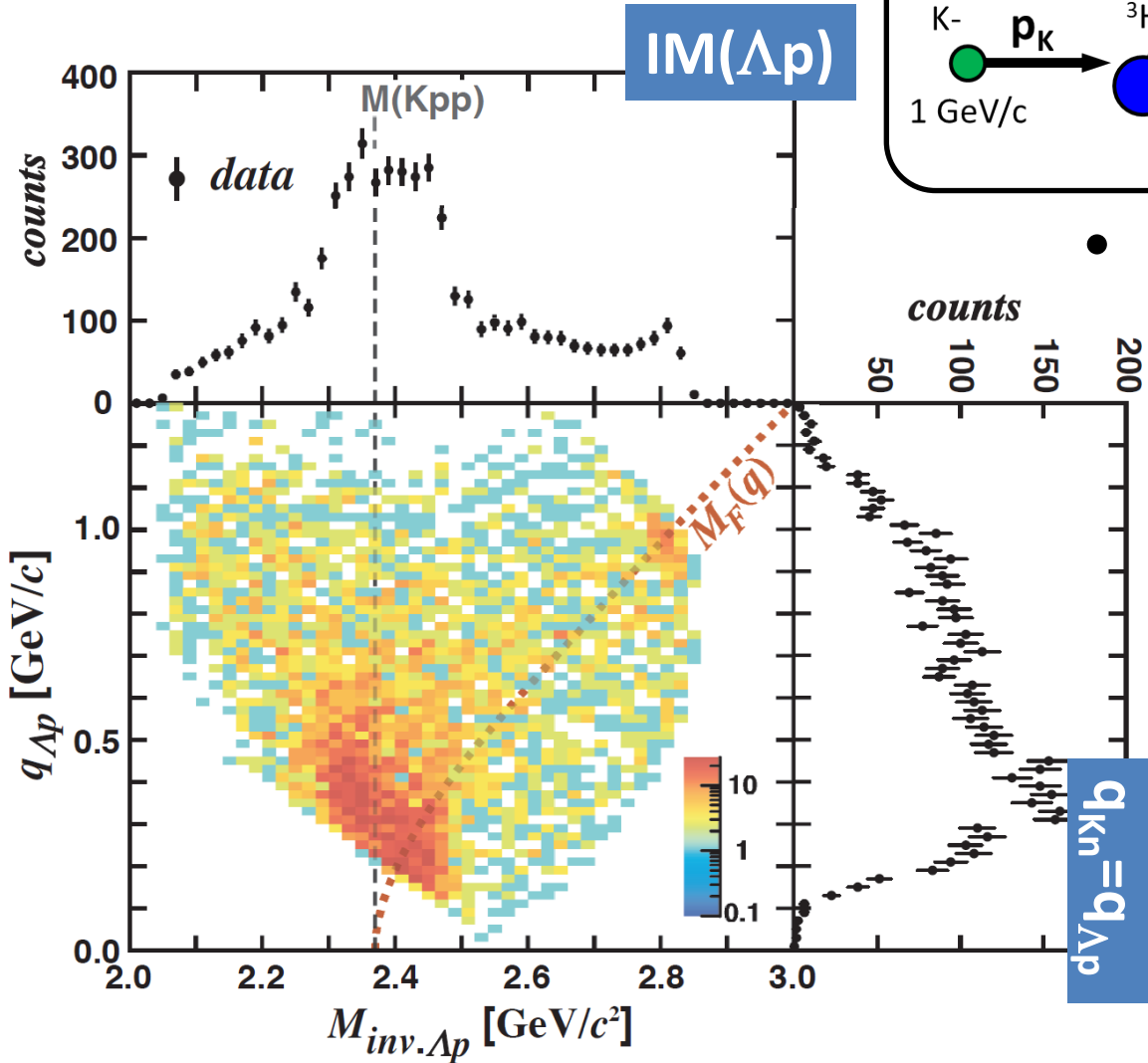


Structures are concentrated in forward-n region
= **small momentum-transfer**



IM(Δp) vs. Momentum Transfer q_{Kn}

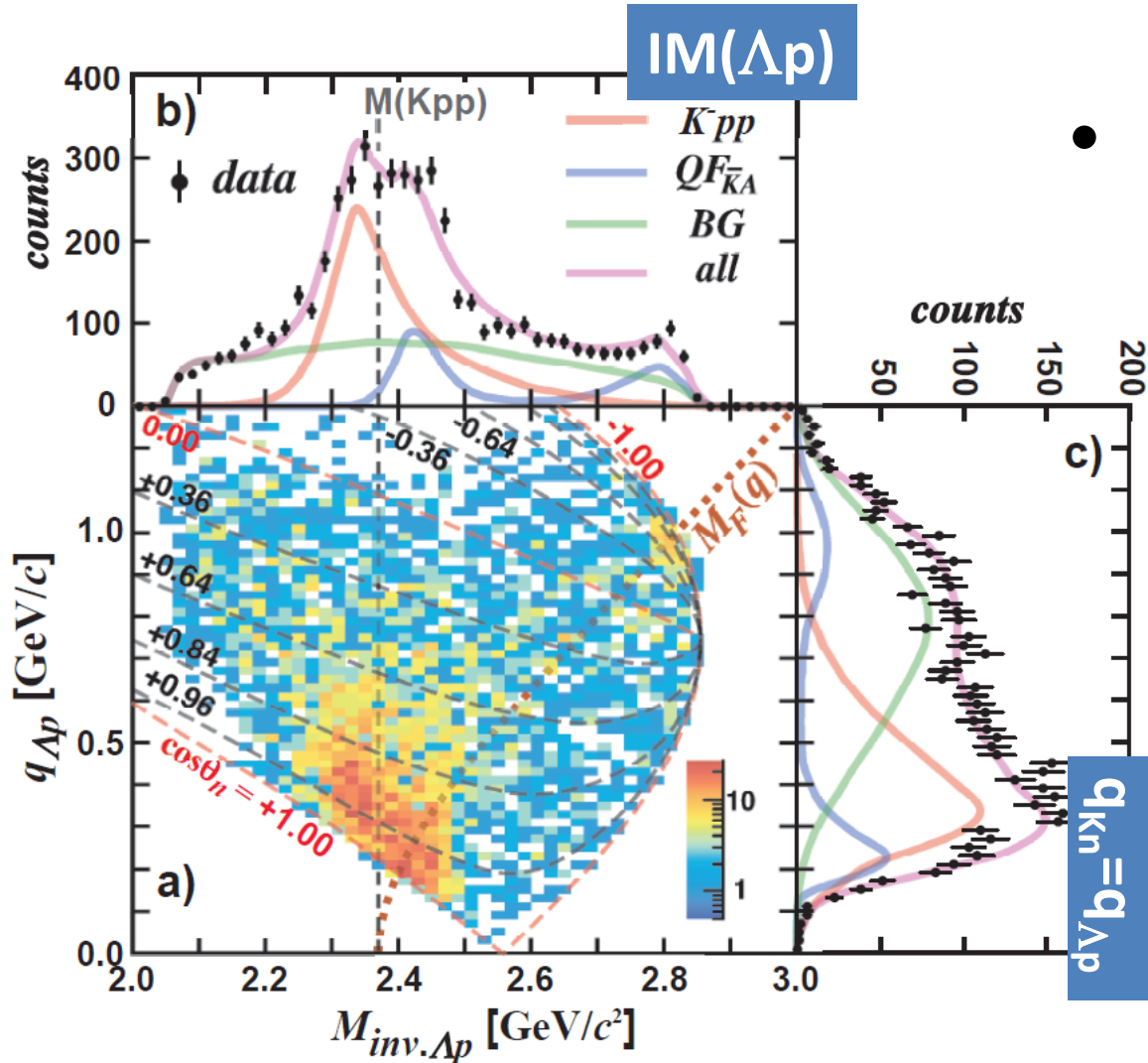
E15 collab., arXiv:1805.12275



- Seems to consist of 3 components
- **Bound state**
 - centroid NOT depend on q_{Kn}
- **Qasi-elastic K^- abs.**
 - centroid depends on q_{Kn}
- **Background**
 - Broad distribution

IM(Λp) vs. Momentum Transfer q_{Kn}

E15 collab., arXiv:1805.12275



- Fit with 3 components

– Bound state

- **centroid NOT depend on q_{Kn}**

- BW*(Gauss form-factor)

$$f_{\{Kpp\}}(M, q) = \frac{A_{Kpp} (\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} e^{-\left(\frac{q}{Q_{Kpp}}\right)^2}$$

– Qasi-elastic K⁻ abs.

- **centroid depends on q_{Kn}**

- Followed by Λp conversion

– Background

- Broad distribution

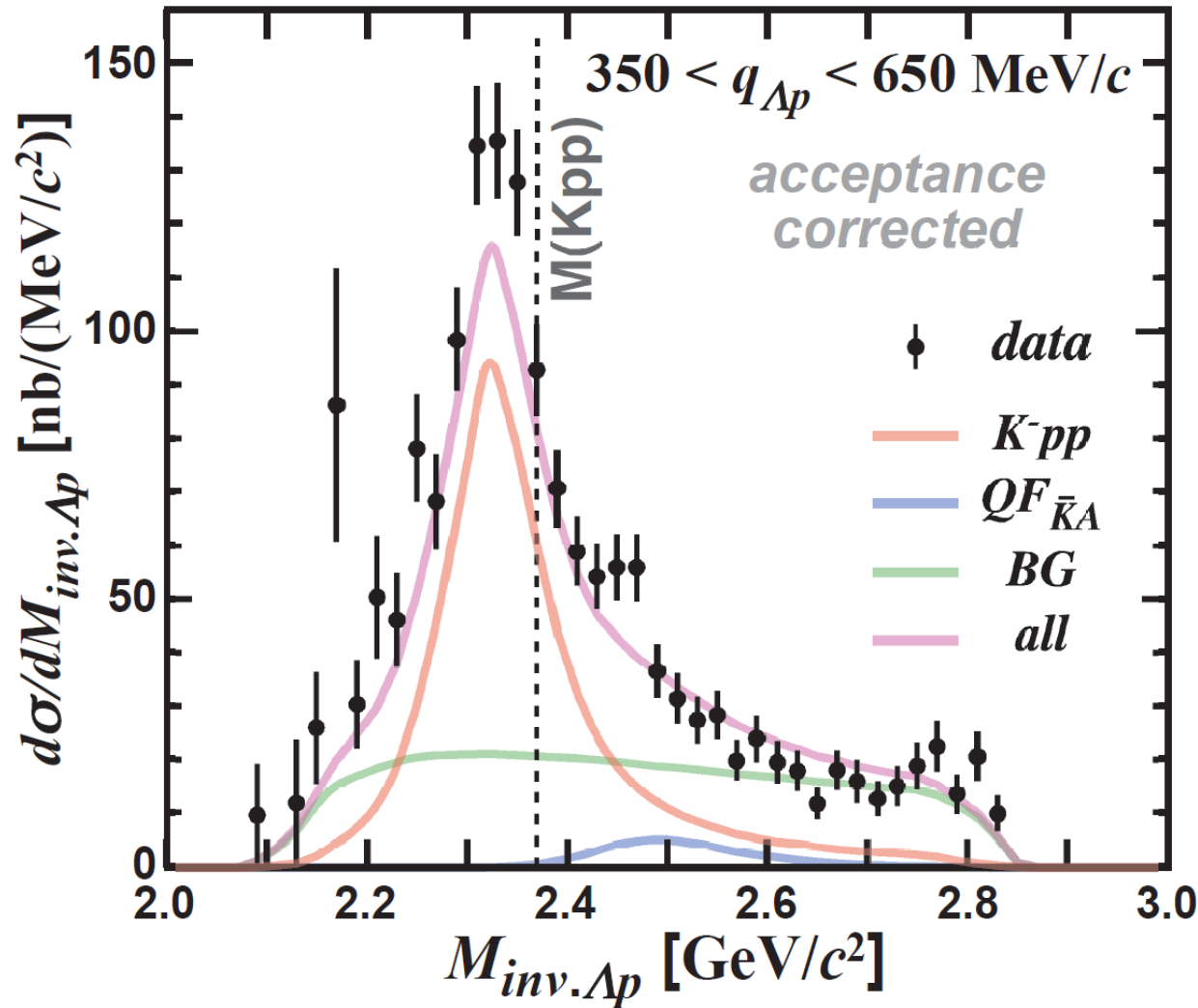
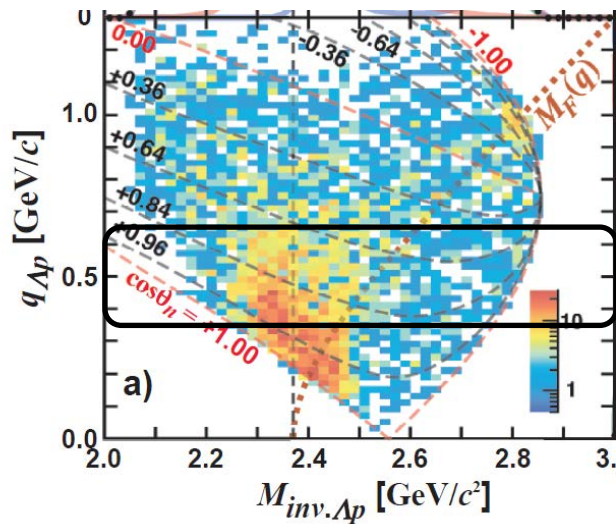
* We conduct the fitting in each 2D bin

