

15th International Workshop on Meson Physics
Kraków, Poland, June 7th - 12th 2018

Scalar Dipole Dynamical Polarizabilities

from

**proton Real Compton
Scattering data**



UNIVERSITÀ
DI PAVIA

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&
INFN, Pavia (Italy)



Stefano Sconfiatti

1 / 24

Outline

- ✓ RCS: theoretical framework
- ✓ Static polarizabilities
- ✓ Low Energy expansion
- ✓ Dynamical polarizabilities

Outline

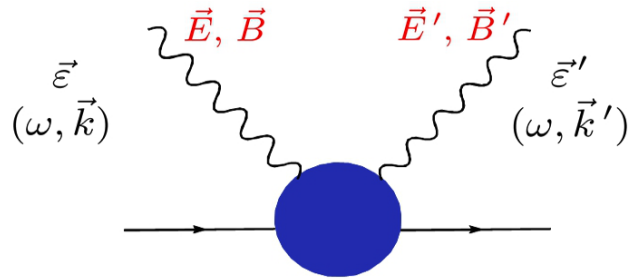
- ✓ RCS: theoretical framework
- ✓ Static polarizabilities
- ✓ Low Energy expansion
- ✓ Dynamical polarizabilities

- ✓ Data set inconsistency
- ✓ Very high correlations
- ✓ Too much parameters
- ✓ New approach: simplex + bootstrap

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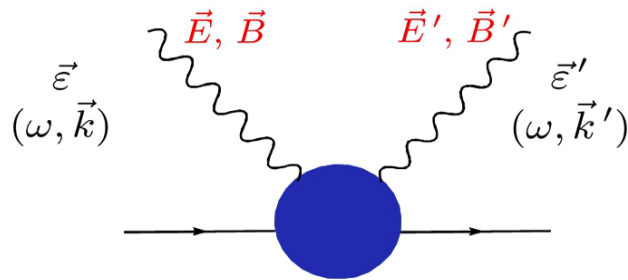
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- ✓ Static polarizabilities: cross check
 - ✓ Systematical errors
 - ✓ Dynamical polarizabilities from data
 - ✓ Conclusions and future perspectives

RCS amplitudes and Dispersion Relations



$$A_i(\nu, t) = A_i^s(\nu, 0) + A_i^t(0, t) + A_i(0, 0)$$

RCS amplitudes and Dispersion Relations



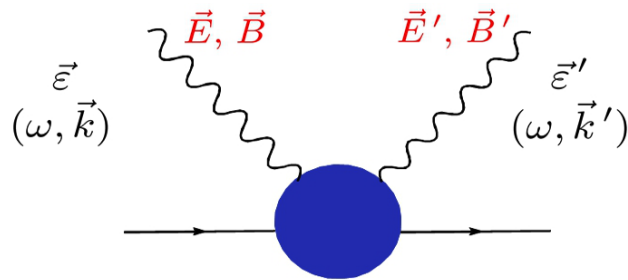
$A_i(0,0) = a_i$ **→** Static polarizabilities

$$A_i(\nu, t) = A_i^s(\nu, 0) + A_i^t(0, t) + A_i(0, 0)$$

Subtracted Dispersion Relations (s-channel)

$$A_i^s(\nu, 0) = \frac{2}{\pi} \nu^2 \text{P} \int_{\nu_{thr}}^{\infty} \text{Im}_s A_i(\nu', t) \frac{d\nu'}{\nu'(\nu'^2 - \nu^2)}$$

RCS amplitudes and Dispersion Relations



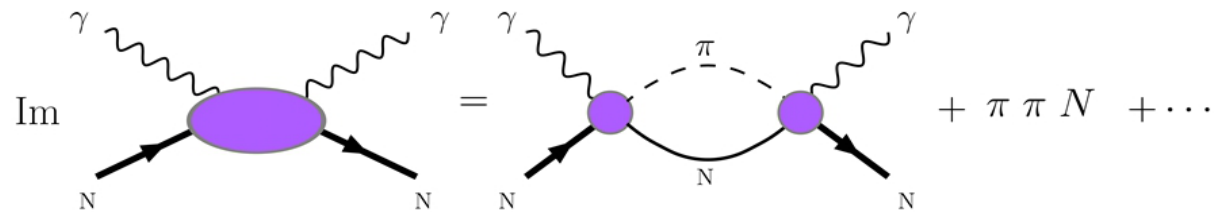
$$A_i(\mathbf{0}, \mathbf{0}) = a_i \longrightarrow \text{Static polarizabilities}$$

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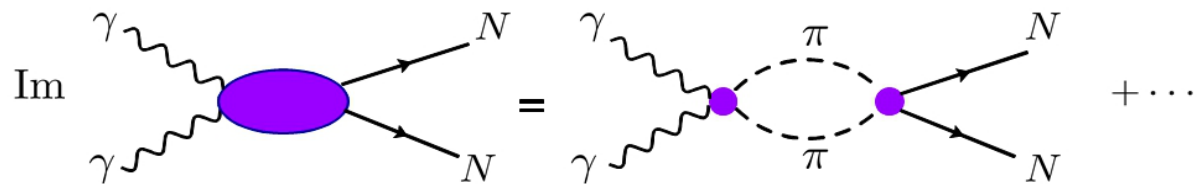
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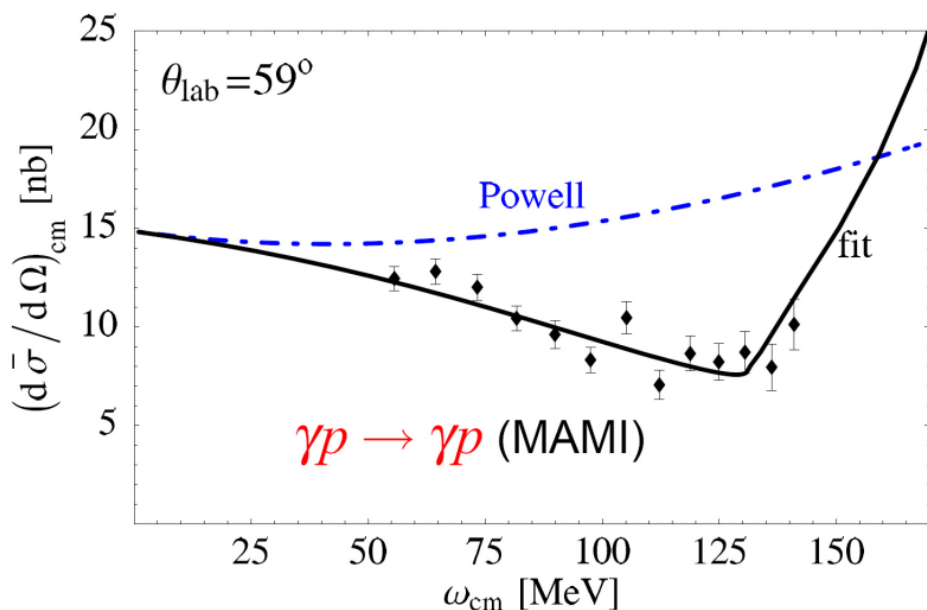
s-CHANNEL



t-CHANNEL



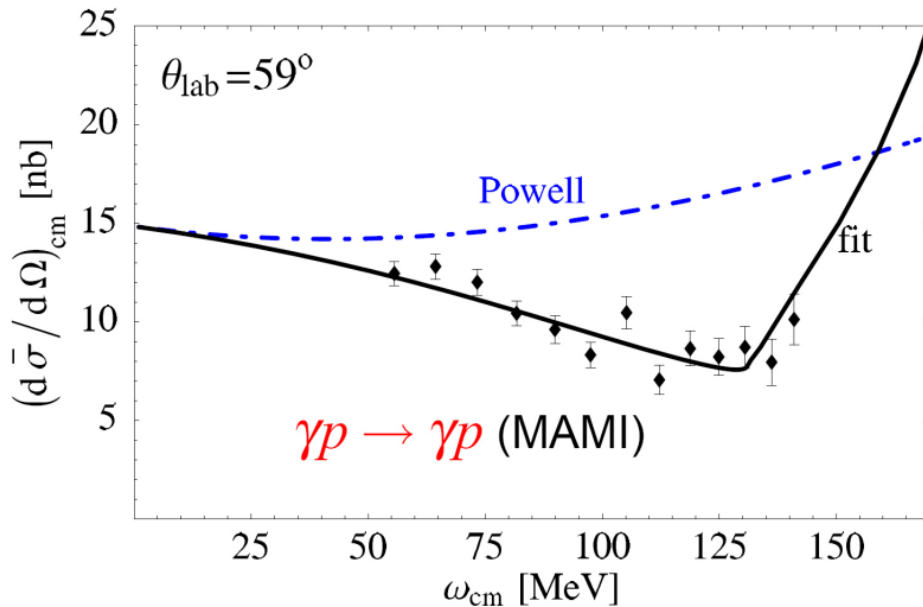
Static polarizabilities



Powell cross section: point-like nucleon with anomalous magnetic moment

Static polarizabilities: response of the internal nucleon degrees of freedom to a **static** electric and magnetic field

Static polarizabilities



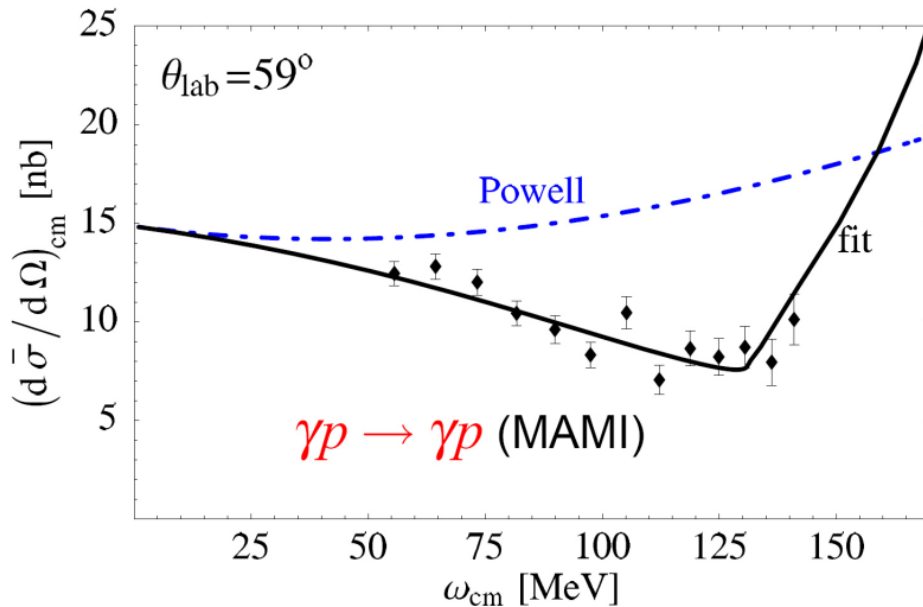
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Static polarizabilities: response of the internal nucleon degrees of freedom to a **static** electric and magnetic field

$$H_{\text{eff}}^{\text{pol}} = -2\pi \left\{ \omega^2 \left[\alpha_{E1} \vec{E}^2 + \beta_{M1} \vec{B}^2 \right] \right\}$$

