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Highlights from the STAR experiment Hanna Zbroszczyk

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supported by National Science Centre, Poland

MESON 2018, Kraków, 9th June 2018



Relativistic Heavy Ion Collider (RHIC) Brookhaven National Laboratory (BNL), New York



2 concentric rings of 1740 superconducting magnets
3.8 km circumference

The Solenoidal Tracker At RHIC

- Tracking and PID (full 2π) TPC: $|\eta| < 1$ TOF: $|\eta| < 1$ BEMC: $|\eta| < 1$ EEMC: $1 < \eta < 2$ HFT (2014-2016): $|\eta| < 1$ MTD (2014+): $|\eta| < 0.5$
- MB trigger and event plane reconstruction BBC: $3.3 < |\eta| < 5$ EPD (2018+): $2.1 < |\eta| < 5.1$ FMS: $2.5 < \eta < 4$ VPD: $4.2 < |\eta| < 5.1$ ZDC: $6.5 < |\eta| < 7.5$
- On-going/future upgrades iTPC (2019+): $|\eta| < 1.5$ eTOF (2019+): $-1.6 < \eta < -1$





RHIC Top Energy p+p, p+Al, p+Au, d+Au, ³He+Au, Cu+Cu, Cu+Au, Ru+Ru, Zr+Zr, Au+Au, U+U QCD at high energy density/temperature Properties of QGP, EoS

Beam Energy Scan Au+Au 7.7-62 GeV QCD phase transition Search for critical point Turn-off of QGP signatures

Fixed-Target Program Au+Au =3.0-7.7 GeV High baryon density regime with 420-720 MeV

- 1. Open heavy flavor $\,D^{0}\,v_{_{1}}^{},\,D^{0}\,R_{_{AA}}^{}\,$ and $R_{_{CP}}^{},\,\Lambda_{_{C}}^{}$
- 2. Quarkonium ΥR_{AA}
- 3. Jet modification and high- $p_{\rm T}$ hadrons di-jet imbalance, di-hadron correlation
- 4. Chirality, vorticity and polarization effects Λ polarization, Φ polarization, CME, CMW
- 5. Initial state physics and approach to equilibrium v_2 and v_3 fluctuations
- 6. Collectivity in small systems v_2 in p+Au and d+Au
- 7. Collective dynamics longitudinal decorrelation, identified particle v_1
- 8. High baryon density and astrophysics v_1 from fixed target
- 9. Correlations and fluctuations femtoscopy
- 10. Phase diagram and search for the critical point net Λ and off-diagonal cumulants
- 11. Thermodynamics and hadron chemistry triton, hypertriton mass
- 12. Upgrades BES-II and forward upgrades



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Results



- The moving **spectators** can produce enormously large **electromagnetic field** (eB $\sim 10^{18}$ G at RHIC)

- Due to **early** production of heavy quarks ($\tau_{CQ} \sim 0.1$ fm/c) positive and negative charm quarks (CQs) can get **deflected** by the initial EM force

- Model predicts **opposite** v_1 for charm and anti-charm quarks induced by this initial EM field

- This induced $\boldsymbol{v}_{_1}$ depends on the balance between E and B fields

- The **magnitude** of such induced v_1 for heavy quarks is much **larger** than the light quarks





- This can induce **larger** v_1 in charm quarks than light

flavors

- Magnitude of charm quark v_1 depends on the drag parameter used in this model

- We can probe the longitudinal profile of the initial matter distribution through heavy flavor $v_{_{\rm 1}}$

(v_1 -slope) Charm-Quark >> (v_1 -slope) Light-Quark

- Charm quarks much more sensitive to the initial tilt than the charged hadrons $D^0(\bar{D}^0) v_1$ can be used to constrain drag coefficients in conjunction with v_2 and R_{AA}





Recent hydro model with initial EM field predicts v_1 -**split** between the D and anti-D meson

D meson v_1 greater than the anti-D Predicted difference in v_1 is about 10 times smaller than the average v_1



Significant suppression at **low** $\mathbf{p}_{_{\mathrm{T}}}$ with no strong centrality dependence, Suppression at **high** $\mathbf{p}_{_{\mathrm{T}}}$ decreases towards more peripheral collisions.

Non-prompt D⁰ R_{AA} study has been performed, need better precision measurements to understand mass dependence of energy loss.