“K-pp” Bound-State

$$f_{\{Kpp\}}(M, q) = \frac{A_{Kpp} (\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} e^{-\left(\frac{q}{Q_{Kpp}}\right)^2}$$

Select $0.35 < q_{Kn} < 0.65$
GeV/c

– BS and QF are well separated



"K⁻pp" Bound-State

$$f_{\{Kpp\}}(M, q) = \frac{A_{Kpp} (\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} e^{-\left(\frac{q}{Q_{Kpp}}\right)^2}$$

Fit values
that reproduce the spectrum:

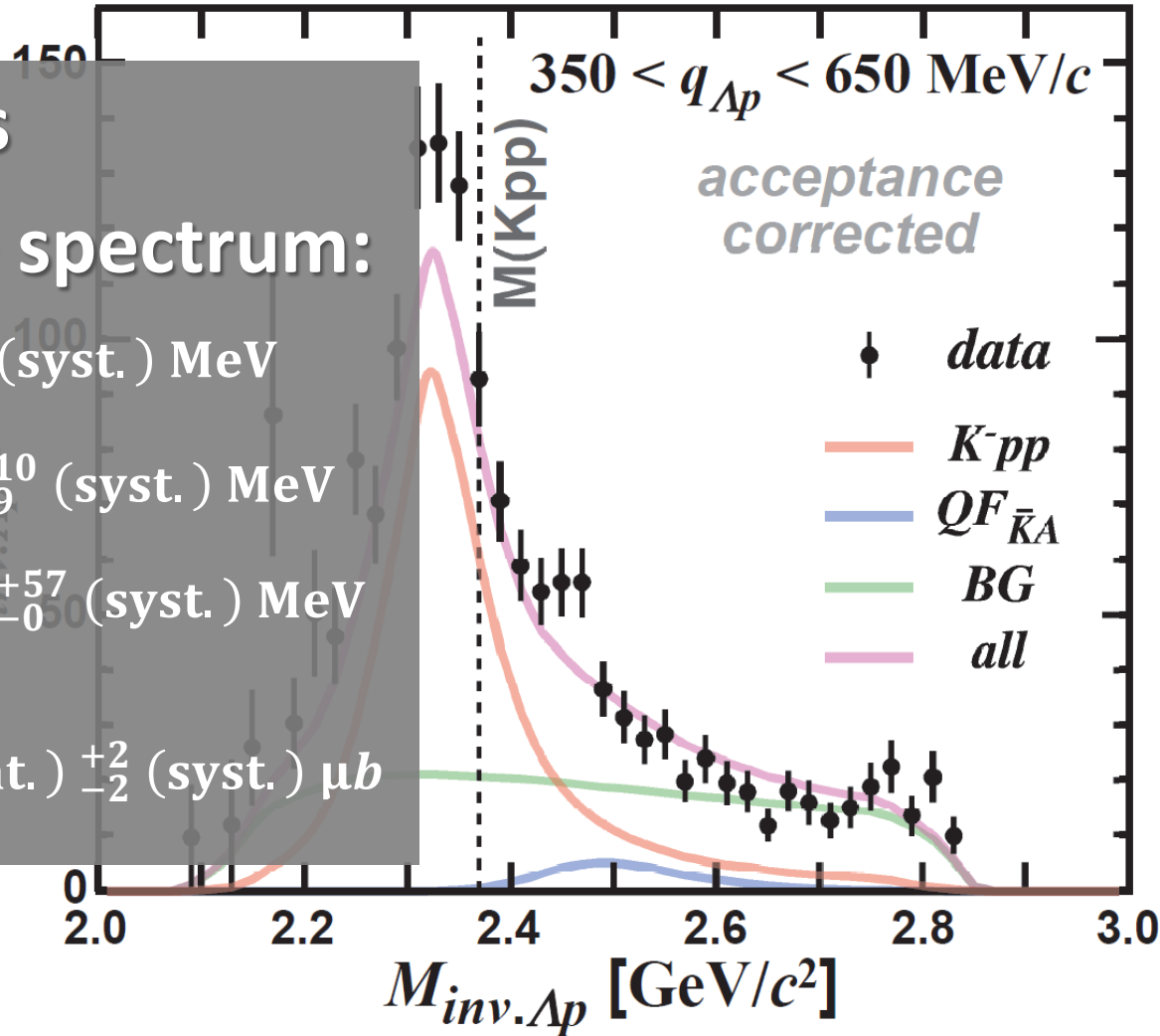
$$B_{"Kpp"} = 47 \pm 3 \text{ (stat.) } +3_{-6} \text{ (syst.) MeV}$$

$$\Gamma_{"Kpp"} = 115 \pm 7 \text{ (stat.) } +10_{-9} \text{ (syst.) MeV}$$

$$Q_{"Kpp"} = 381 \pm 14 \text{ (stat.) } +57_{-0} \text{ (syst.) MeV}$$

at below the $M(K^-pp)$

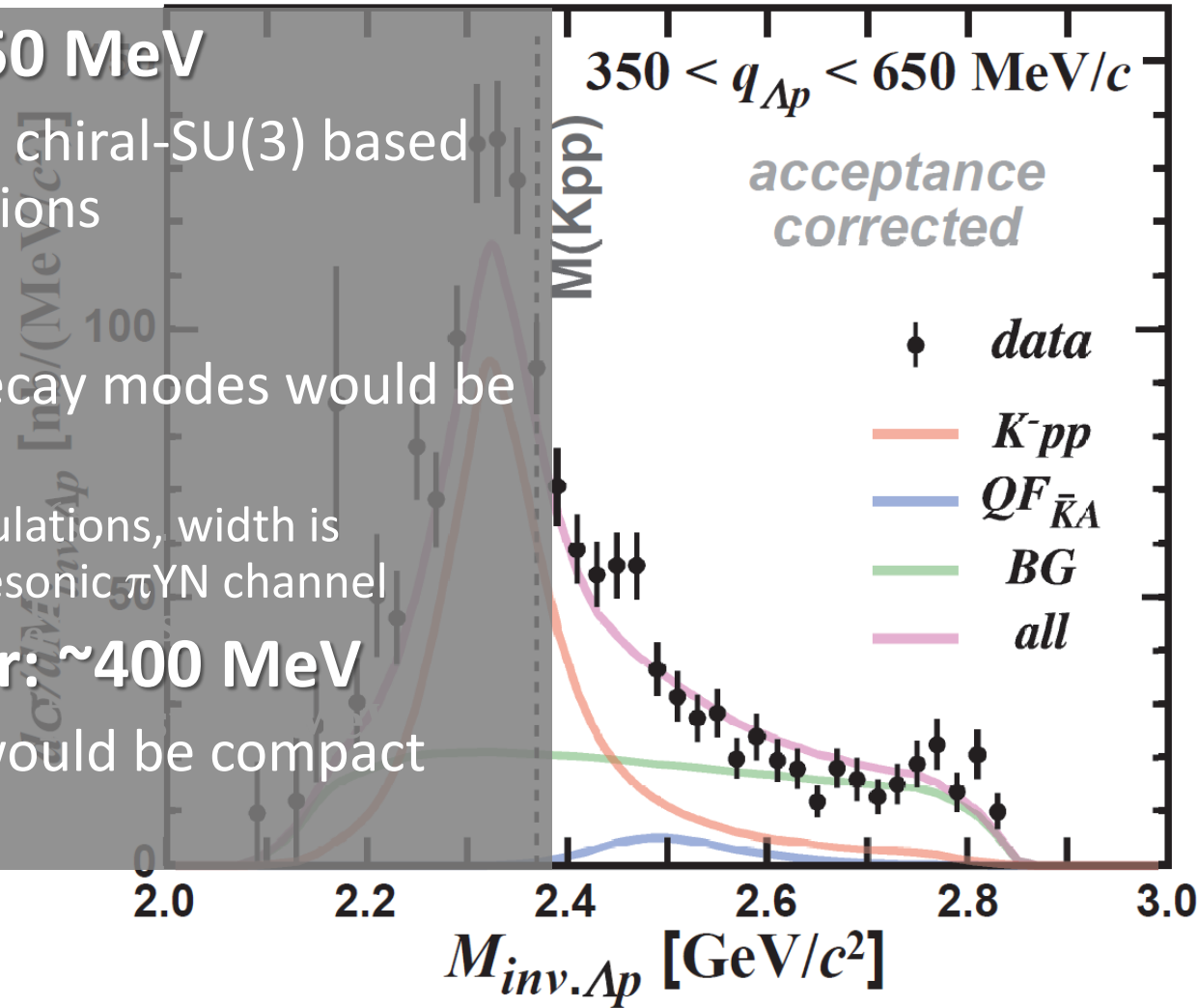
$$\sigma_{"Kpp"} \cdot Br_{\Lambda p} = 15 \pm 1 \text{ (stat.) } +2_{-2} \text{ (syst.) } \mu b$$