spin-independent dipole

Static polarizabilities



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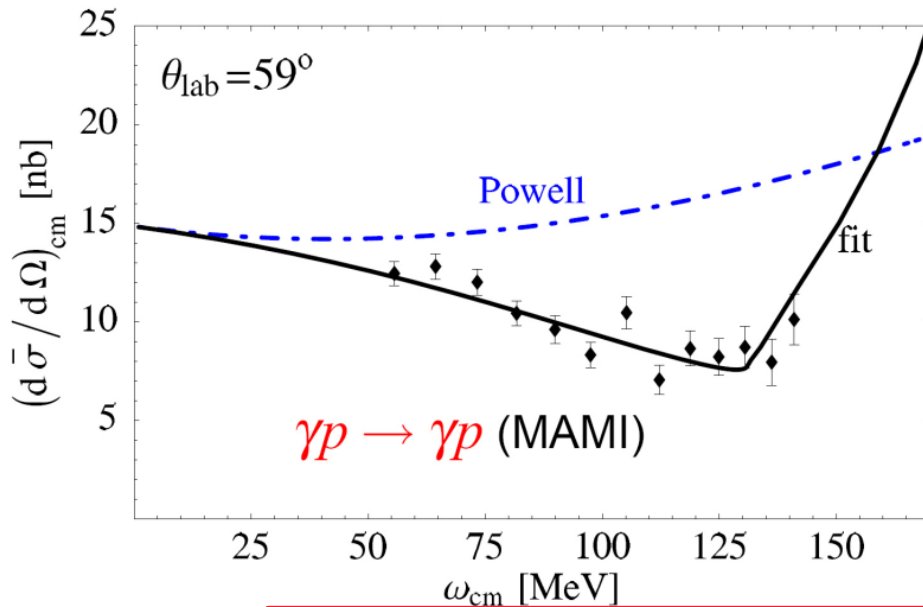
$$+ \omega^3 \left[\gamma_{E1E1} \vec{\sigma} \cdot (\vec{E} \times \dot{\vec{E}}) + \gamma_{M1M1} \vec{\sigma} \cdot (\vec{B} \times \dot{\vec{B}}) \right.$$

spin-dependent dipole

$$\left. -2 \gamma_{M1E2} \sigma_i B_j E_{ij} + 2 \gamma_{E1M2} \sigma_i E_j B_{ij} \right] + \mathcal{O}(\omega^3) \left. \right\}$$

spin-dependent dipole-quadrupole

Static polarizabilities



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Static polarizabilities: response of the internal nucleon degrees of freedom to a **static** electric and magnetic field

$$\begin{aligned}
 H_{\text{eff}}^{\text{pol}} = & \boxed{-2\pi \left\{ \omega^2 \left[\alpha_{E1} \vec{E}^2 + \beta_{M1} \vec{B}^2 \right] \right.} && \text{spin-independent dipole} \\
 & + \omega^3 \left[\gamma_{E1E1} \vec{\sigma} \cdot (\vec{E} \times \dot{\vec{E}}) + \gamma_{M1M1} \vec{\sigma} \cdot (\vec{B} \times \dot{\vec{B}}) \right. && \text{spin-dependent dipole} \\
 & \left. \left. - 2 \gamma_{M1E2} \sigma_i B_j E_{ij} + 2 \gamma_{E1M2} \sigma_i E_j B_{ij} \right] + \mathcal{O}(\omega^3) \right\} && \text{spin-dependent dipole-quadrupole}
 \end{aligned}$$

Low Energy Expansion (LEX) (I)

$$A_i(\nu, t) = A_i(\nu, t)|_{(0,0)} + \left. \frac{\partial A_i(\nu, t)}{\partial \nu^2} \right|_{(0,0)} \nu^2 + \left. \frac{\partial A_i(\nu, t)}{\partial t} \right|_{(0,0)} t + \frac{1}{2} \left(\left. \frac{\partial^2 A_i(\nu, t)}{\partial \nu^4} \right|_{(0,0)} \nu^4 + \left. \frac{\partial^2 A_i(\nu, t)}{\partial t^2} \right|_{(0,0)} t^2 + 2 \left. \frac{\partial^2 A_i(\nu, t)}{\partial \nu^2 \partial t} \right|_{(0,0)} \nu^2 t \right)$$

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ν and t as independent variables

Lorentz invariant amplitudes

Need to choose a ref-frame: CM

$R_i(\mathbf{A}_i)$

(ready for multipole expansion)

Multipole expansion and DYNAMICAL polarizabilities

$$R_1 = \sum_{l \geq 1} \{ [(l+1)f_{EE}^{l+} + lf_{EE}^{l-}](lP'_l + P''_{l-1}) - [(l+1)f_{MM}^{l+} + lf_{MM}^{l-}]P''_l \}$$

$$R_2 = \sum_{l \geq 1} \{ [(l+1)f_{MM}^{l+} + lf_{MM}^{l-}](lP'_l + P''_{l-1}) - [(l+1)f_{EE}^{l+} + lf_{EE}^{l-}]P''_l \}$$

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DYNAMICAL POLARIZABILITIES

$$\alpha_{El} = a(l) \frac{(l+1)f_{EE}^{l+} + lf_{EE}^{l-}}{\omega^{2l}} \quad \beta_{Ml} = a(l) \frac{(l+1)f_{MM}^{l+} + lf_{MM}^{l-}}{\omega^{2l}}$$

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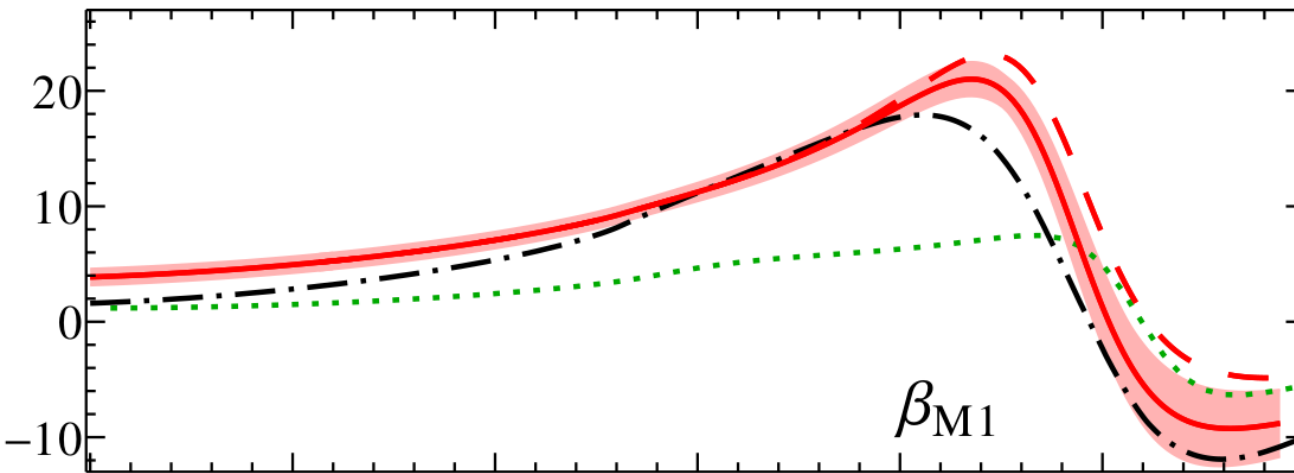
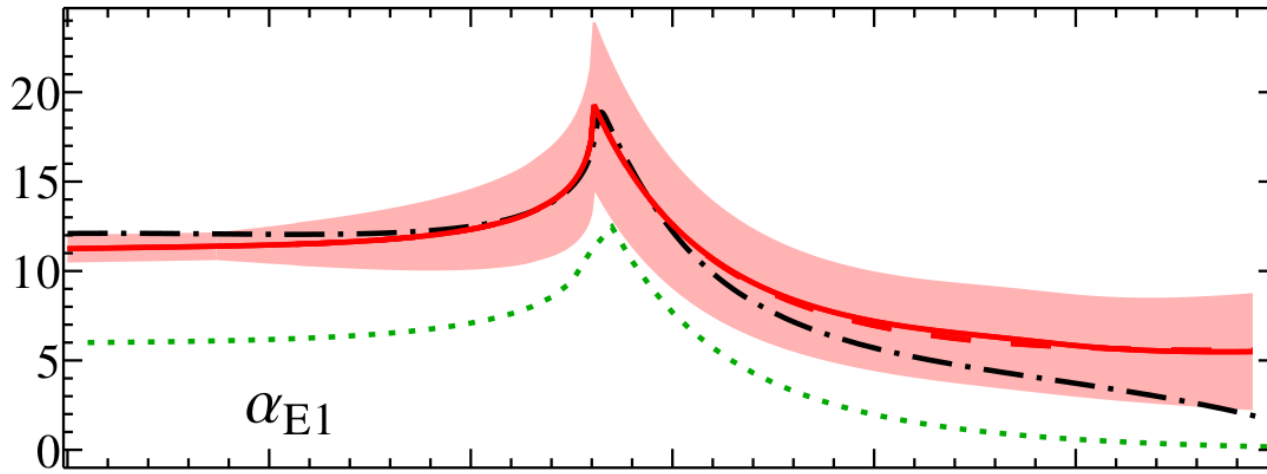
DIPOLE DYNAMICAL POLARIZABILITIES

(DDPs)

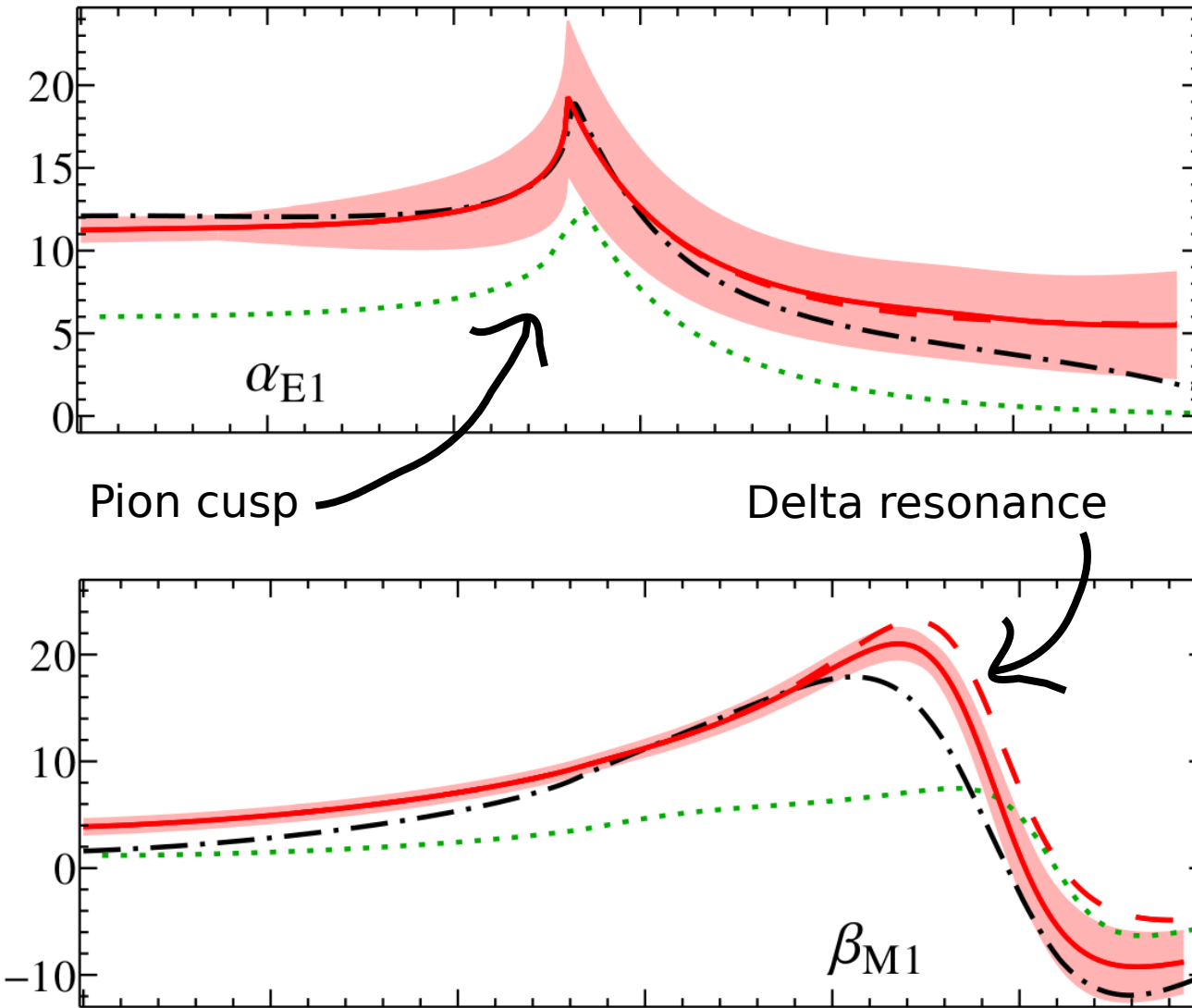
$$\alpha_{E1}(\omega)$$

$$\beta_{M1}(\omega)$$

DDPs: physical meaning



DDPs: physical meaning



DDPs:
response of the
internal nucleon
degrees of freedom to
an electric and
magnetic field with
and explicit
dependence on
energy