STAR data was re-analysed due to
error found durring analysis
→ erratum will be published soon

First evidence_of non-zero directed flow for heavy flavor Both D^0 and D^0 show **negative** v_1 -slope near mid-rapidity

Heavy flavor $v_1 > \text{light flavor } v_1$

Data can be used to probe **initial** matter distribution

Current precision is **not sufficient** to draw conclusion on magnetic field induced charge separation of heavy quarks

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$$E\frac{d^{3}N}{dp^{3}} = \frac{1}{2\pi} \frac{d^{2}N}{p_{T}dp_{T}dy} \left(1 + 2\sum_{n=1}^{\infty} v_{n} \cos\left[n(\phi - \Phi_{R})\right]\right)$$

Q-cumulant method (traditional)

$$\langle 2 \rangle_n = \left\langle e^{in \, (\phi_1 - \phi_2)} \right\rangle$$

$$v_n^4\{4\} = \langle 4 \rangle_{nn} - 2 \langle 2 \rangle_n \langle 2 \rangle_n$$

$$\langle 4 \rangle_{nm} = \left\langle e^{in (\phi_1 - \phi_2) + im(\phi_3 - \phi_4)} \right\rangle$$

$$NSC(n,m) = \frac{\langle 4 \rangle_{nm} - \langle 2 \rangle_n \langle 2 \rangle_m}{\langle 2 \rangle_n^{Sub} \langle 2 \rangle_m^{Sub}}$$

$$\Phi \text{ - azimuthal angle}$$

Two-subevent method

$$\langle 2 \rangle_n^{Sub} = \left\langle e^{in \left(\phi_A - \phi_B\right)} \right\rangle \qquad v_n^2 \{2\} = c_n \{2\} = \langle 2 \rangle_n^{Sub}$$

✓ Short-range non-flow contribution in v_2 {2} is suppressed by $|\Delta \eta| > 0.7$



$$v_n^4\{4\} = 2\langle v_n^2 \rangle^2 - \langle v_n^4 \rangle$$

$$\left[\frac{v_n\{4\}}{v_n\{2\}}\right]^4 = 2 - \frac{\langle v_n^4 \rangle}{\langle v_n^2 \rangle^2}$$

Sensitive to flow fluctuations



Strong dependence of v_2^{2} and v_2^{4} on collision centrality more significant for higher collision energies

Weak dependence of $v_2\{2\}/v_2\{4\}$ on collision centrality



Weak dependence of $v_2\{2\}$, $v_2\{4\}$ and $v_2\{2\}/v_2\{4\}$ on transverse momentum



Significant dependence of $v_2\{2\}$, $v_2\{4\}$ and $v_2\{2\}/v_2\{4\}$ on collision centrality for different A+A collisions

Anisotropic flow magnitude is sensitive to: - initial-state spatial anisotropy

- flow fluctuations and correlations
- viscous attenuation ($\propto \eta/s$ (T))



on collision centrality for various systems.

Are dynamical final-state fluctuations significantly less than the initial-state fluctuations?

Strong dependence of v_2 {2}, v_2 {4} on collision centrality, collision energy, transverse momentum

Weak dependence of $v_2\{4\}/v_2\{2\}$ and $v_2\{2\}/\varepsilon_2\{2\}$ (elliptic flow fluctuations) on the size of colliding system and: collision centrality, collision energy, transverse momentum

Flow flucuations are dominated by the fluctuations of the **initial state eccentricity**

Similar viscous coefficient for different colliding systems

- Search for turn-off QGP

signatures



cover $\sqrt{s_{_{NN}}}$ from 3.0 GeV to

7.7 GeV

21



Detector efficiency Detector acceptance (each rapidity window) Energy loss



Negavtive pions spectra **are consistent** with AGS results.



Directed flow for pions and protons with fit describing midrapidity region.

Directed flow of protons **agrees** with AGS results.



Directed flow for Λ and $K^0_{\ s}$ particles and their fits describing mid-rapidity region.



HBT radii for pions are **consistent** with AGS results.

- **STAR is ready** to operate with the Fixed Target mode
- Spectra and particle yields agree with AGS results
- **Proton directed flow** v_1 agrees with AGS results
- HBT radii agree with AGS results

High-baryon density regime will be accessible with the Fix Target mode in STAR!