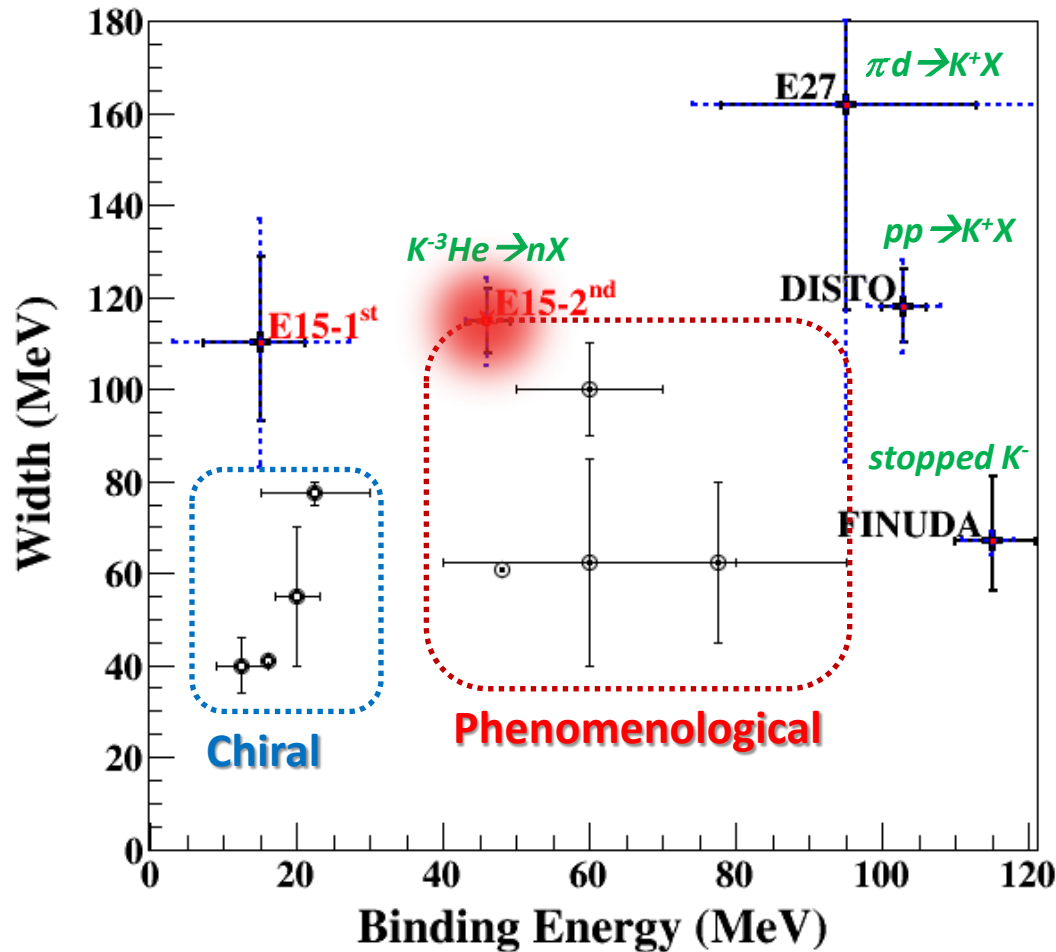
“K-pp” Bound-State

$$f_{\{Kpp\}}(M, q) = \frac{A_{Kpp} (\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} e^{-\left(\frac{q}{Q_{Kpp}}\right)^2}$$

- **Binding energy: ~50 MeV**
 - Much deeper than chiral-SU(3) based theoretical predictions
- **Width: ~100 MeV**
 - Non-mesonic YN decay modes would be dominant
 - in theoretical calculations, width is evaluated with mesonic π YN channel
- **S-wave form factor: ~400 MeV**
 - $K^- + {}^3\text{He}$ reaction would be compact (~0.5 fm)



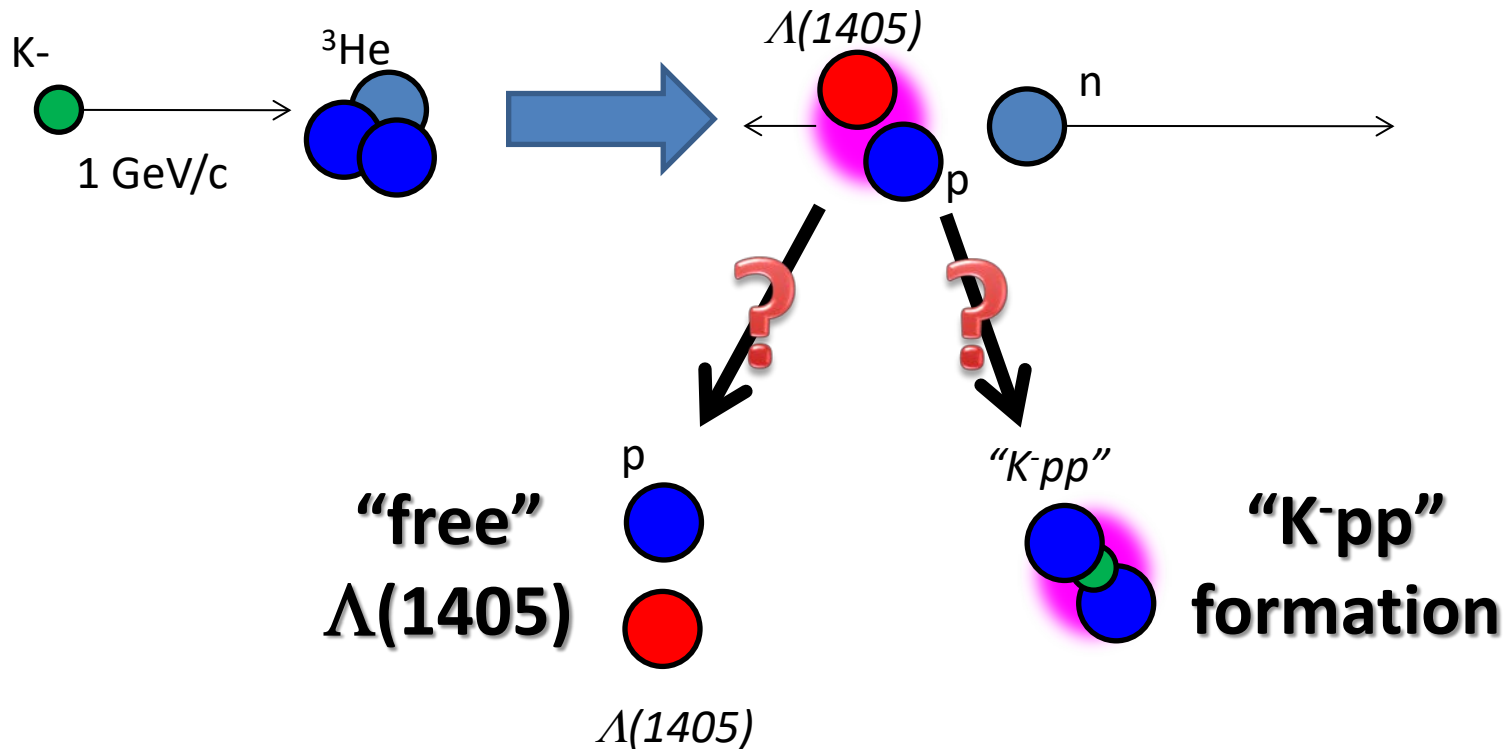
Present Status of “K⁻pp”



For further understandings:

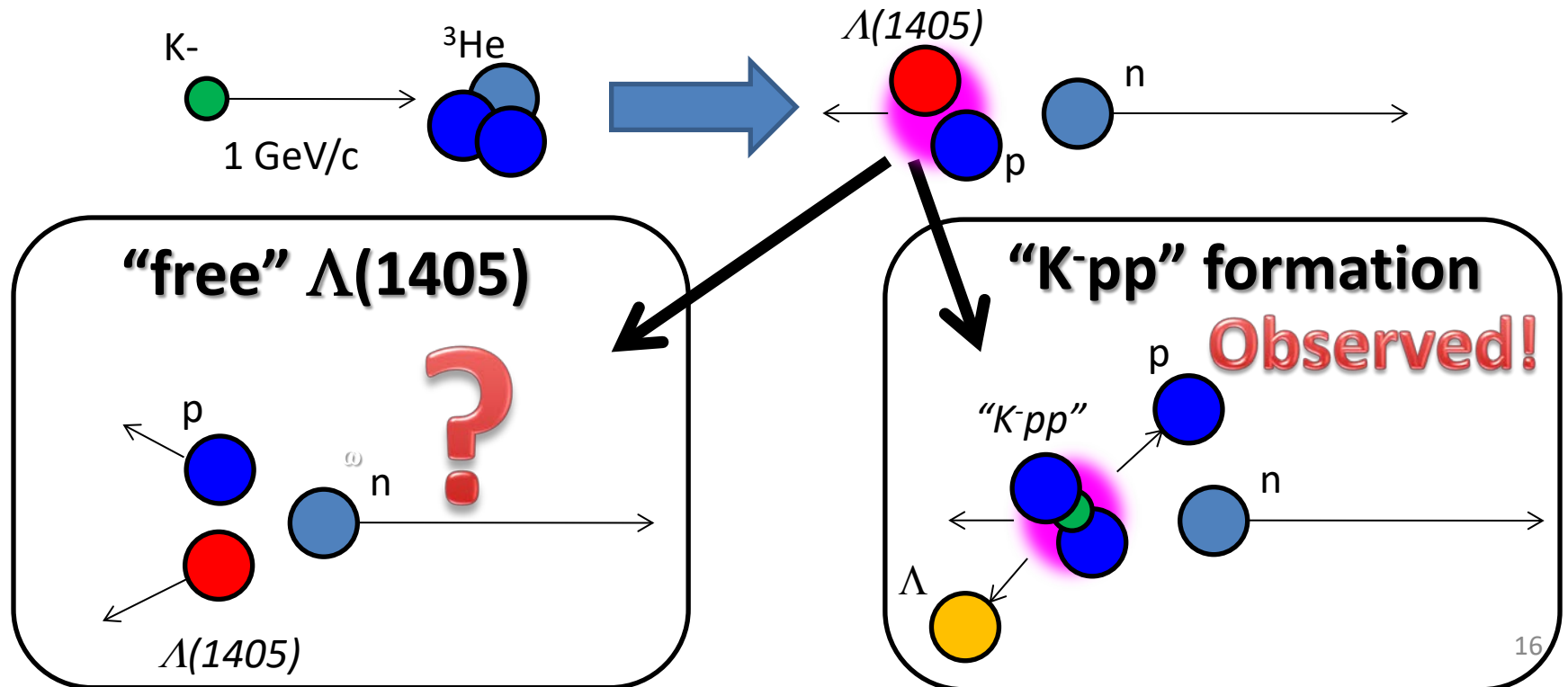
- ✓ $\Lambda(1405)$ production $\rightarrow \Lambda^*N$ doorway
- ✓ $\pi\Sigma N$ decay channel \rightarrow new info. of $K^{\text{bar}}NN$

$\Lambda(1405)$ in ${}^3\text{He}(K^-, \pi\Sigma p)n$



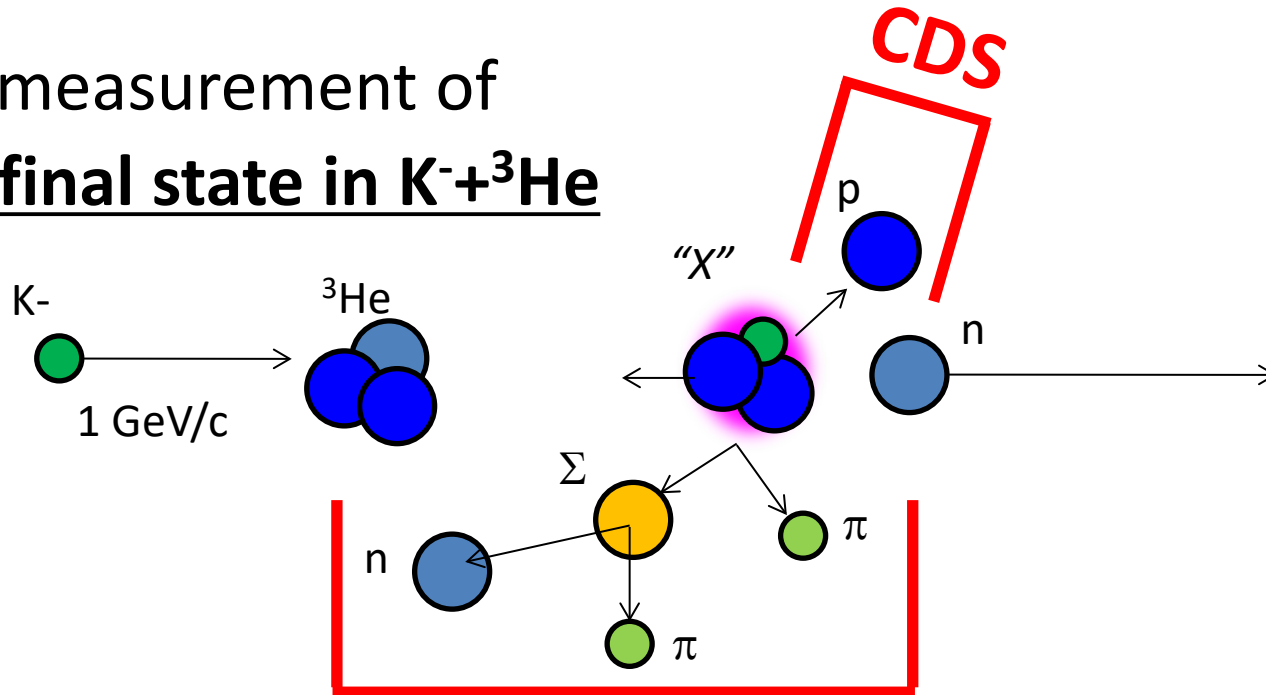
$\Lambda(1405)p$ and “K⁻pp”