Low Energy Expansion (LEX) (II)

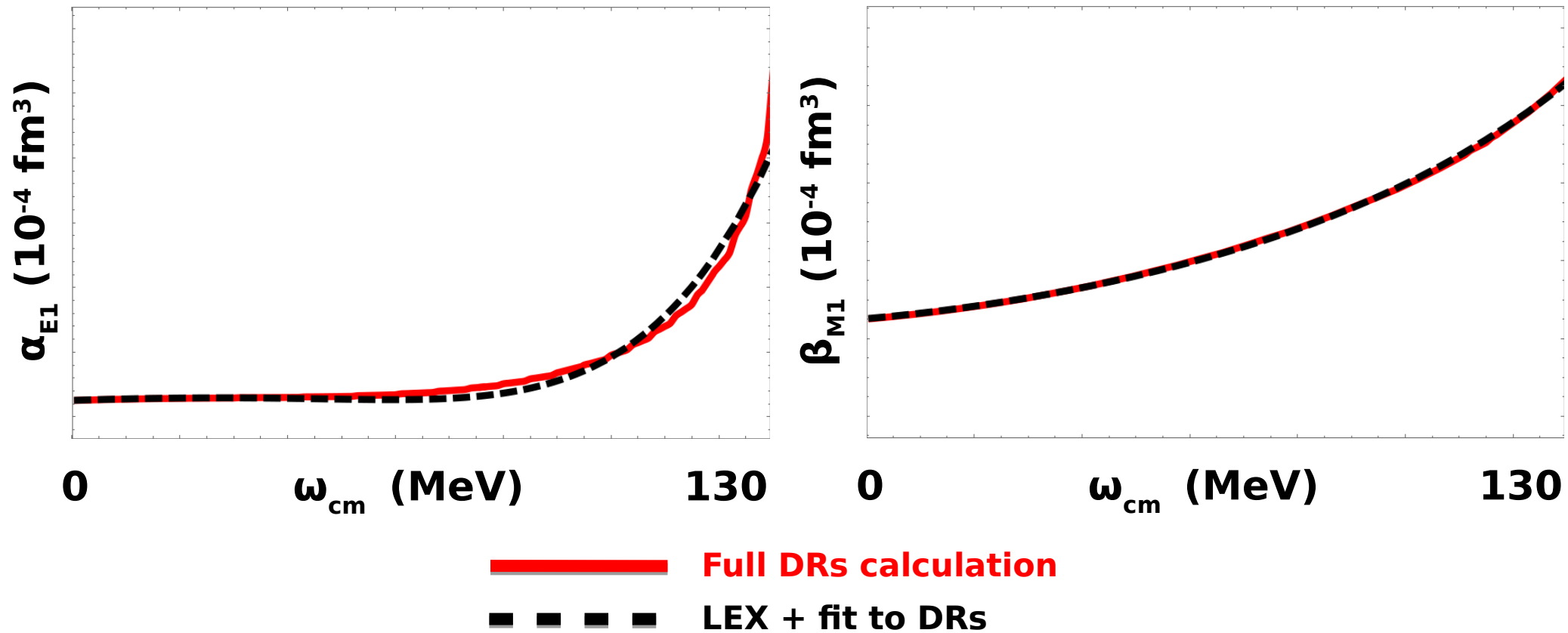
$$\alpha_{E1}(\omega) = \alpha_{E10} + \alpha_{E11} \omega + \alpha_{E12} \omega^2 + \alpha_{E13} \omega^3 + \alpha_{E14} \omega^4 + \alpha_{E15} \omega^5$$

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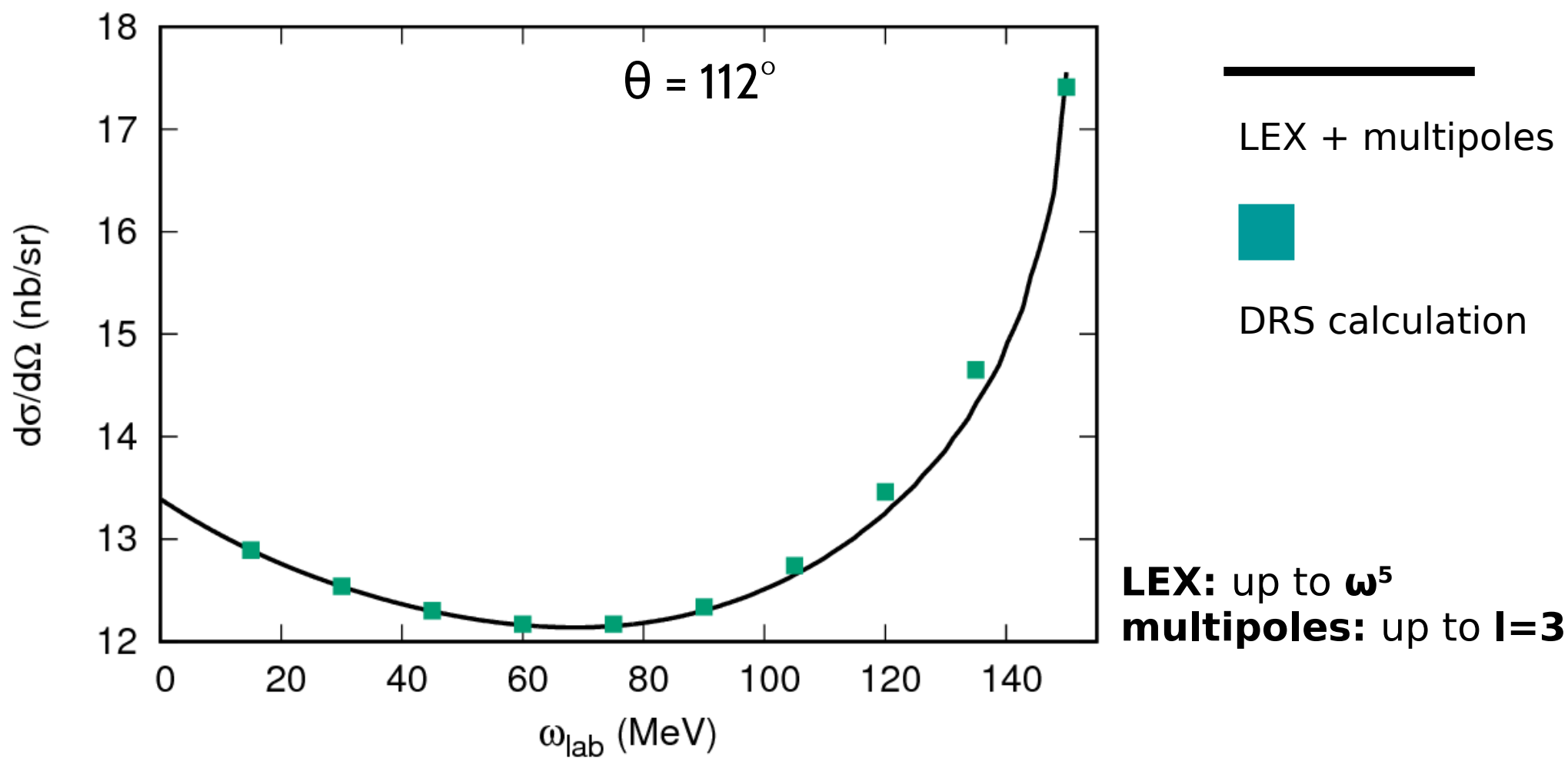
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DRs vs (LEX + multipoles)



The GOAL

Extract scalar
**Dipole Dynamical
Polarizabilities**
(DDPs)
from RCS data

Complications

Gradient method to find the χ^2 minimum

VERY **high correlations** between parameters!

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```
MINUIT WARNING IN HESSE  
===== MATRIX FORCED POS-DEF BY ADDING  
0.13727E-01 TO DIAGONAL.
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Gradient method to find the χ^2 minimum

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VERY **low sensitivity** of the data to dynamical polarizabilities

NO WAY to find the “right” minimum and to define “right” errors on fit parameters

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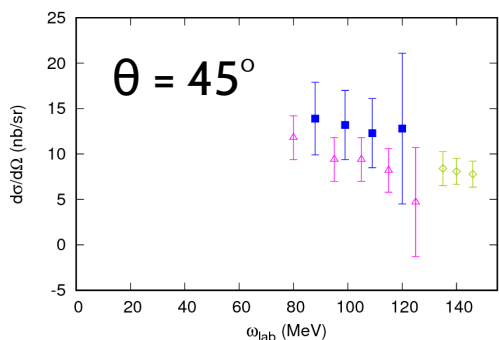
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Combination of **SIMPLEX** method and **BOOTSTRAP** technique

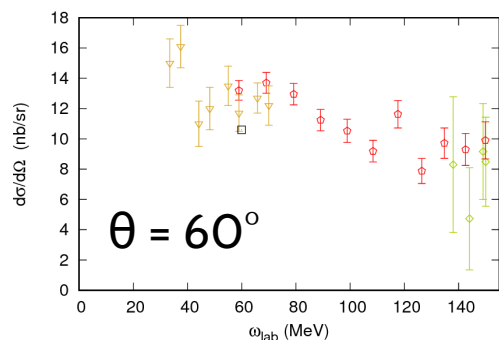
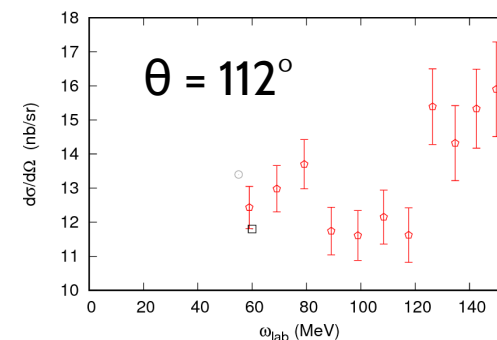
The DATA set