Single- and two- particle distributions

$$P_{1}(p) = E \frac{dN}{d^{3}p} = \int d^{4}x S(x, p)$$

$$S(x,p) - \text{emission function: the distribution of source density probability of finding particle with x and p
$$P_{2}(p_{1}, p_{2}) = E_{1}E_{2}\frac{dN}{d^{3}p_{1}d^{3}p_{2}} = \int d^{4}x_{1}S(x_{1}, p_{1})d^{4}x_{2}S(x_{2}, p_{2})\Phi(x_{2}, p_{2}|x_{1}, p_{1})$$

$$P_{2}(p_{1}, p_{2}) = E_{1}E_{2}\frac{dN}{d^{3}p_{1}d^{3}p_{2}} = \int d^{4}x_{1}S(x_{1}, p_{1})d^{4}x_{2}S(x_{2}, p_{2})\Phi(x_{2}, p_{2}|x_{1}, p_{1})$$

$$P_{1}(p_{1}, p_{2}) = \frac{P_{2}(p_{1}, p_{2})}{P_{1}(p_{1})P_{1}(p_{2})}$$

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centrality	$R_{in u}p-p$ [fm]	$R_{inv}\overline{p}-\overline{p}$ [fm]	$R_{inv} p - \overline{p}$ [fm]	No significant difference
0-10%	4 . 00 \pm 0.15 \pm 0.02	3 . 83 \pm 0.20 \pm 0.03	3 . 39 \pm 0.12 \pm 0.14	between proton-proton
10-30%	3 . 61 \pm 0.13 \pm 0.17	3 . 68 \pm 0.15 \pm 0.11	2 . 69 \pm 0.10 \pm 0.12	correlation functions
30-70%	$2.72 \pm 0.07 \pm 0.07$	$2.95 \pm 0.11 \pm 0.08$	2 . 56 \pm 0.09 \pm 0.12	









- Clear centrality dependence of source size at BES energies
- Visible energy dependence of source size at BES energies
- No visible difference between proton-proton and antiproton-antiproton correlation functions at $\sqrt{s_{_{NN}}} = 39$ GeV
- Correlation functions contaminated by residual correlations residual correction required





Hyperon-Nucleon:

- play an important role in neutron star and QCD theory

- measurements of masses of hypertriton and anti-hypertriton provide insight into H-N interactions and the CPT symmetry

- measurements sensitive to the temperature and nucleon phase-space of the system freezeout.

- excellent tool to explore the QCD properties
- R. O. Gomes, V. Dexheimer, S. Schramm, and C. A. Z. Vasconsellos, The Astrophys. J. 808, 8 (2015).
- [2] L. L. Lopes and D. P. Menezes, Phys. Rev. C 89, 025805 (2014).
- [3] J. Antoniadis et al., Science 340, 448 (2013).
- [4] László P. Csernai, Joseph I. Kapusta, Phys. Reps. 131, 223 (1986).
- [5] A. Z. Mekjian, Phys. Rev. C 17, 1051 (1978).
- [6] Kaijia Sun et al., Phys. Lett. B 774, 103 (2017).





 $^{3}_{\Lambda}$ H has many decay channels:

- ✓ Non-meson decay channels: $^{3}_{\Lambda}H \rightarrow d + n$ $^{3}_{\Lambda}H \rightarrow p + n + n$

Good PID of charged particles in STAR detector.

Reconstructing ${}^{3}_{\Lambda}H (\frac{3}{\Lambda}\overline{H})$ through: ${}^{3}_{\Lambda}H \rightarrow {}^{3}He + \pi^{-}$ ${}^{3}_{\Lambda}H \rightarrow d + p + \pi^{-}$







Measurements of the massover-charge ratio differences between light nuclei and antinuclei.