- Theoretically, “K⁻pp” is expected to be produced via $\Lambda(1405)+p \rightarrow$ “K⁻pp” door-way process
 - comparison between $\Lambda(1405)p$ and “K⁻pp” production would give us an important information



$K^- \ ^3\text{He} \rightarrow \pi \Sigma \text{pn} @ \text{E15}$

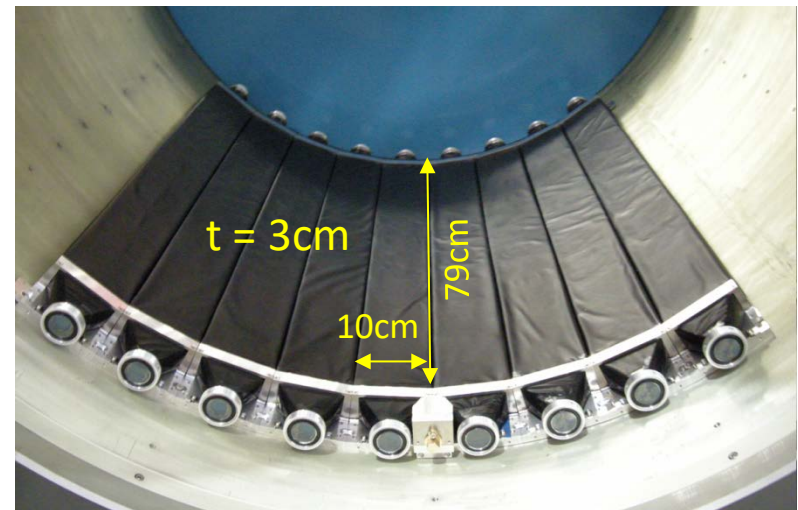
- Exclusive measurement of $\pi^\pm \Sigma^\mp \text{pn}$ final state in $K^- + ^3\text{He}$



CDS

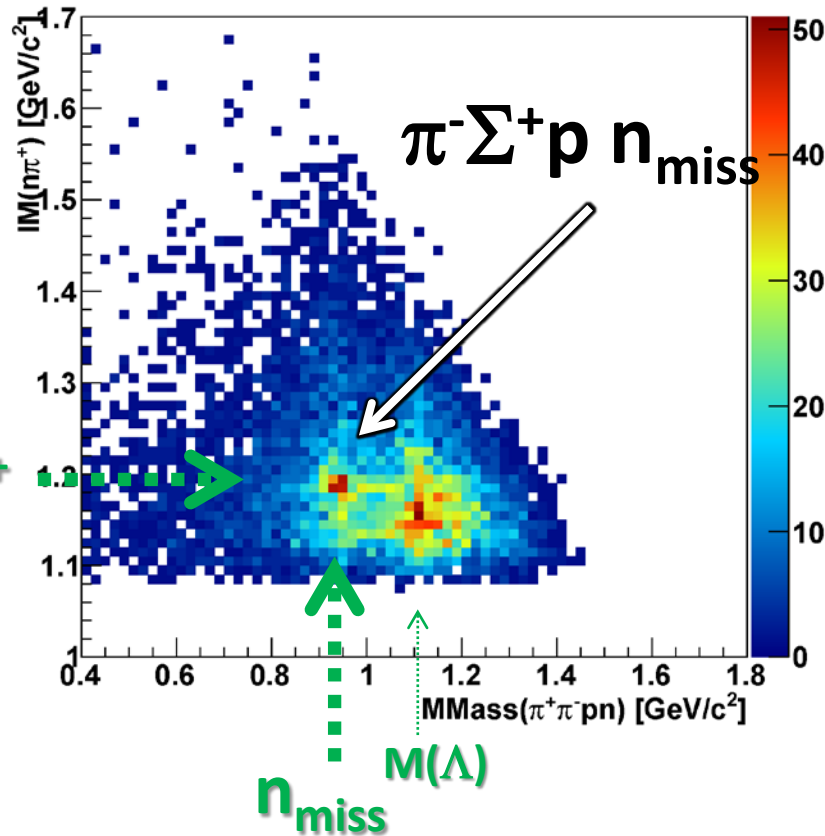
- Experimental challenge of neutron detection with thin scintillation counter ($t=3\text{cm}$)

n detection efficiency $\sim 3\%$

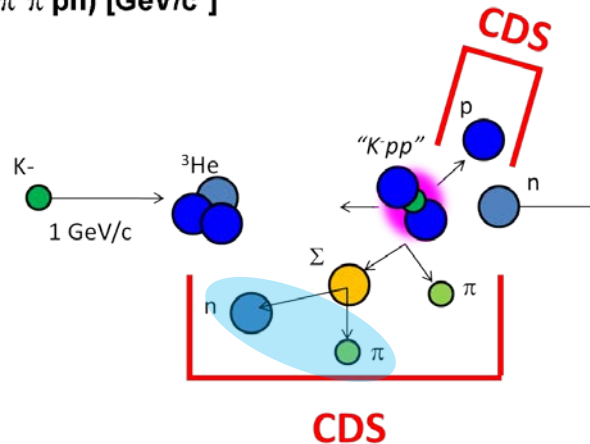
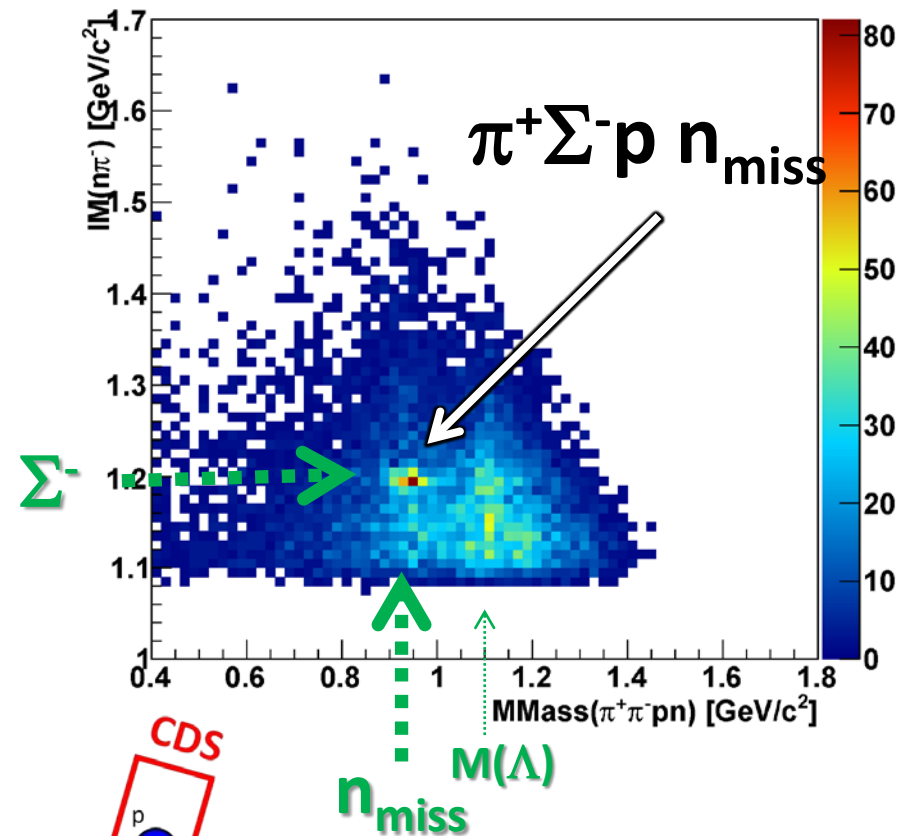


$\pi\Sigma\rho n$ Events

IM($n\pi^+$) vs MM($\pi^+\pi^-pn$)



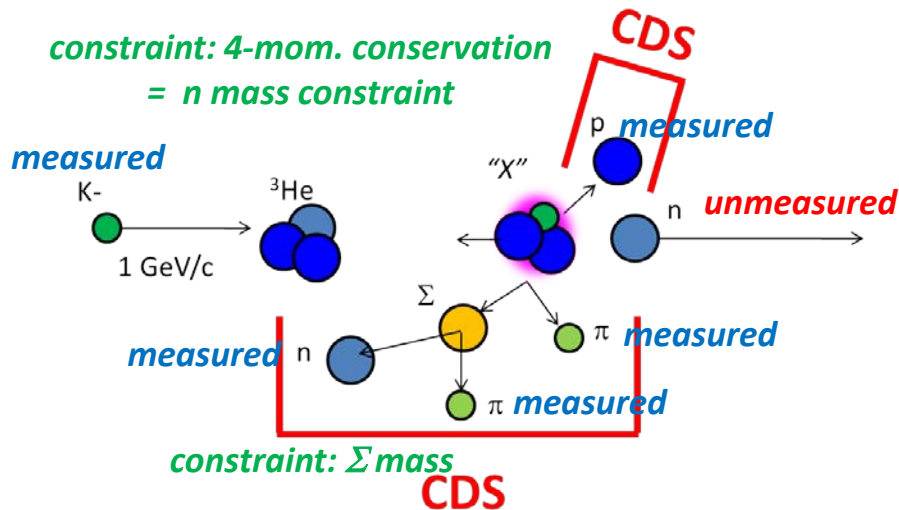
IM($n\pi^-$) vs MM($\pi^+\pi^-pn$)