***Half of the Spartans
that King Leonidas led
to the Battle of
Thermopylae...***

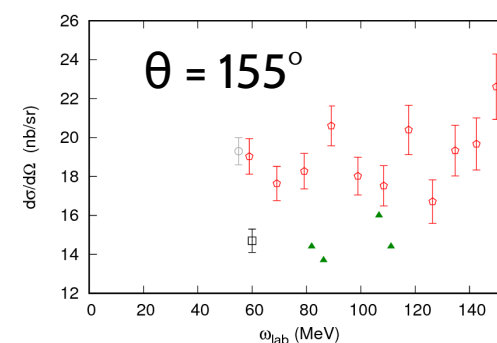
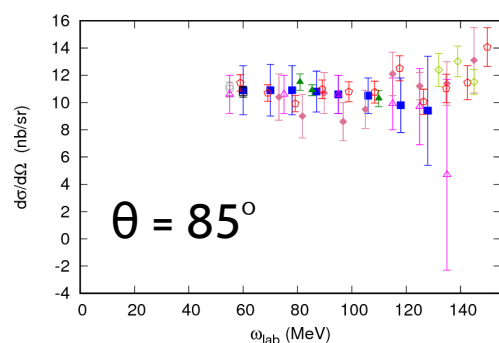
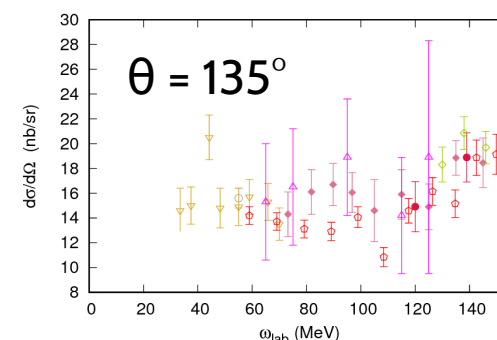
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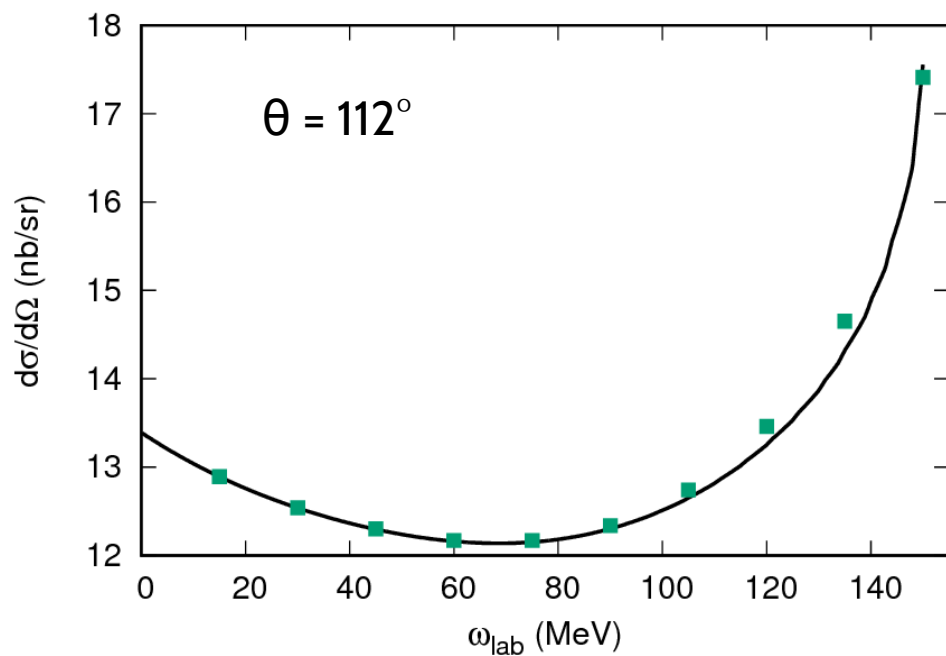
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Symbol	Set	Ref
	GOL 60	Goldansky et al.
	OdL 01	Olmos de León
	HAL 93	Hallin et al.
	HYM 59	Hyman et al.
	PUG 67	Pugh et al.
	FED 91	Federspiel et al.
	BER 61	Bernardini et al.
	BAR 74	Baranov et al.
	OXL 58	Oxley
	MAC 95	MacGibbon et al.

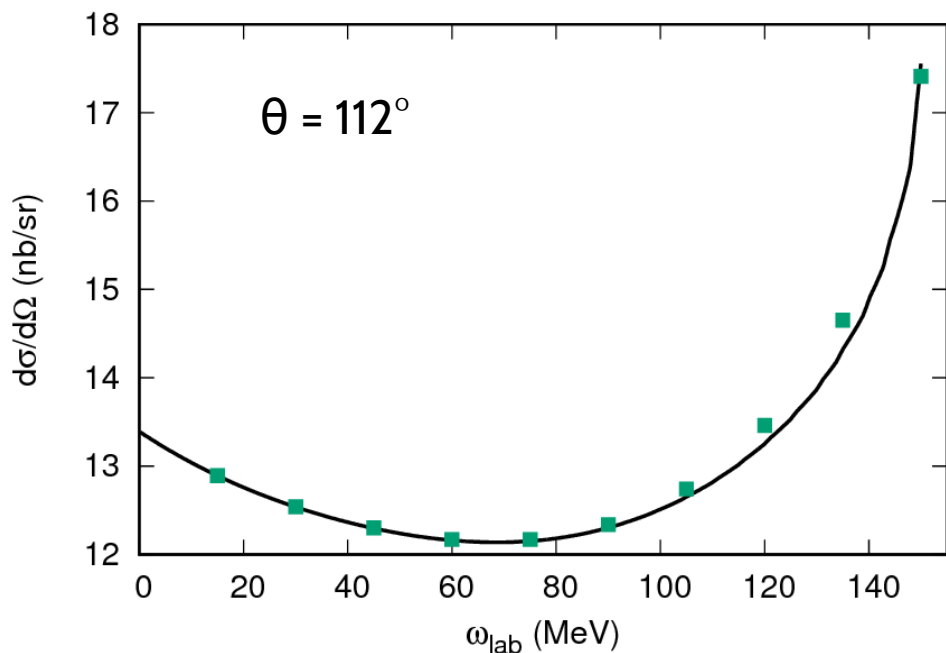


DRs vs (*LEX* + multipoles)



FIRST cross check:
comparison with
LEX + multipoles
and ***DRs***

DRs vs (*LEX* + multipoles)



FIRST cross check:
comparison with
LEX + multipoles
and *DRs*

$\alpha_{E1} (10^{-4} \text{ fm}^3)$

$\beta_{M1} (10^{-4} \text{ fm}^3)$

DRs

11.9 ± 0.2

1.9 ± 0.2

LEX + MULTIPOLES

11.8 ± 0.2

2.0 ± 0.2

Bootstrap sampling and systematics

$$S_{i,exp}^{boot} = S_{i,exp} \pm \gamma \sigma_{i,exp}$$

Gaussian distributed

Bootstrap sampling and systematics

$$\mathbf{S}_{i,exp}^{boot} = \mathbf{S}_{i,exp} \pm \gamma \boldsymbol{\sigma}_{i,exp} \quad \text{Gaussian distributed}$$

How can we include systematical errors?

$$\chi_{mod}^2 = \sum_{i=1}^{N_{tot}} \left[\frac{\mathcal{N} \mathcal{S}_{i,exp} - \mathcal{S}_{i,theory}}{\mathcal{N} \sigma_{i,exp}} \right]^2 + \left(\frac{\mathcal{N} - 1}{\sigma_{i,sys}} \right)^2$$

...one normalization factor per data set is needed!

Bootstrap sampling and systematics

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$$\mathbf{S}_{i,exp}^{boot} = \xi_i [\mathbf{S}_{i,exp} \pm \gamma \boldsymbol{\sigma}_{i,exp}]$$

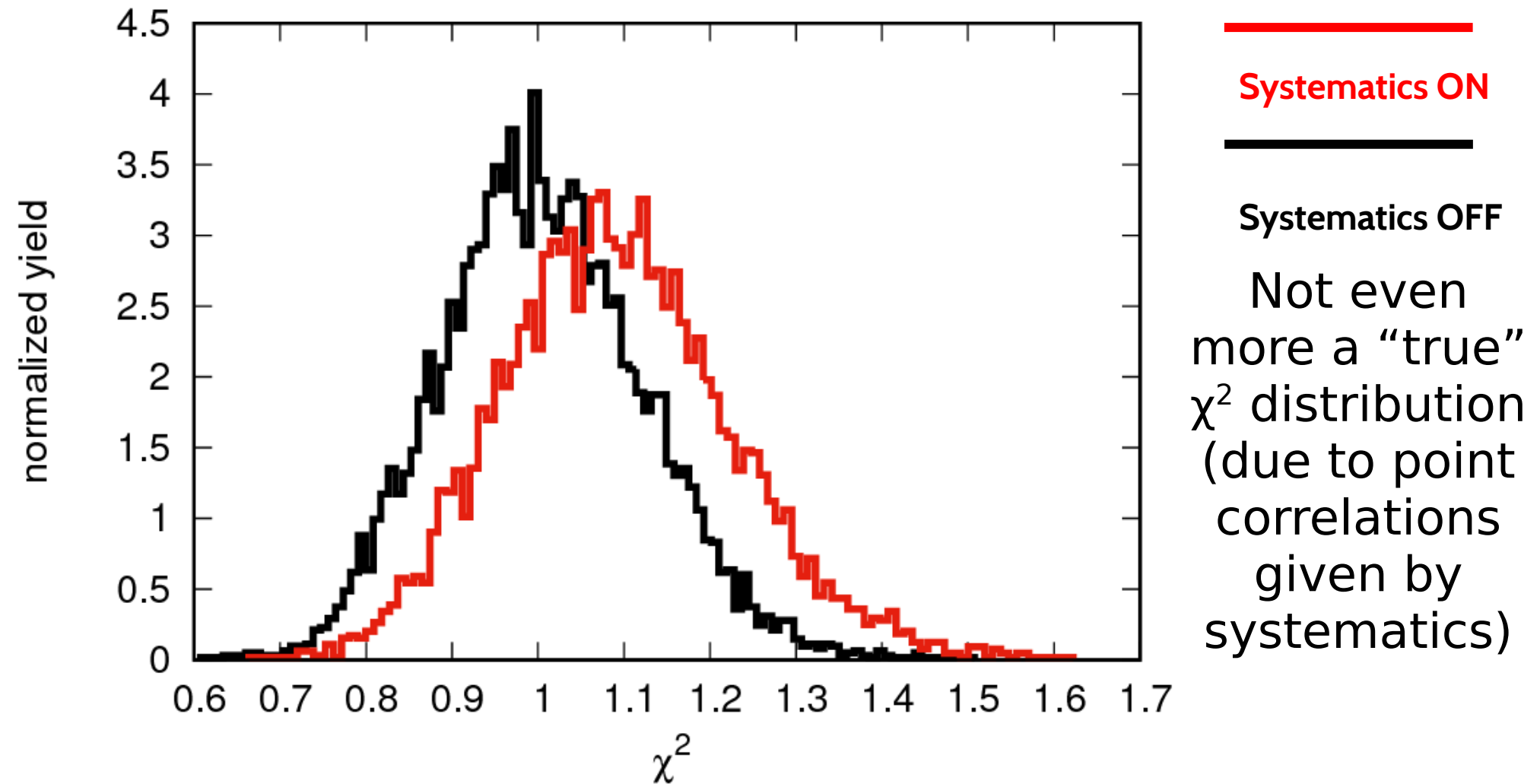
At every bootstrap cycle the systematical errors for each set can vary independently!