Conclusions & Summary

Summary

1. Open heavy flavor - $D^0 v_1$, $D^0 R_{AA}$ and R_{CP} , Λ_C

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Upgrades

iTPC Upgrade:

STAR

- Improves tracking and acceptance at low pT and extra y acceptance
- Ready in 2019

STAR Note 0644: Technical Design Report for the iTPC Upgrade eTOF Upgrade:

- Improves PID and acceptance
- Ready in 2019

arXiv:1609.05102v1 [nucl-ex]

inner TPC upgrade

endcap TOF

EPD Upgrade:

Event Plane Detector

- Improves event plane resolution and centrality definition
- Taking data in 2018 run

STAR Note 0666: An Event Plane Detector for STAR

STAR Note 0696: STAR Collaboration Beam Use Request for Run 19+ (Scenario 1)

Single Beam Energy (GeV/nucleon)	√ <i>S</i> _{NN} (GeV)	Run Year	Run Time	Species	Min-Bias Events Number
5.75	3.5 (FXT)	2020	2 days	Au+Au	100M
7.3	3.9 (FXT)	2019	2 days	Au+Au	100M
9.8	4.5 (FXT)	2019	2 days	Au+Au	100M
13.5	5.2 (FXT)	2020	2 days	Au+Au	100M
19.5	6.2 (FXT)	2020	2 days	Au+Au	100M
31.2	7.7 (FXT)	2019	2 days	Au+Au	100M

- iTPC & eTOF upgrades will be available
- Need 100M events at each energy to match sensitivity of BES-II:
 2 days per energy (3.5 GeV 7.7 GeV)
- Data rate is DAQ limited
- Data at 7.7 GeV will provide an overlap energy with collider mode

FXT in Run 18

Trigger commissioning occurring now

1 Billion events at 7.2 GeV

100 Million events at 3.0 GeV

EPD ready and available for flow analyses

Can obtain fluctuation measurement at energies below BES-I

Thank you!

Backup slides



Strong dependence on collision energy

Weak dependence on collision centrality



Low multiplicity subtraction scaled by short-range near-side ($|\Delta \eta| < 0.5$) jet yield

$$V_{n,n}^{HM}(subtracted) = V_{n,n}^{HM} - V_{n,n}^{LM} \times \frac{N_{asso.}^{LM}}{N_{asso.}^{HM}} \times \frac{Y_{jet,near-side}^{HM}}{Y_{jet,near-side}^{LM}}$$

ATLAS:PRC90(2014)044906 CMS:PLB765(2017)193 STAR: PLB743(2015)333

Short-range near-side jet modification = long-range away-side jet modification

Template fit



$$\begin{split} Y_{templ.}(\Delta \phi) &= \mathsf{F} \times Y_{LM}(\Delta \phi) + Y_{ridge}(\Delta \phi) \\ \text{where} \\ Y_{ridge}(\Delta \phi) &= \mathsf{G} \times (1 + 2 \times \sum_{n=2}^{4} V_{n,n} \times \cos(n\Delta \phi)) \end{split}$$

ATLAS:PRL(116)172301

A new method by ATLAS Collaboration away-side jet shape can be measured in Low Multiplicity (LM) events scaled by "F" parameter (due to jet modification)

 v_2 without subtraction is **larger** than that with subtraction for both methods.

The subtraction of non-flow contributions are very **important** for STAR results are comparable with PHENIX results, except at high pT.

At lowet $p_T v_2$ from Low Multiplicity subtraction is **35% lower** than from template fit

At intermediate p_{T} they **agree** with each other

STAR results are **comparable** with PHENIX ones.



 v_2 in p+Au collisions without subtraction is **larger** than v_2 in d+Au collisions that with subtraction for both methods.

v₂ in p+Au collisions from Low Multiplicity subtraction is **lower** than from template fit.

STAR results are **comparable** with PHENIX results, except at high pT. The STAR data is clearly lower than PHENIX for p_T >1.5 GeV/c





Large **difference** between subtraction method and template fit

v₂ from subtraction method is **negative** at lower collision energies (different kinematics between near-side and away-side jet-like correlations?)

 $\mathbf{v}_{_2}$ from template fit **increases** with collision centrality



Large **difference** between v_2 from two methods has been observed at low energy \rightarrow large uncertainties in the non-flow subtraction in small systems.

We do see **similar** \mathbf{v}_2 between p+Au and d+Au collisions for same multiplicity $\rightarrow \mathbf{v}_2$ is not only driven by initial geometry.

The integral v_2 extracted by a template fit shows an **universal** trend as a function of $\langle dN/d\eta \rangle$ for different small systems at different energies \rightarrow multiplicity plays an important role in small systems.



Directed flow for identified particles **agrees** with AGS results.

Triton from Au+Au Collision