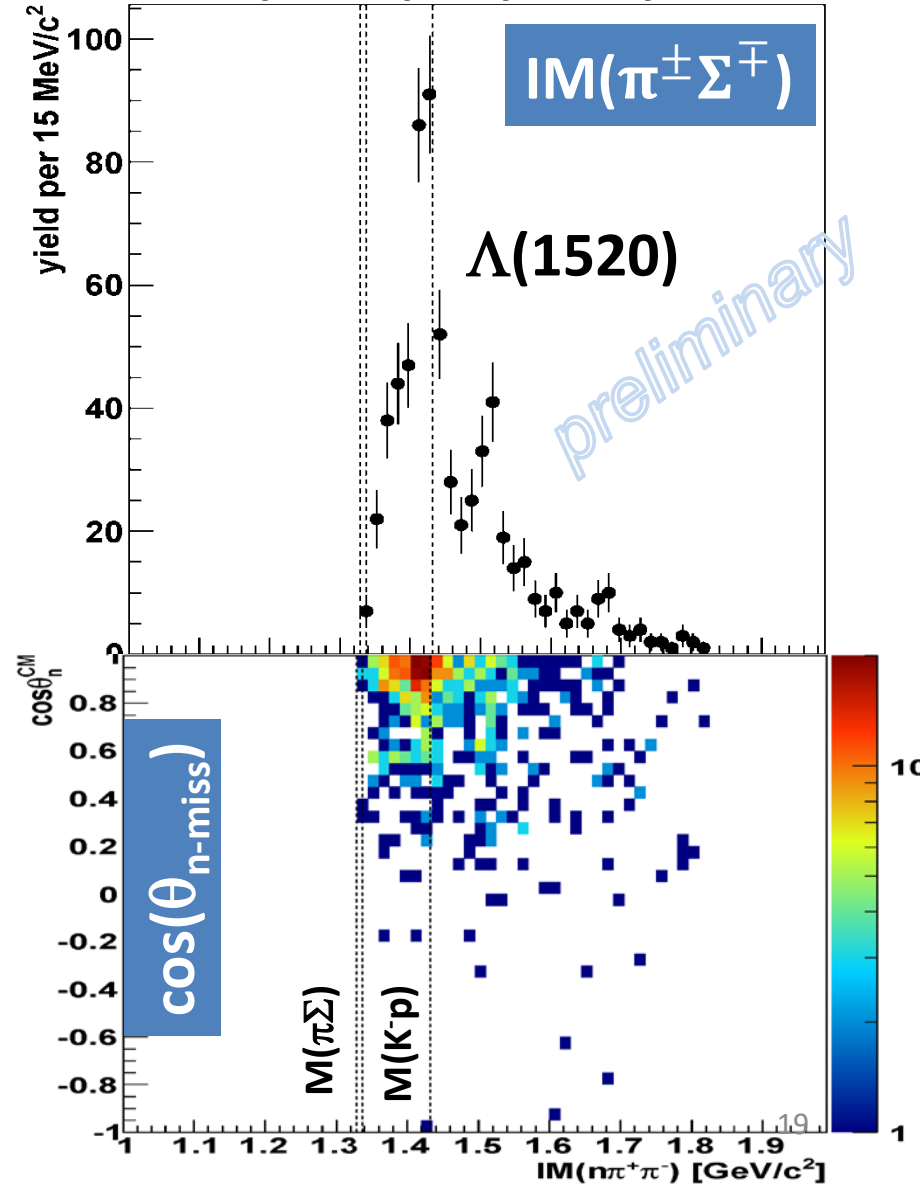
IM($\pi\Sigma$) vs. $\cos(\theta_n^{\text{CM}})$

- $\pi^\pm \Sigma^\mp$ events are separated using kinematical-fit

- Constraints:
 - $M(\Sigma \rightarrow n\pi)$
 - 4-momentum conservation
- Event selection by χ^2 probability ($0.01 < p$)

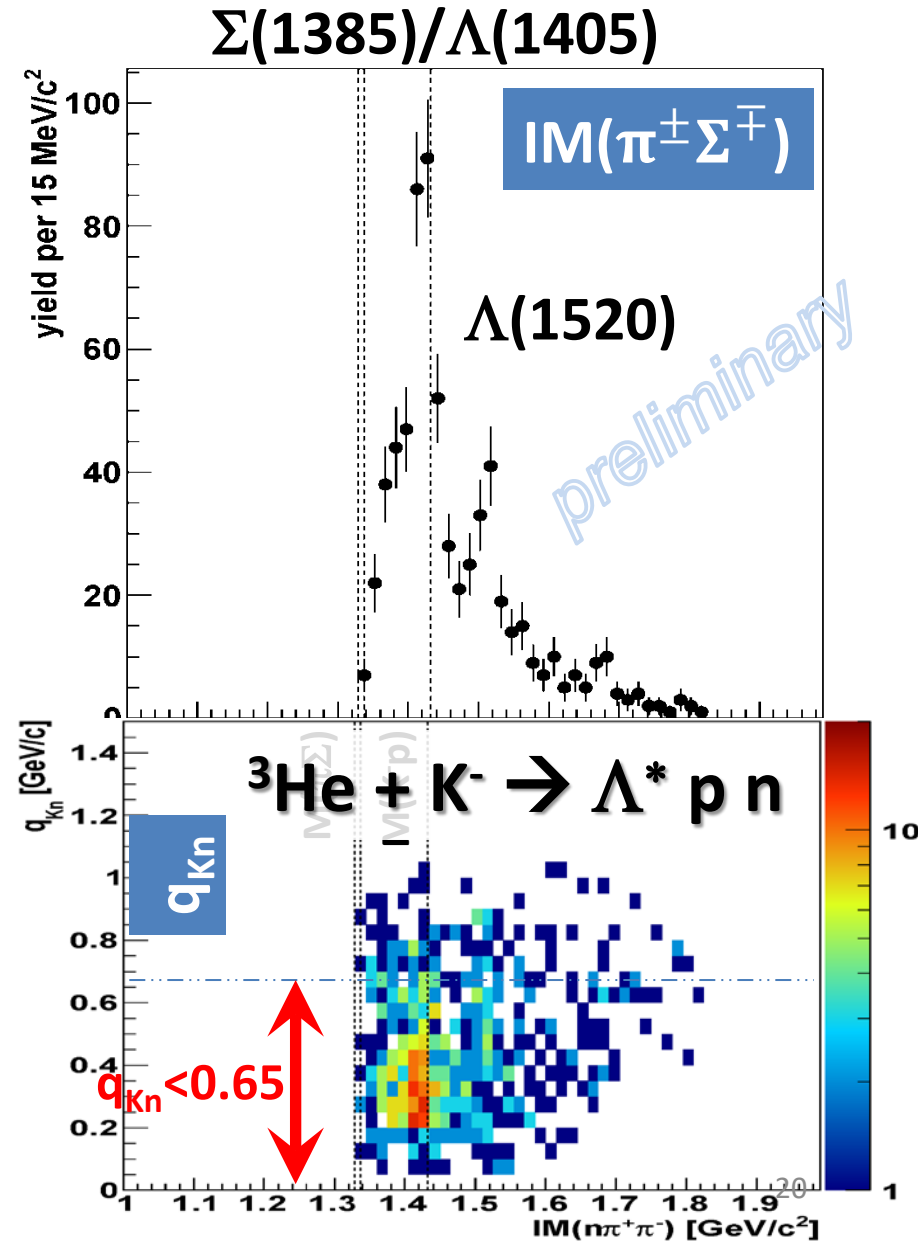
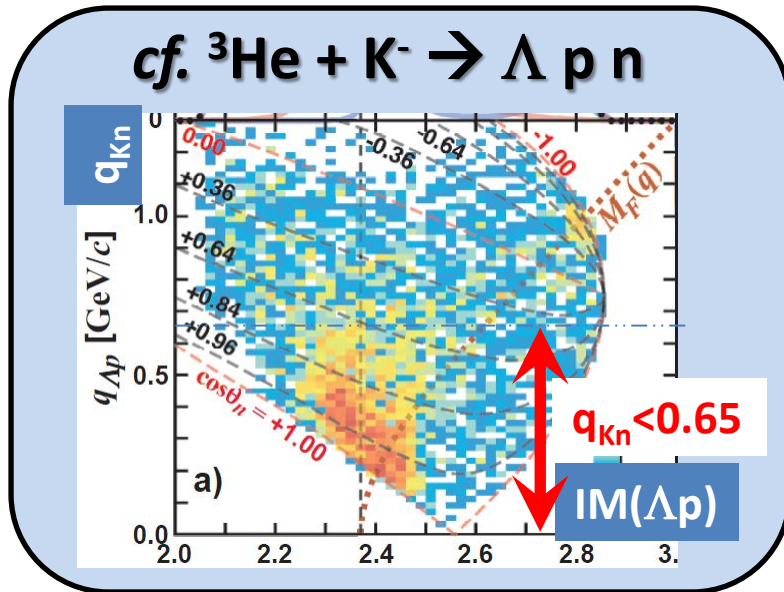


$\Sigma(1385)/\Lambda(1405)$



IM($\pi\Sigma$) vs. Momentum Transfer q_{Kn}

- To compare “K⁻pp” and Λ^* production CS’s, we select $q_{Kn} < 0.65$ GeV/c region
 - “K⁻pp” and Λ^* signals can be seen in this region



Υ^* CS ($q_{K_n} < 0.65$)

$\Lambda(1405)$

$\sim 130\text{-}140 \mu\text{b}$

Flatté param.:

$m_R \sim 1418 \text{ MeV}/c$

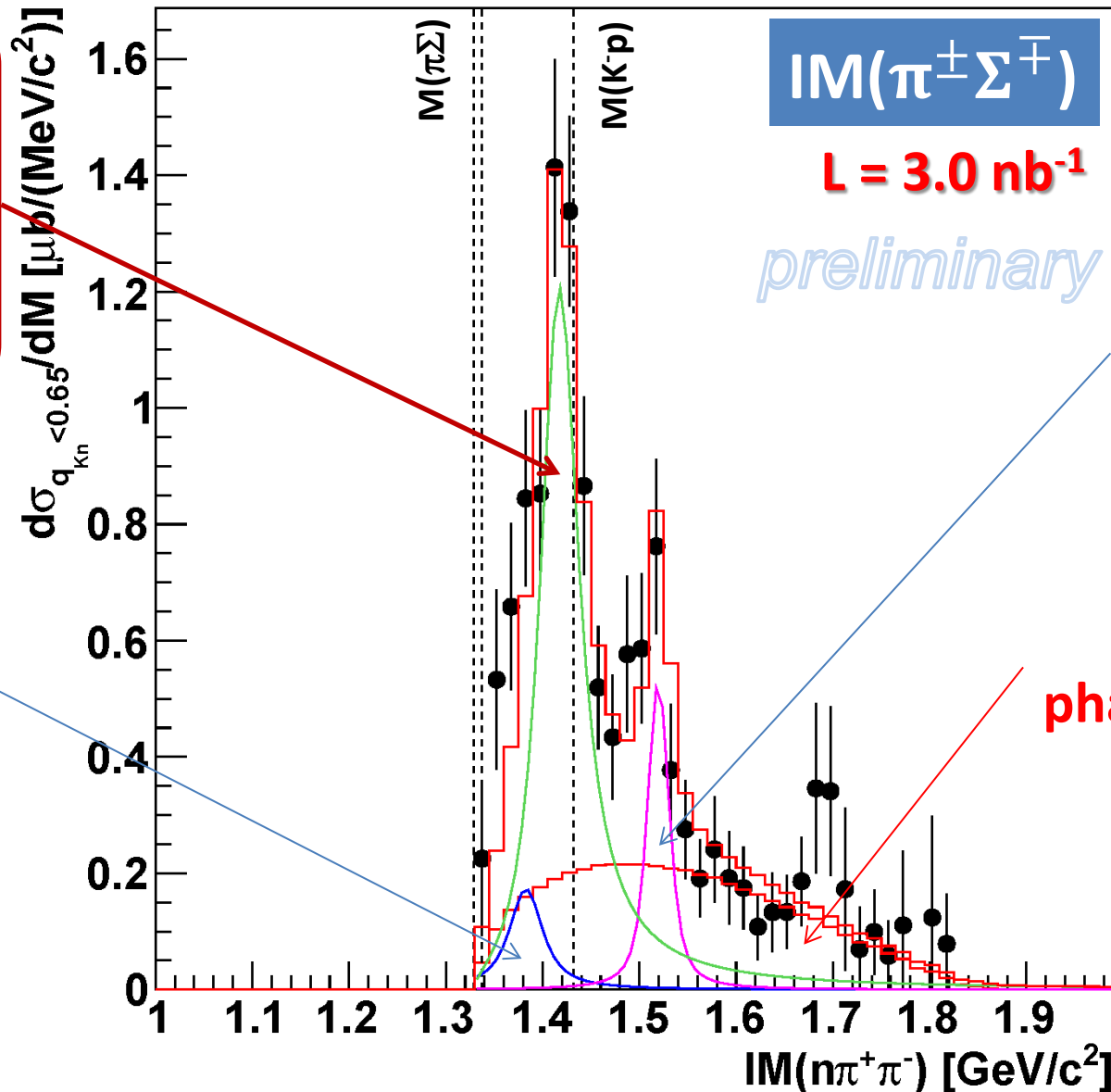
$g_{\pi\Sigma} \sim 1.9\text{E-}1$

$g_{KN} \sim 1.7\text{E-}2$

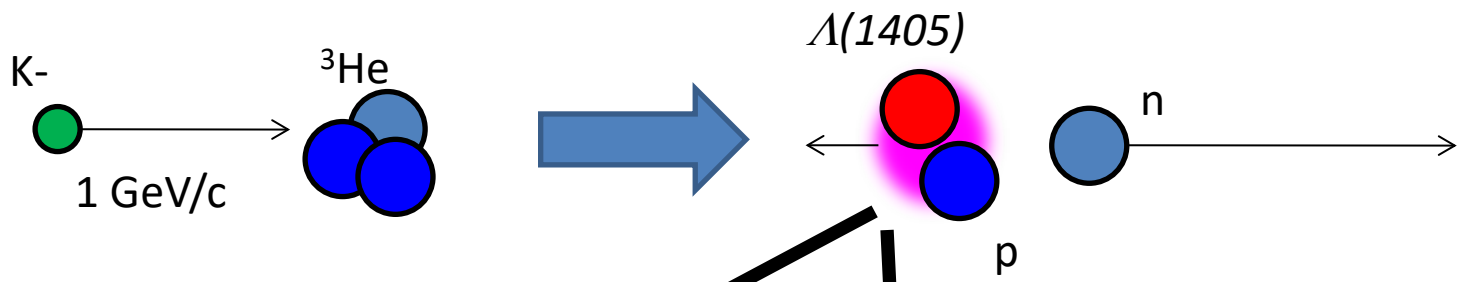
$\Sigma^0(1385)$

$\sim 40\text{-}80 \mu\text{b}$

[evaluated from
 $\Sigma^+(1385) \rightarrow \pi^+ \Lambda$
measurement]

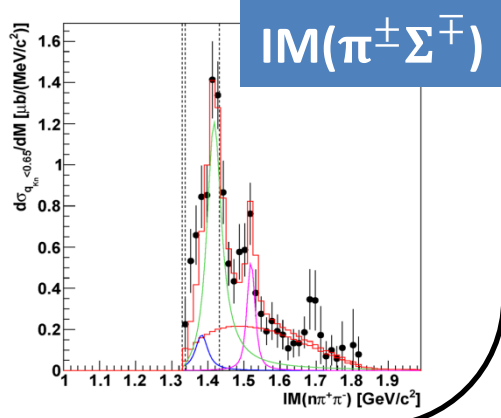
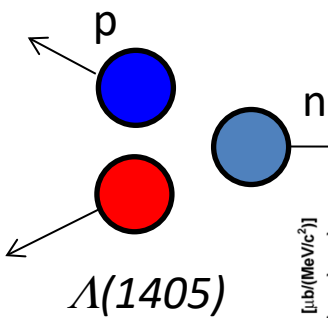


Production of $\Lambda(1405)p$ and "K-pp"



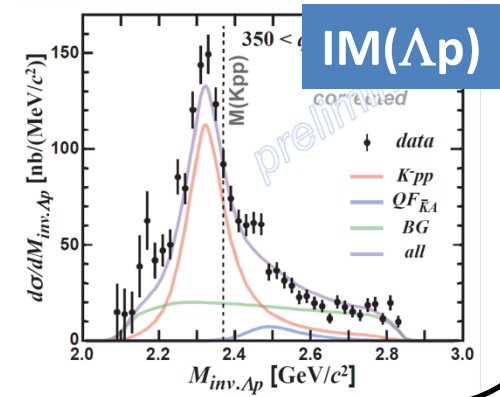
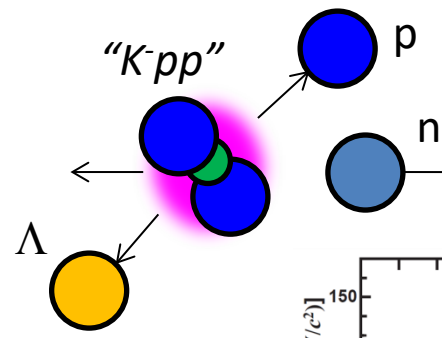
"free" $\Lambda(1405)$

$\sim 130\mu\text{b}$



"K-pp" → Λp

$\sim 20\mu\text{b}$



Large CS of Λ^* compared to "K-pp" formation

“K⁻pp” in ³He(K⁻, πΣp)n

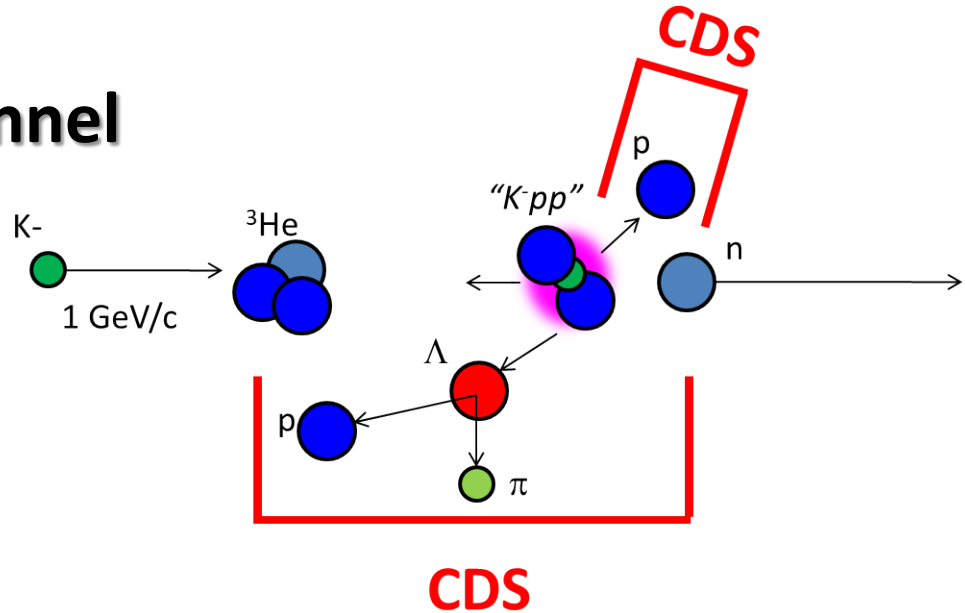
Search for “K⁻pp” → πΣN decay channel

Two Decay Mode of “K⁻pp”

1. “K⁻pp” search via Λp channel

→ Non-mesonic channel

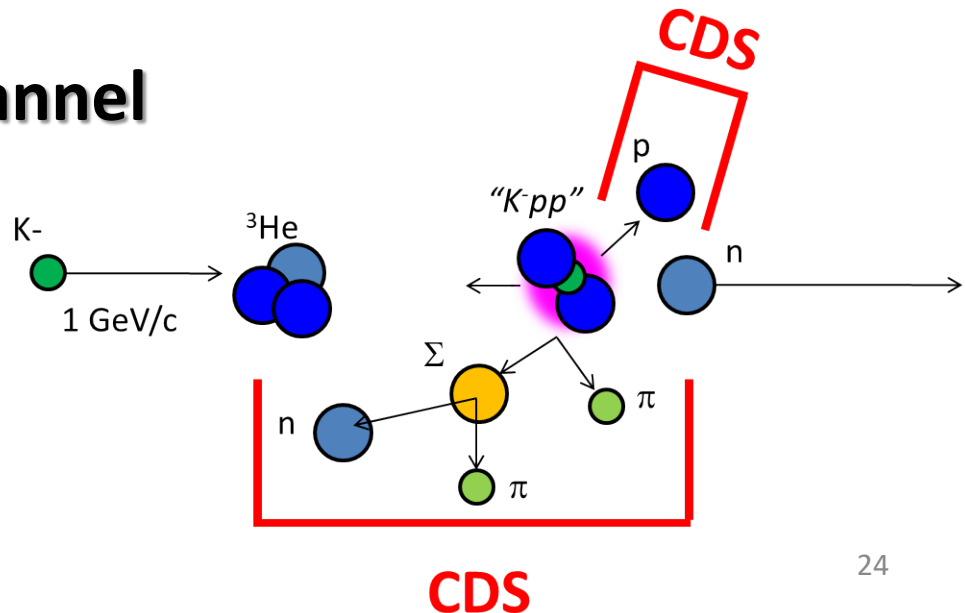
FINUDA/DISTO/E27/E15...



2. “K⁻pp” search via $\pi\Sigma p$ channel

→ Mesonic channel

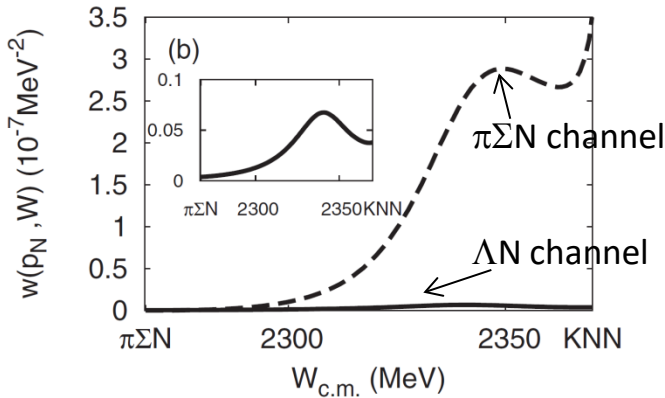
NO measurement so far



Two Decay Mode of “K⁻pp”

kaon absorption probability
of $\Lambda^* N \rightarrow \pi \Sigma N / \Lambda N$