Bootstrap vs Gradient: systematics ON

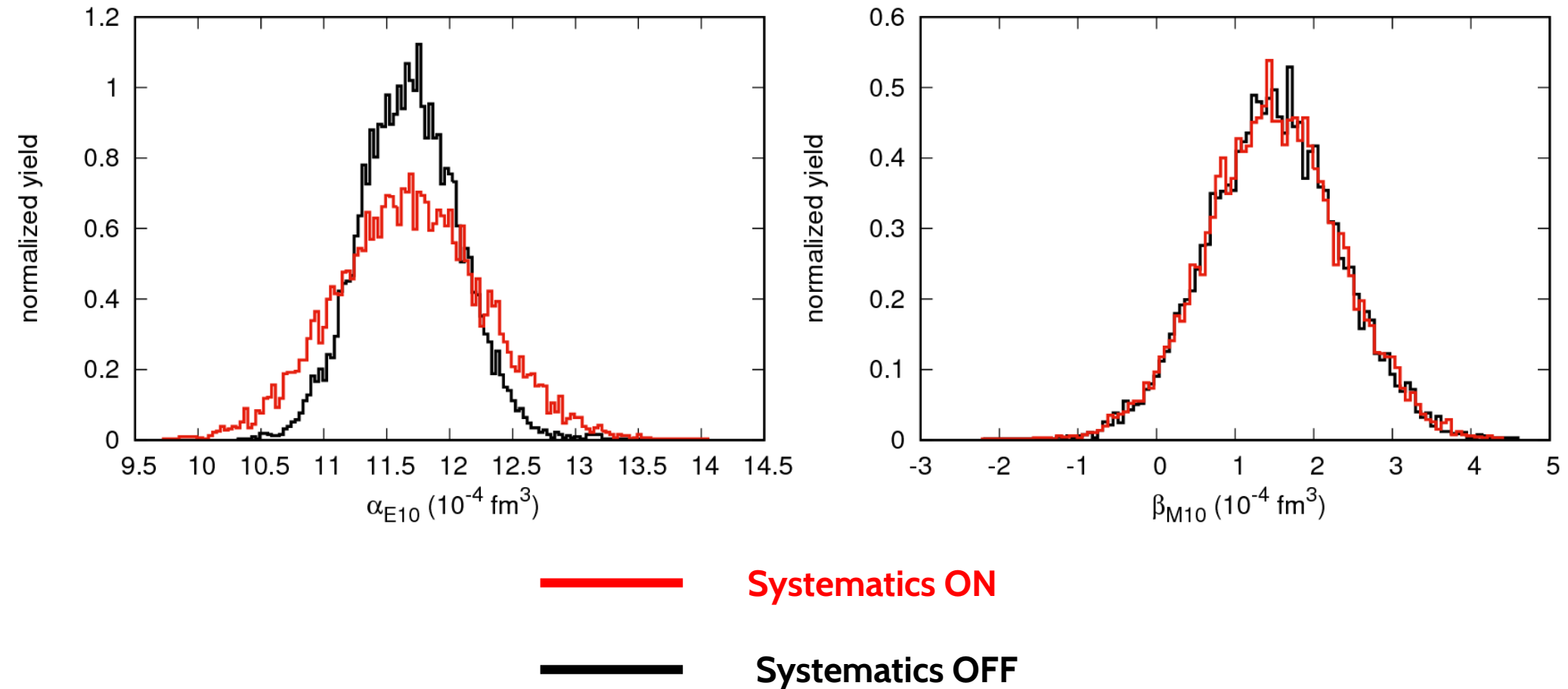
	α_{E1}	β_{M1}
BOOTSTRAP	11.8 \pm 0.2	2.0 \pm 0.2
LEX + MULTIPOLES	11.8 \pm 0.2	2.0 \pm 0.2
BOOTSTRAP SYS ON	11.8 \pm 0.3	2.0 \pm 0.3

Systematical errors enlarge the error band of polarizabilities!

“ χ^2 ” probability distribution in bootstrap framework (static pol.)



The effect of systematics (static spin-independent polarizabilities)



Expected Gaussian shape + systematics enlarging

Fit conditions

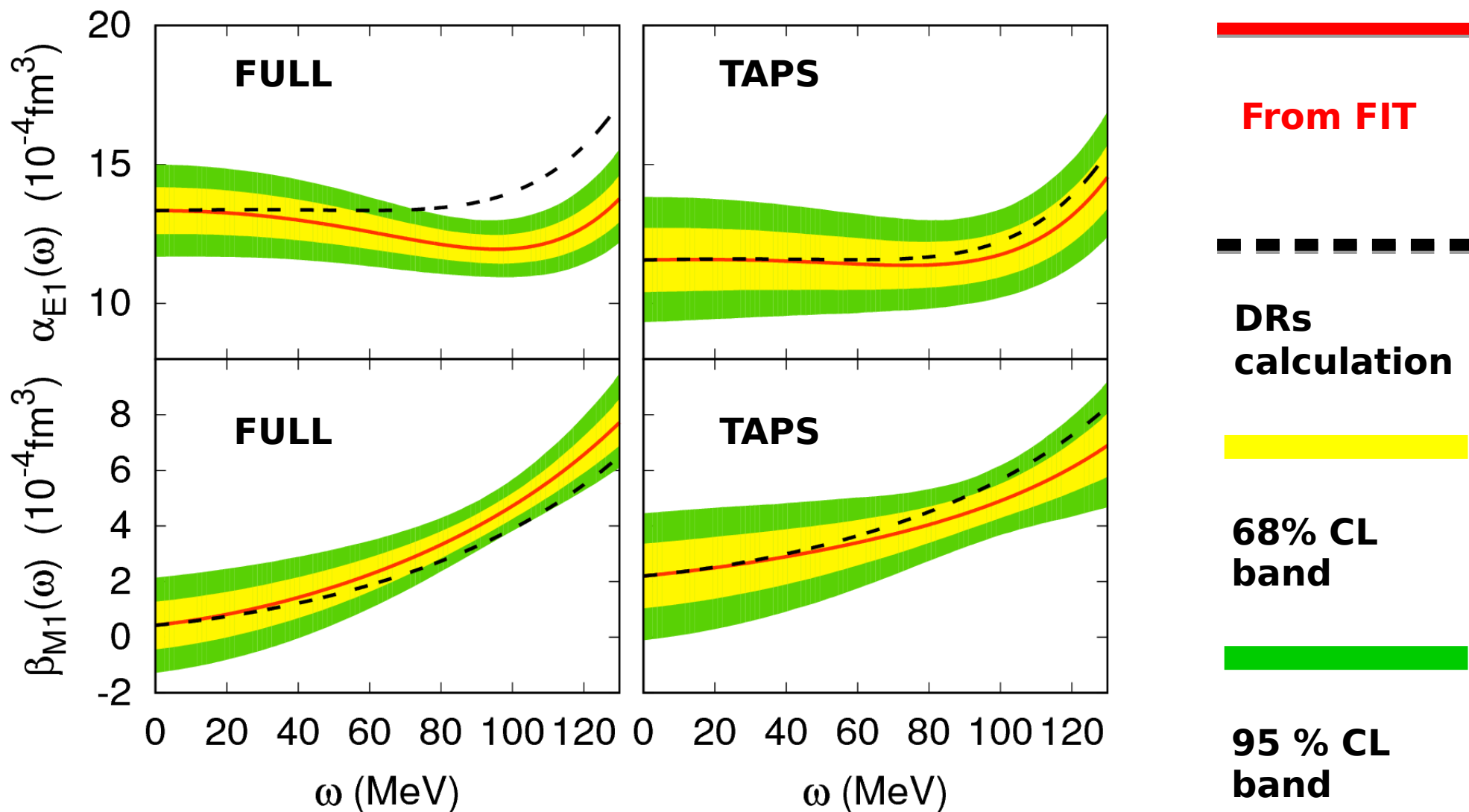


Fit conditions

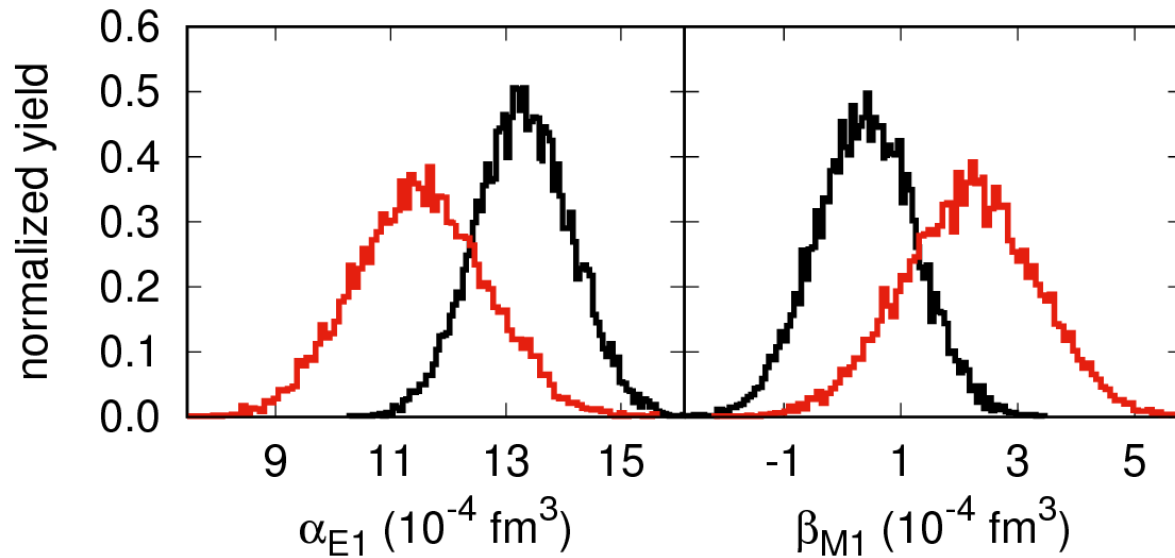
- ✓ Baldin's sum rule
- ✓ Systematical errors ON
- ✓ FULL data set (150 data)
- ✓ TAPS data set (55 data)
- ✓ Errors on Baldin's sum rule and γ_n included in the procedure



DDPs from the fit (II)

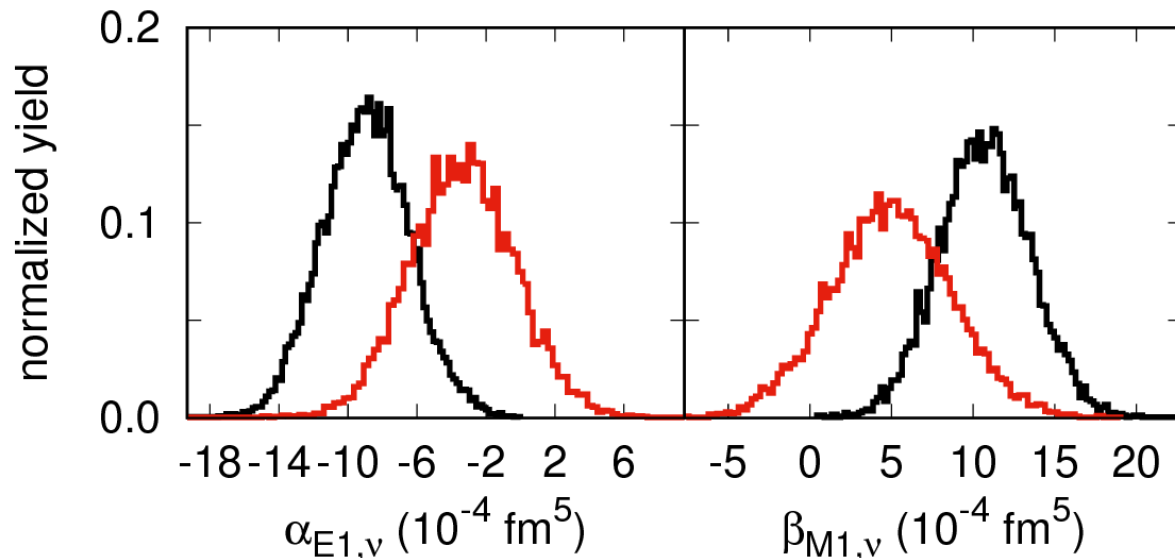


DDPs from the fit: probability distributions



TAPS data set

FULL data set



Probability distributions given by our technique (**not** a priori assumed)

Global results: numerical values

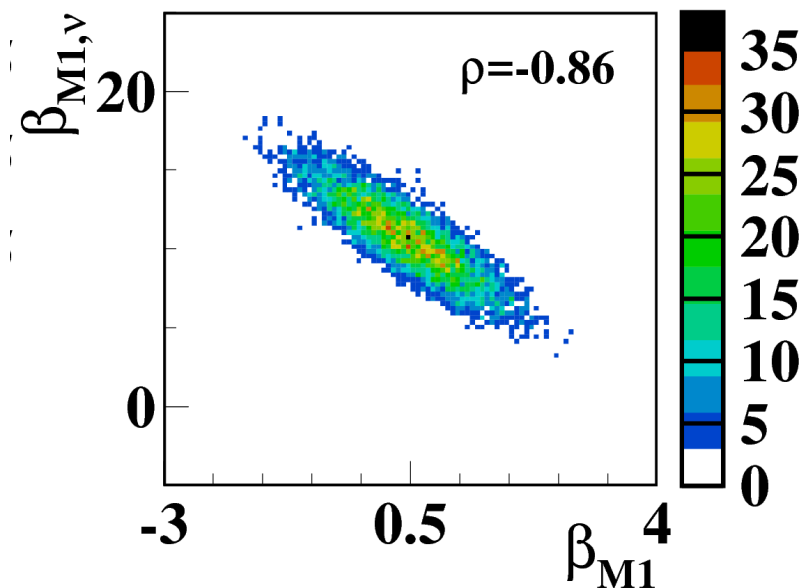
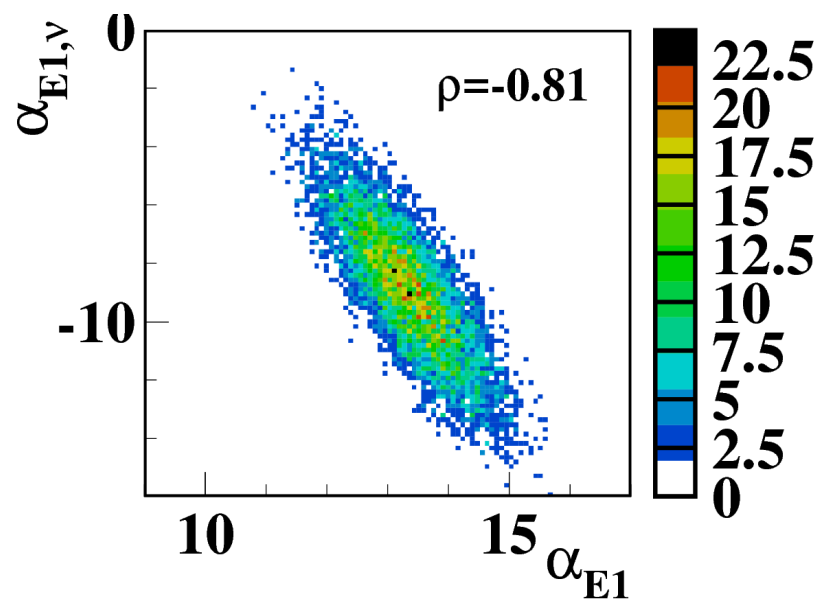
		FULL	TAPS
α_{E1}	(10^{-4}fm^3)	13.3 ± 0.8	11.6 ± 1.1
$\alpha_{E1,\nu}$	(10^{-4}fm^5)	-8.8 ± 2.5	-3.2 ± 3.1
β_{M1}	(10^{-4}fm^3)	0.4 ∓ 0.9	2.2 ∓ 1.1
$\beta_{M1,\nu}$	(10^{-4}fm^5)	10.8 ± 2.8	5.1 ± 3.7

Global results: numerical values

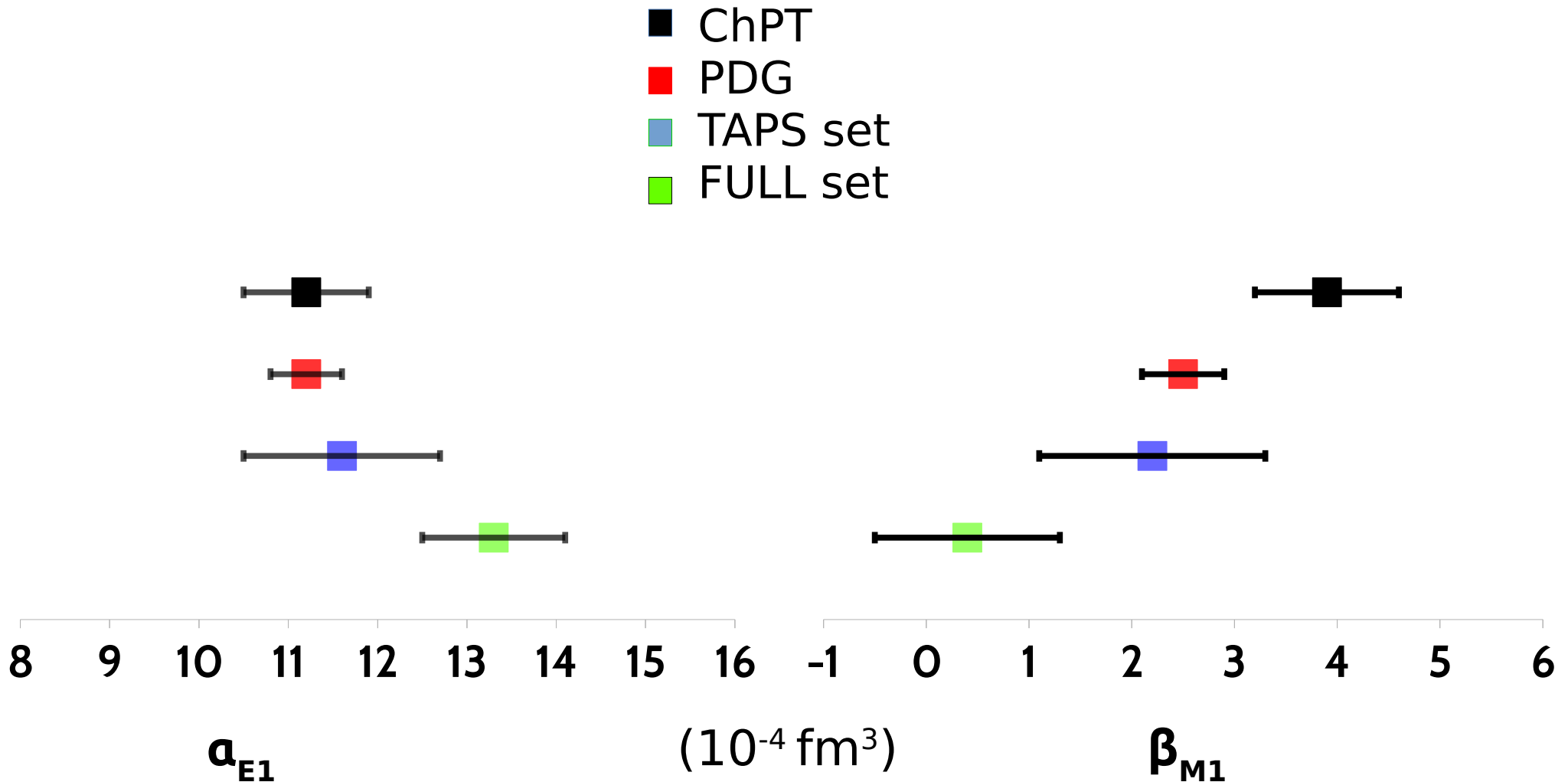
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Very STRONG dependence on data set (maybe due to different angular regions...)