S. Ohnishi et al., PRC88(2013)025204

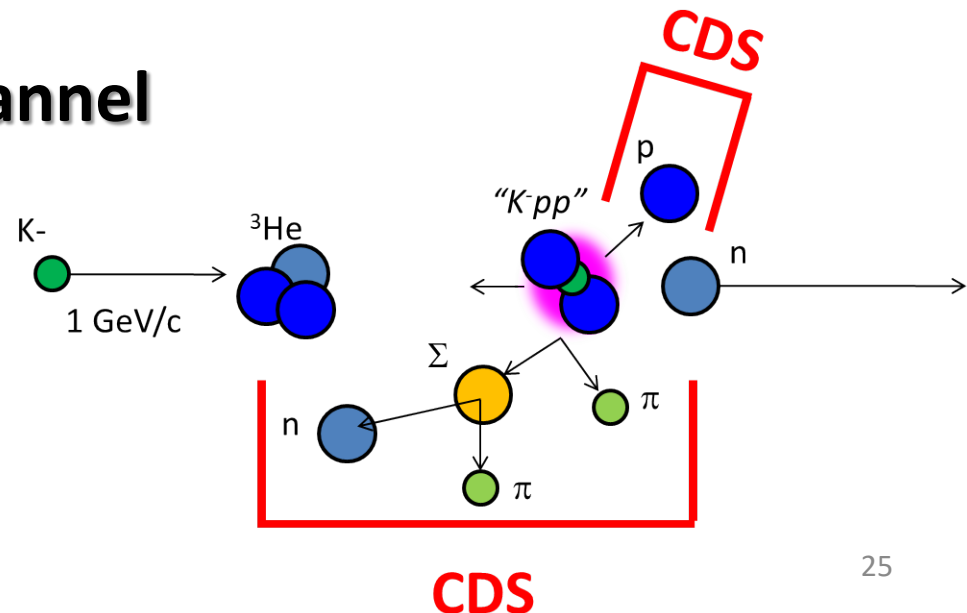


Theoretically,
 $\pi \Sigma N$ decay is expected to be
the dominant channel

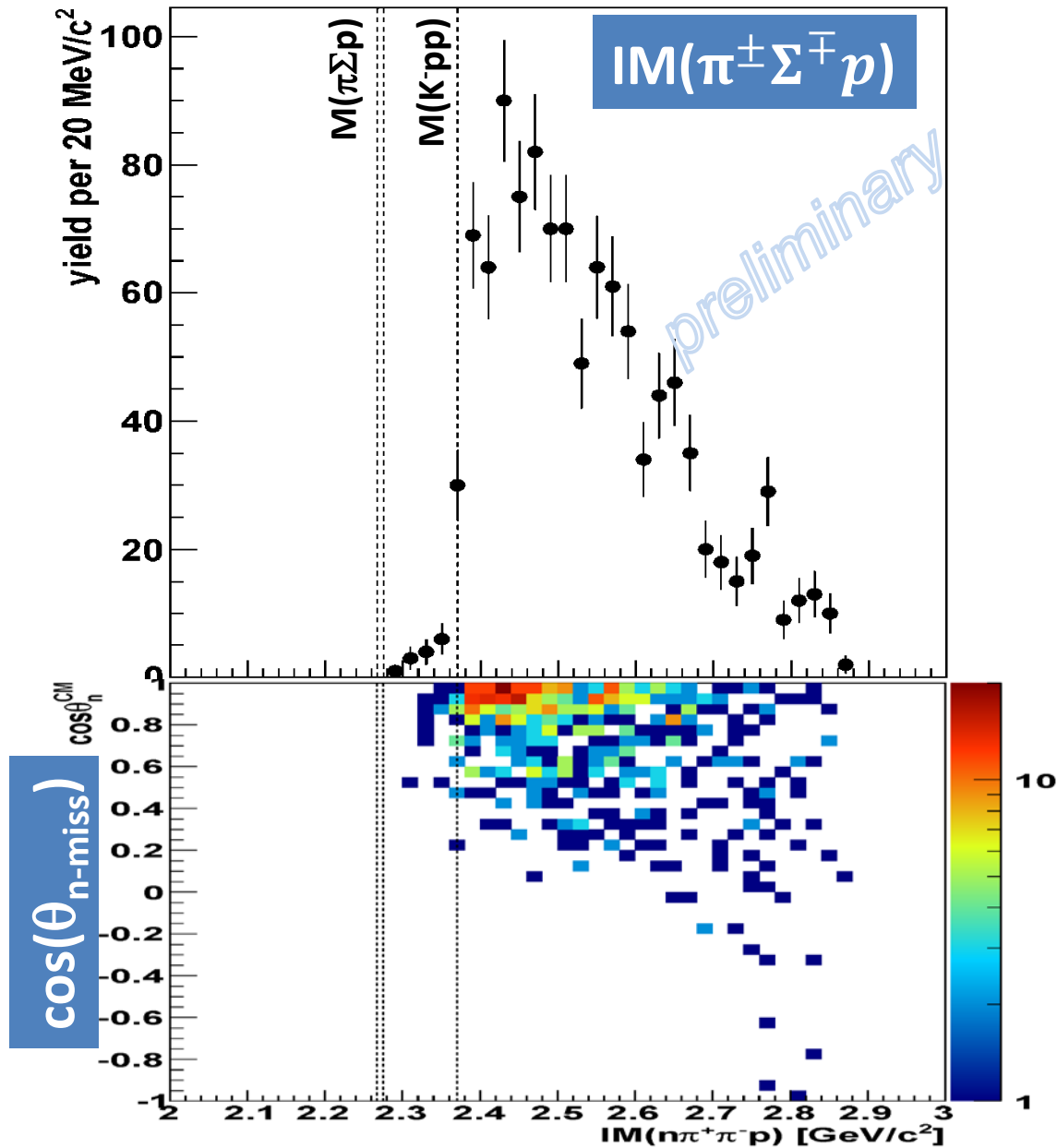
2. “K⁻pp” search via $\pi \Sigma p$ channel