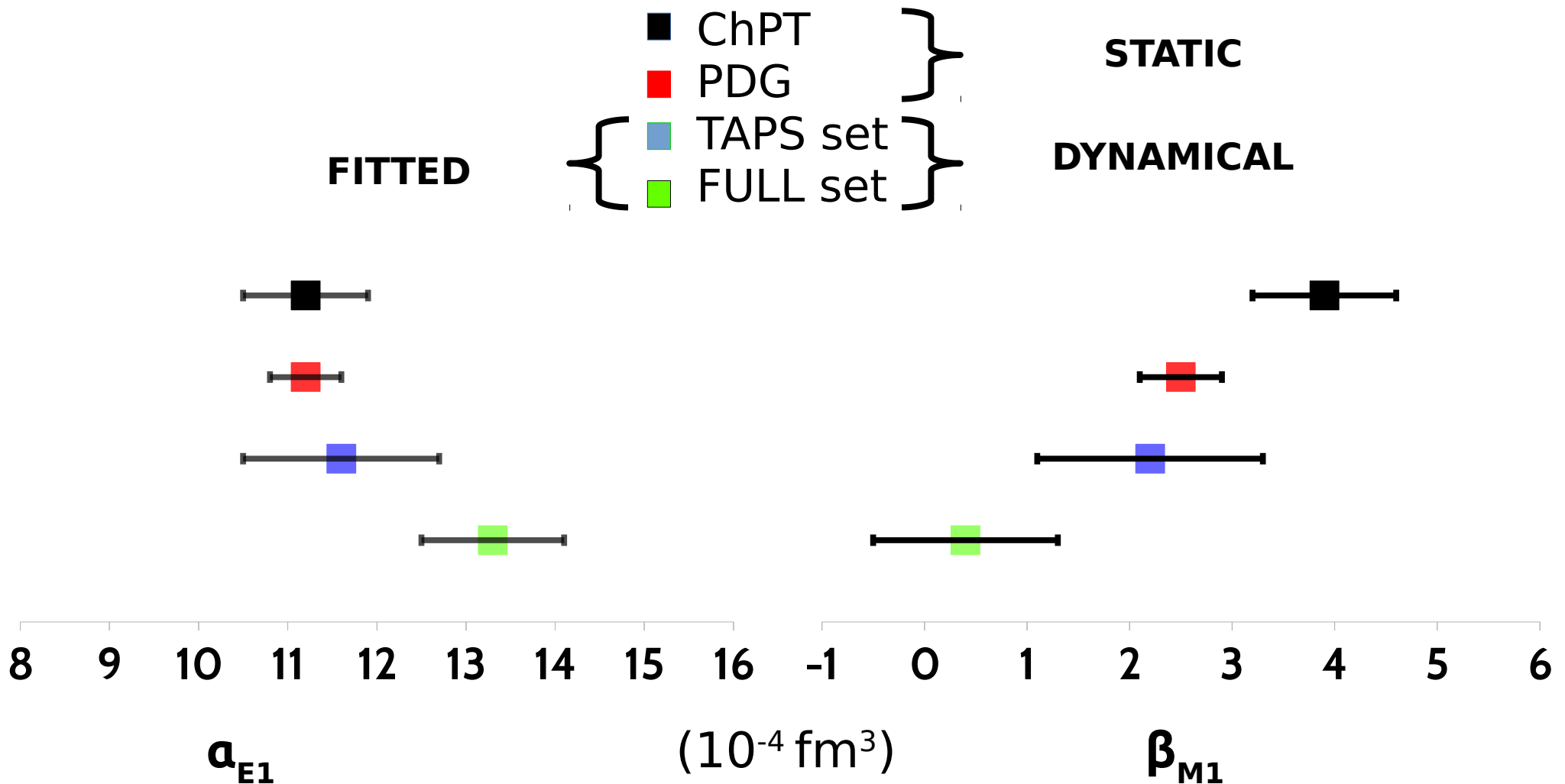
Very HIGH correlations among parameters



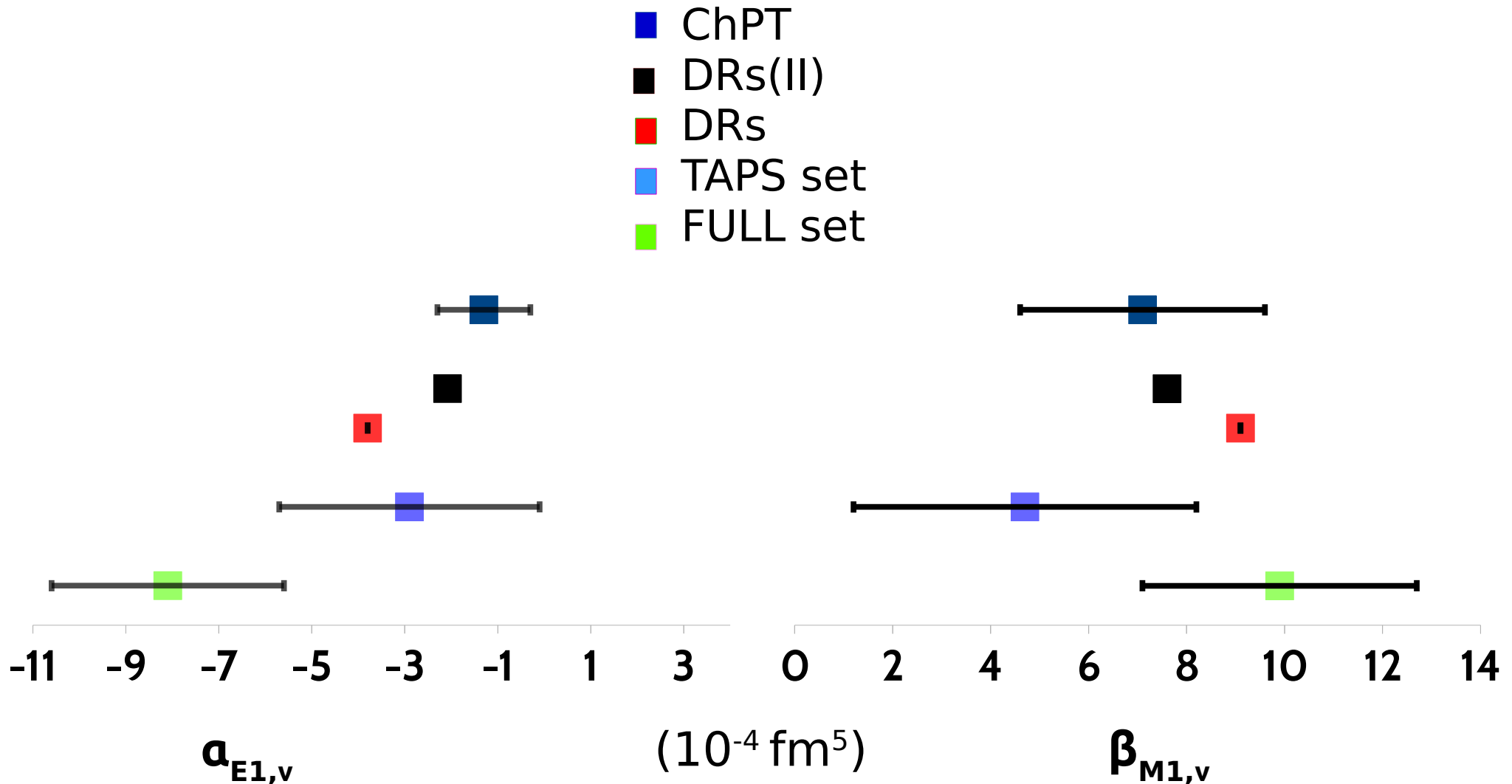
Global results: α_{E1} and β_{M1}



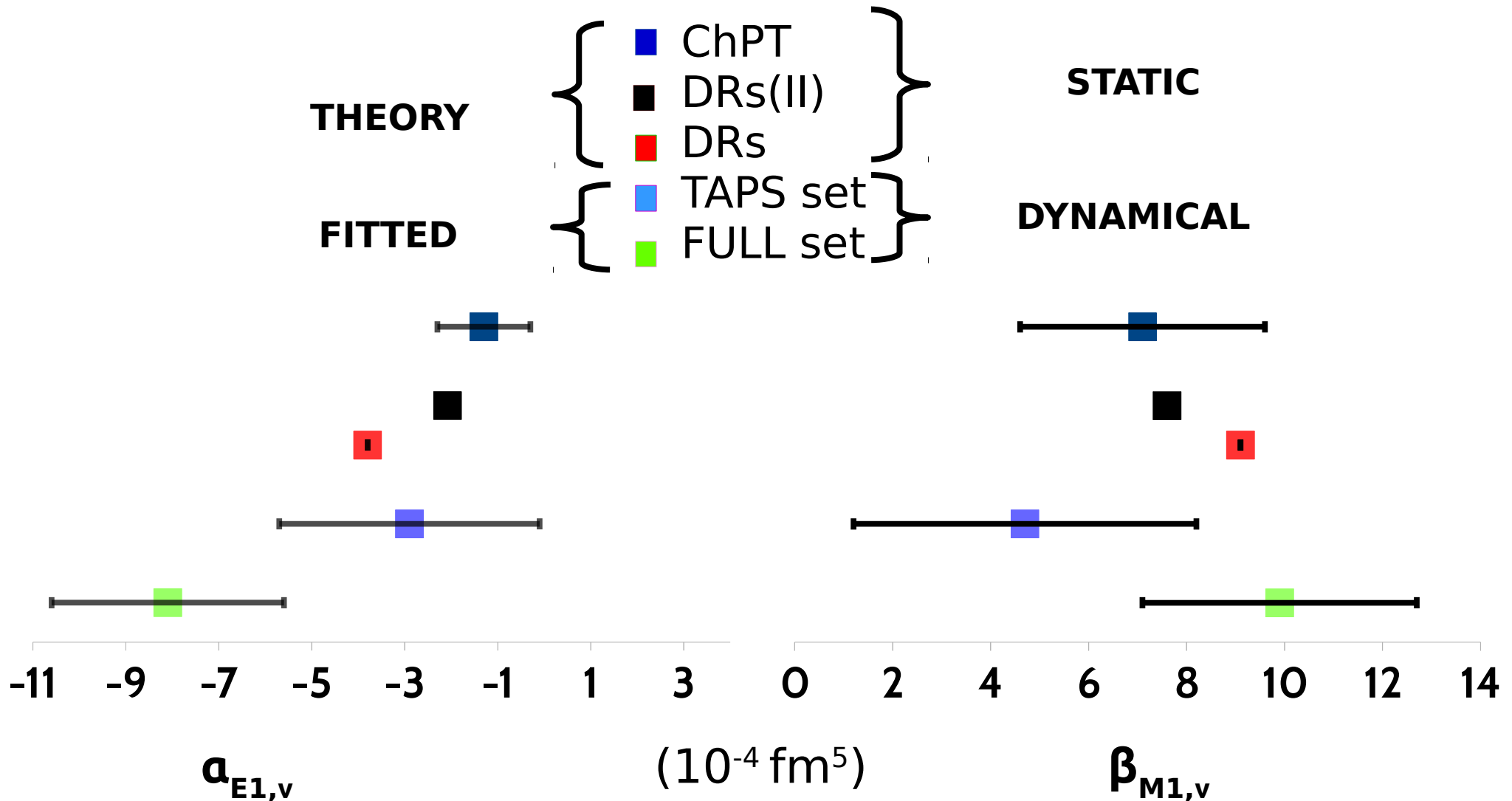
Global results: α_{E1} and β_{M1}



Global results: $\alpha_{E1,v}$ and $\beta_{M1,v}$



Global results: $\alpha_{E1,v}$ and $\beta_{M1,v}$



Conclusions & perspectives

Very useful and versatile technique for data analysis

Effect of systematic sources of uncertainties on the fitted parameters

Waiting for new data in order to reduce the uncertainties of the fitted parameters (MAMI)

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DDPs without LEX (double subtraction in DRs)

Fit of polarized observables in RCS with the same technique

BACKUP SLIDES



Some comments on the data set

Differential cross section

TAPS vs FULL data set

LEX is very slow...

χ^2 curvature close to its
minimum

Outliers identifications: rescaling of
the errors

Some comments on the data set

Strong correlation between parameters

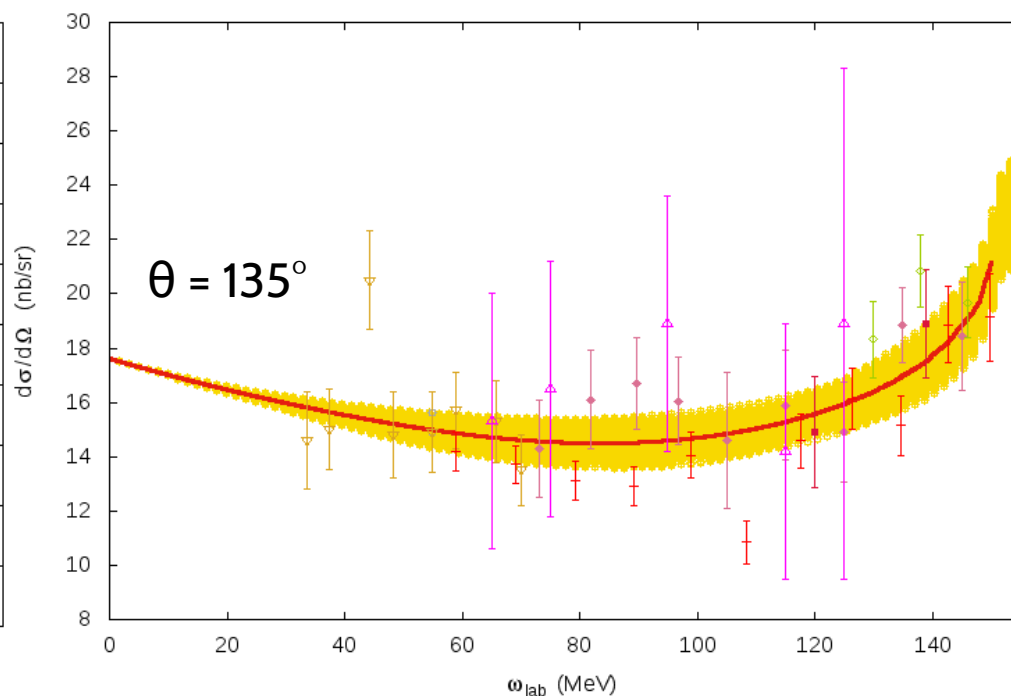
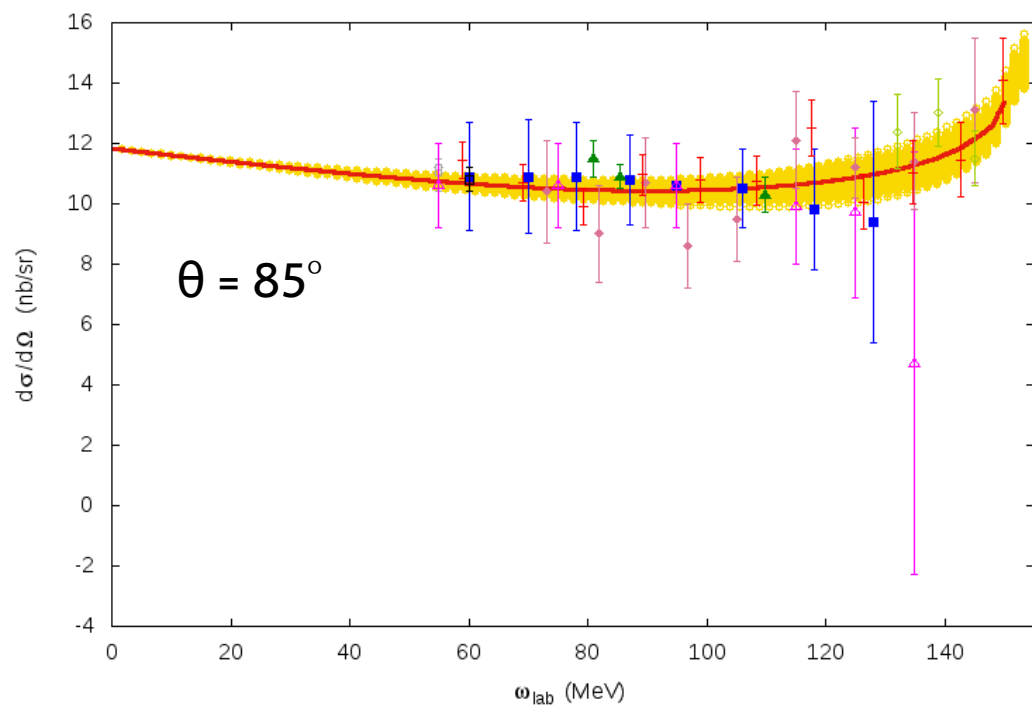
Reduction of parameters number thanks to sum rules

Identifications of the *outliers* (rescaling for statistic errors?)