→ Mesonic channel

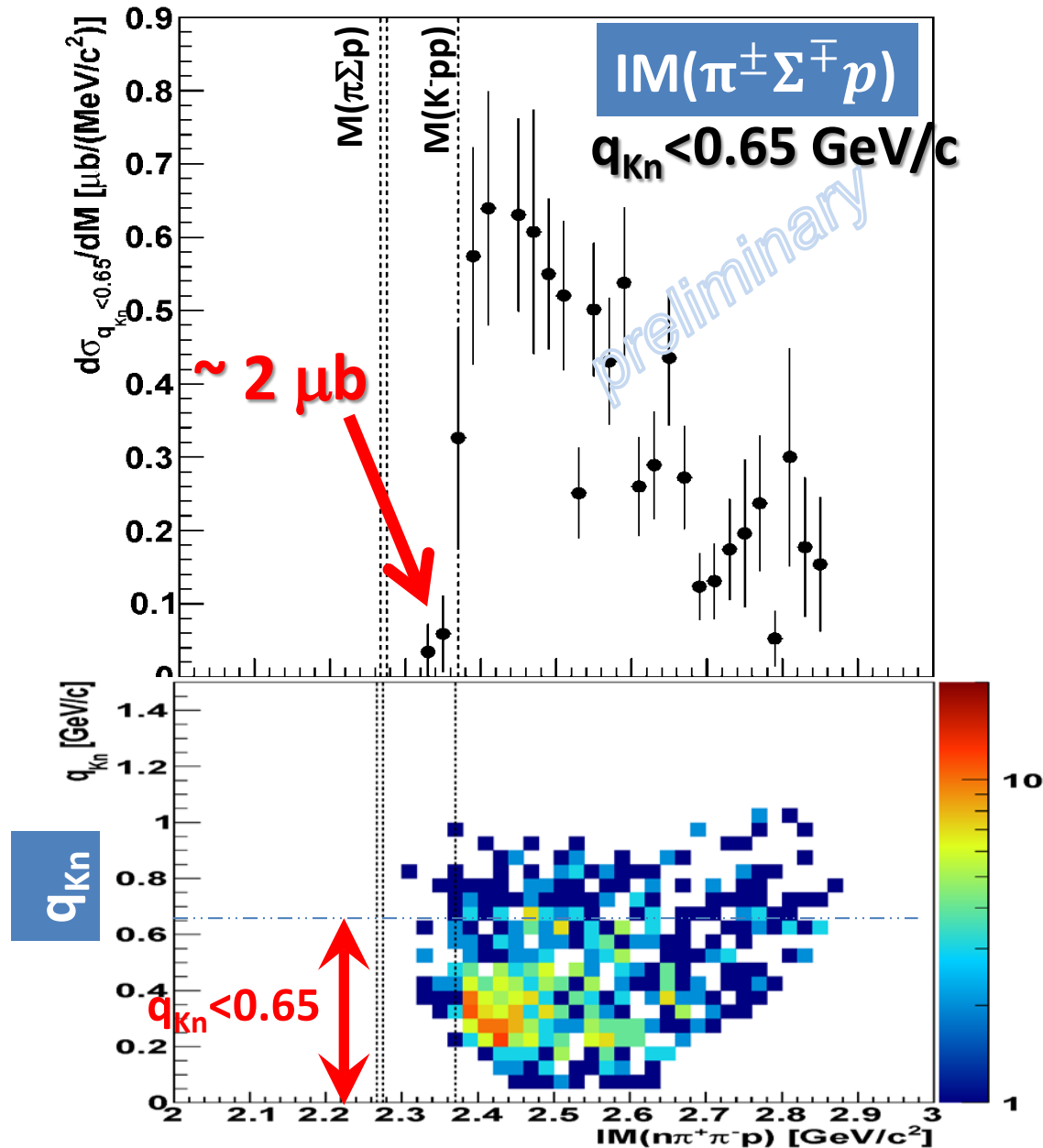
NO measurement so far

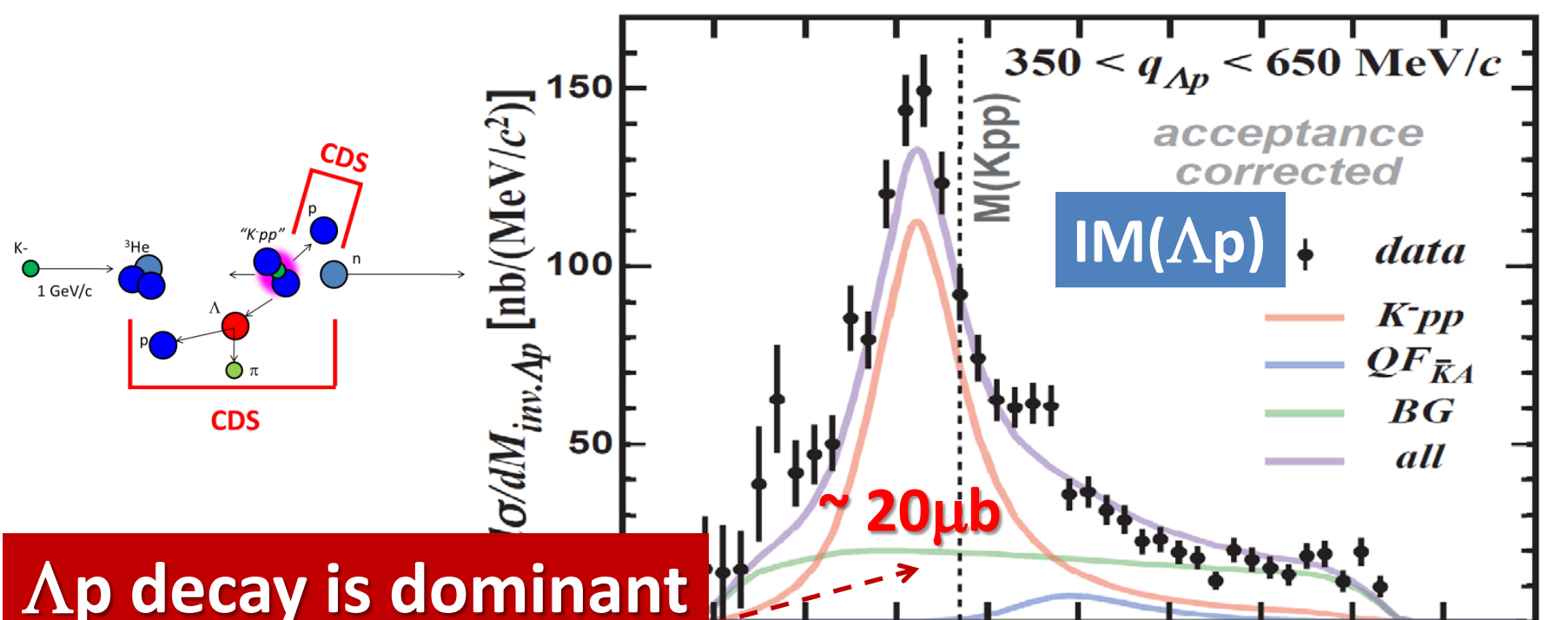


$IM(\pi\Sigma p)$ vs. $\cos(\theta_n^{CM})$

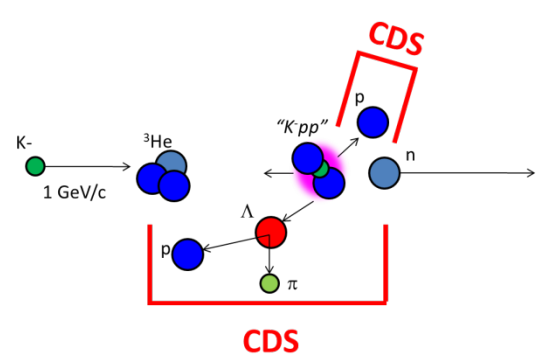
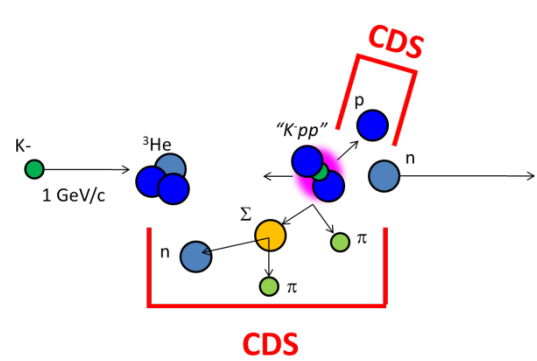
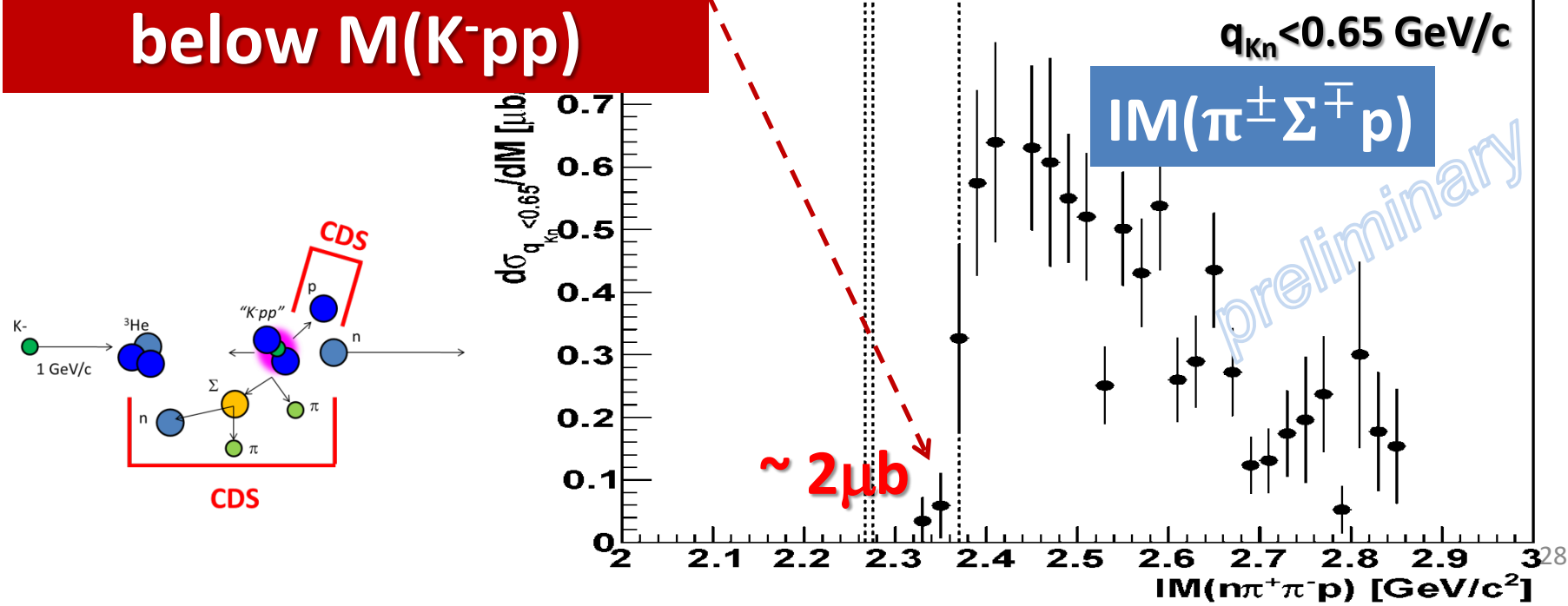


IM($\pi\Sigma p$) vs. Momentum Transfer q_{Kn}

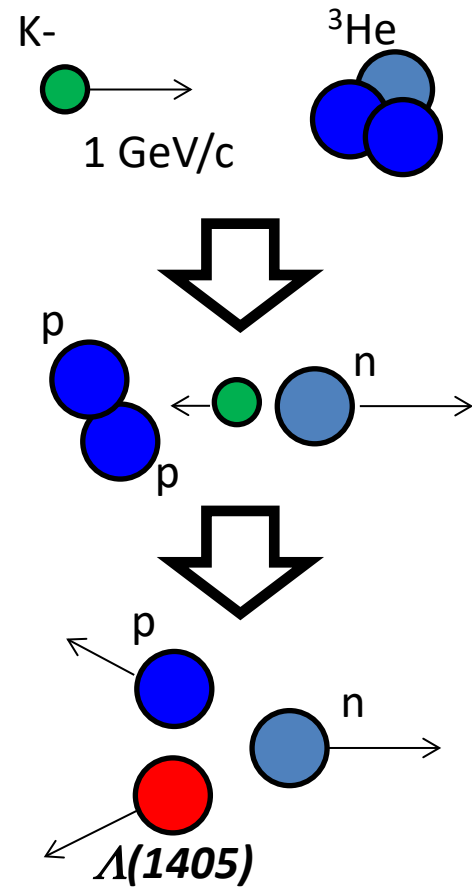
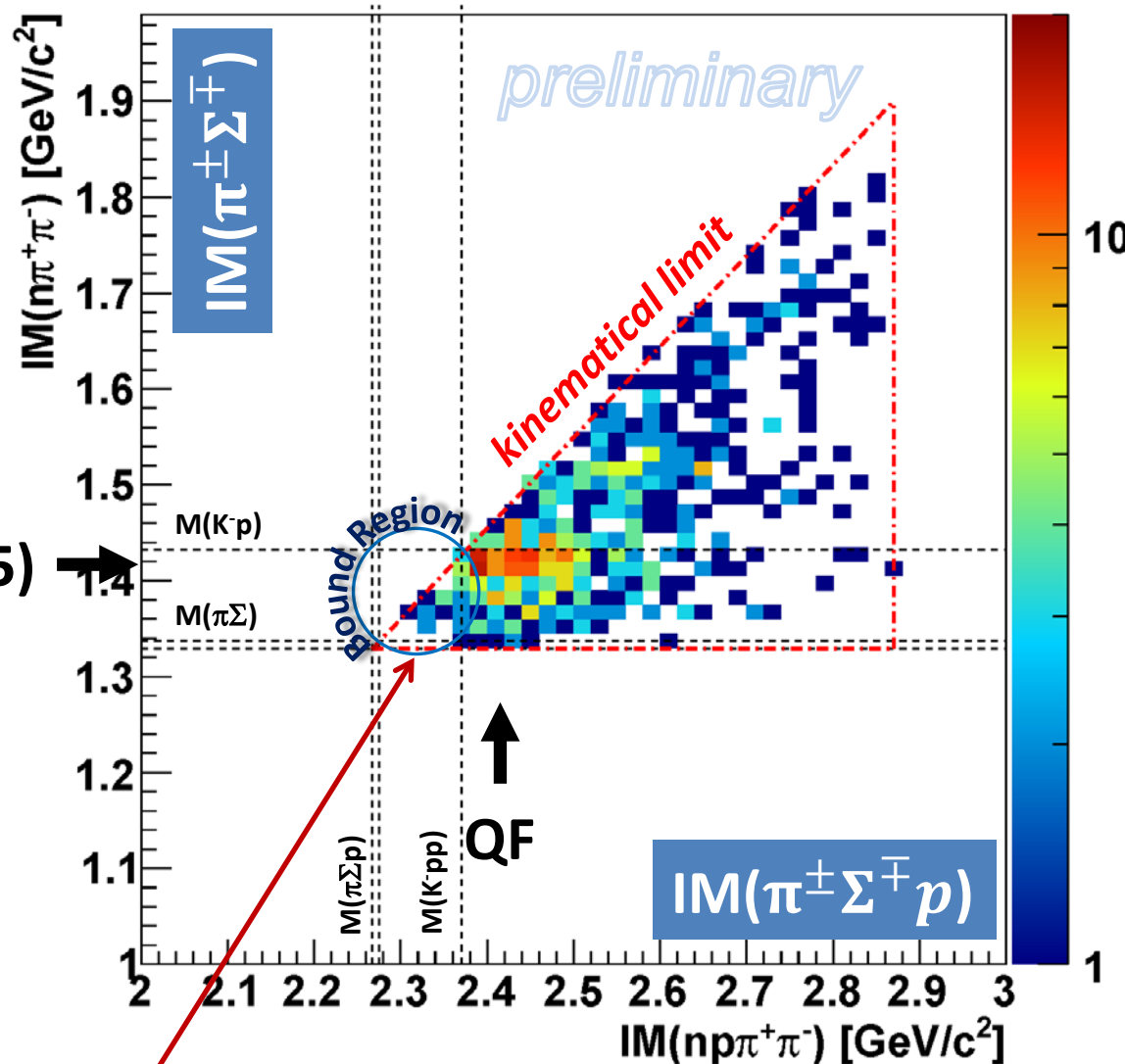




Λp decay is dominant below $M(K^-pp)$



IM($\pi^\pm \Sigma^\mp$) vs. IM($\pi^\pm \Sigma^\mp p$)



Small phase-space of “ K^-pp ” $\rightarrow \pi\Sigma N$

Conclusions

- We have observed a resonance peak below the K^-pp threshold in ${}^3\text{He}(K^-, \Lambda p)n$, “ K^-pp ”

- Binding energy: ~ 50 MeV
- Width: ~ 100 MeV
- S-wave form factor: ~ 400 MeV

← E15 collab., arXiv:1805.12275

- $\Lambda(1405)$ was clearly observed in $\pi^\pm \Sigma^\mp p$ n_{miss} final state

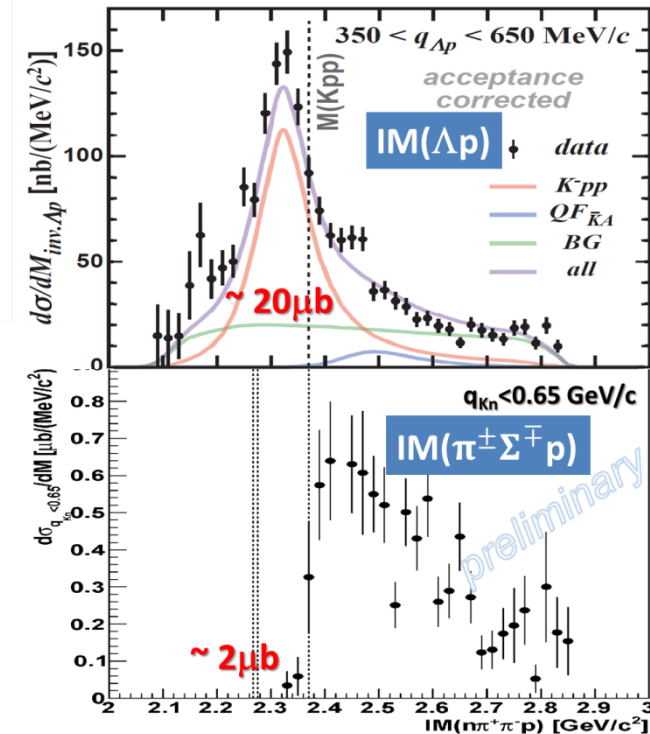
- Large CS of Λ^* compared to “ K^-pp ” formation

← need theoretical feedbacks

- Weak structure below the K^-pp threshold is seen in $\text{IM}(\pi^\pm \Sigma^\mp p)$

- Non-meonic YN decay modes would be dominant

← need further investigation of “ K^-pp ” → $\pi\Sigma N$



What we have to do next























- **More quantitative studies of the “K⁻pp”**
 - J^P
 - Angular distributions are consistent with a J^P=0⁻ assumption in current statistics
 - πΣp decay mode
 - Due to phase-space, or, detector acceptance(?)
- **Series of the kaonic nuclei searches:**
 - “K⁻ppn” via [K⁻ + ⁴He], “K⁻ppnn/K⁻ppppn” via [K⁻ + ⁶Li], etc.
 - “K⁻K⁻pp” via [p^{bar} + ³He annihilation]

**We need a 4π detector system
with γ/n sensitive detectors**

Thank You!

J-PARC E15 Collaboration

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