The χ^2 is not the only *quality indicator* → no “definition” of data set

Waiting for new data (MAMI)

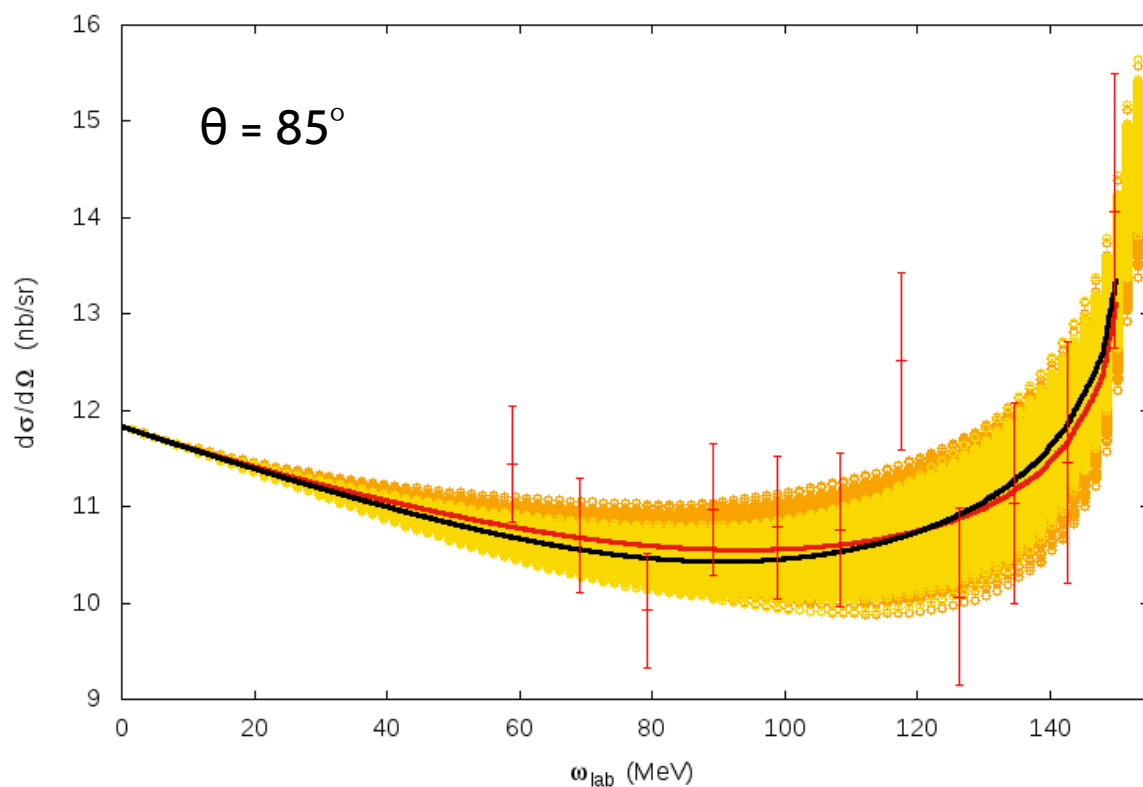
Differential cross section



$d\sigma/d\Omega$ VS lab energy

100% error band from the bootstrap fit

TAPS vs FULL data set



TAPS

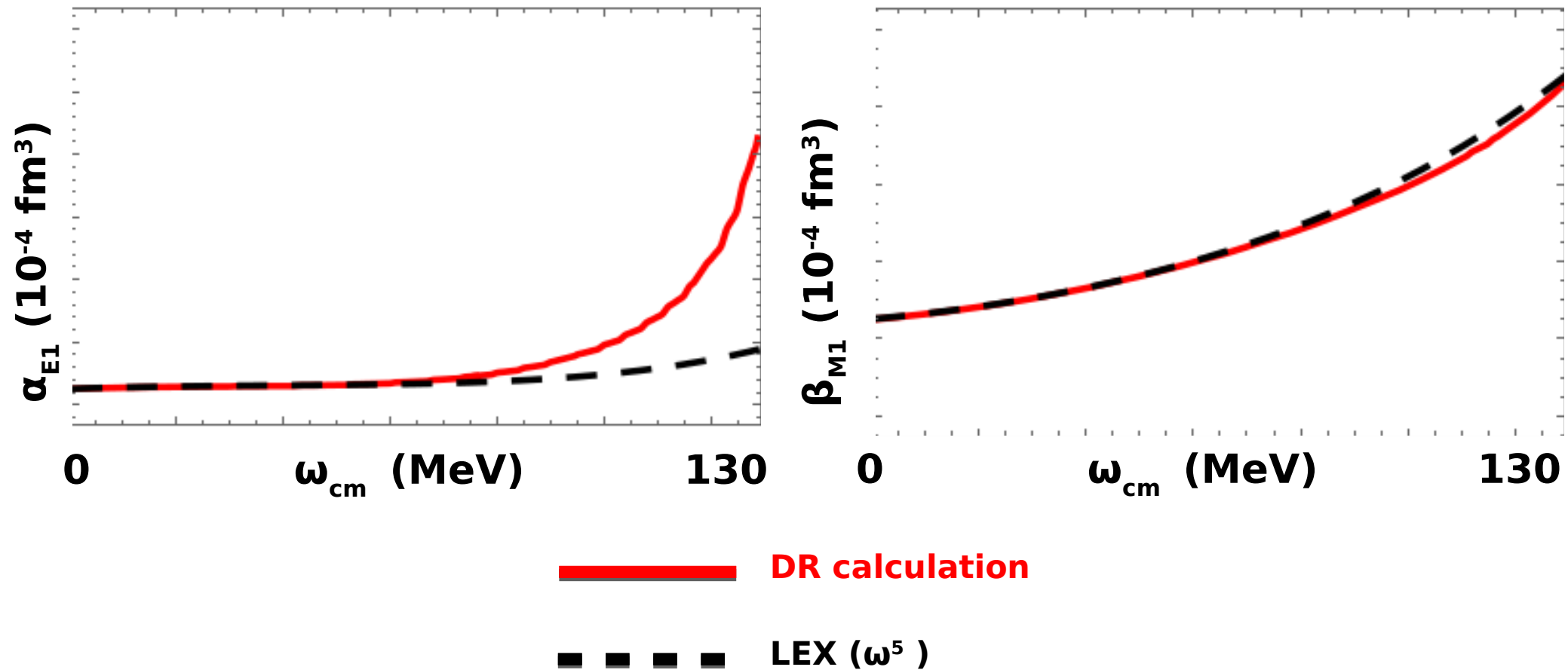
FULL

TAPS 100% error band

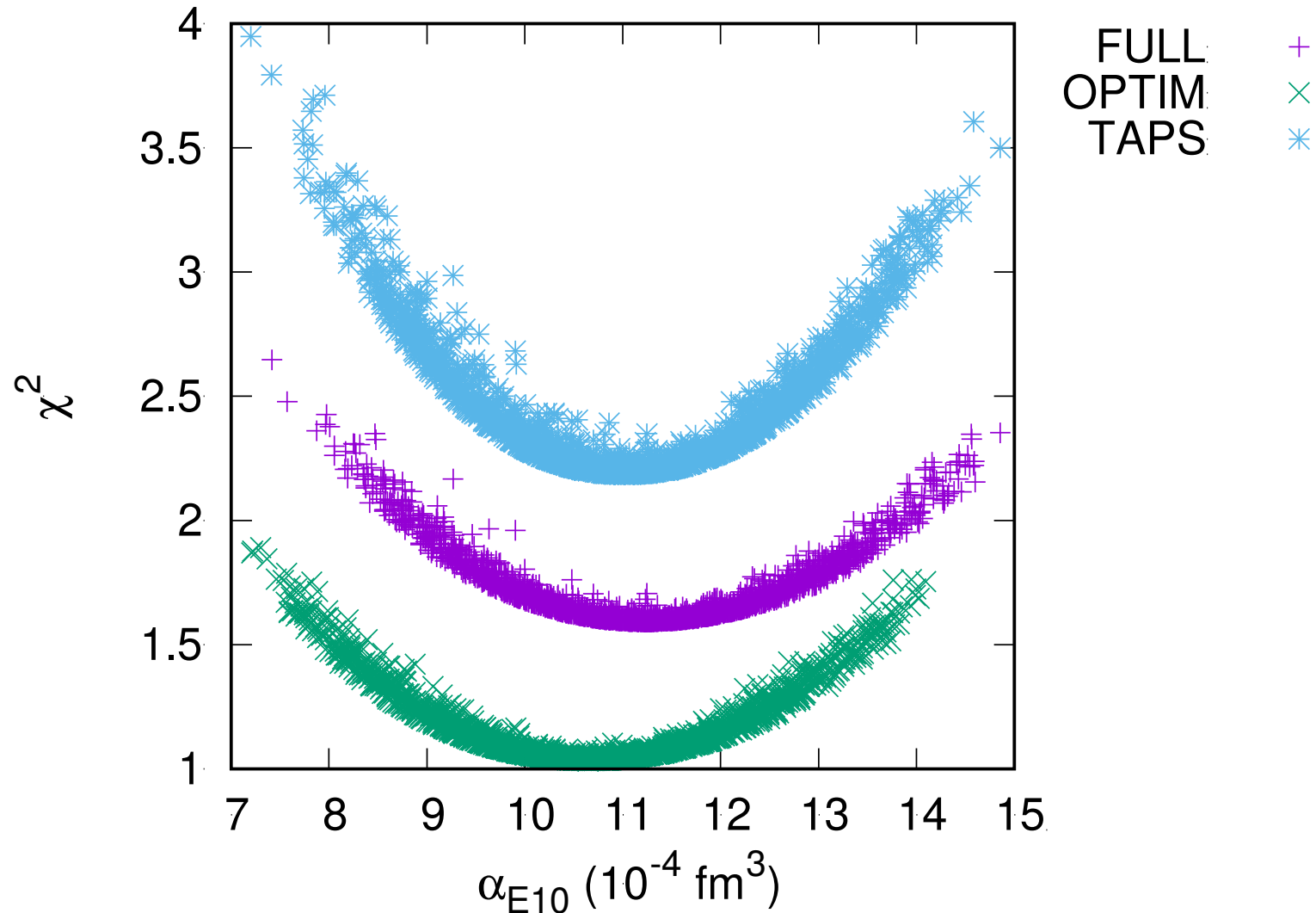
FULL 100% error band

VERY small difference both in calculation and in error band

LEX is very slow ...



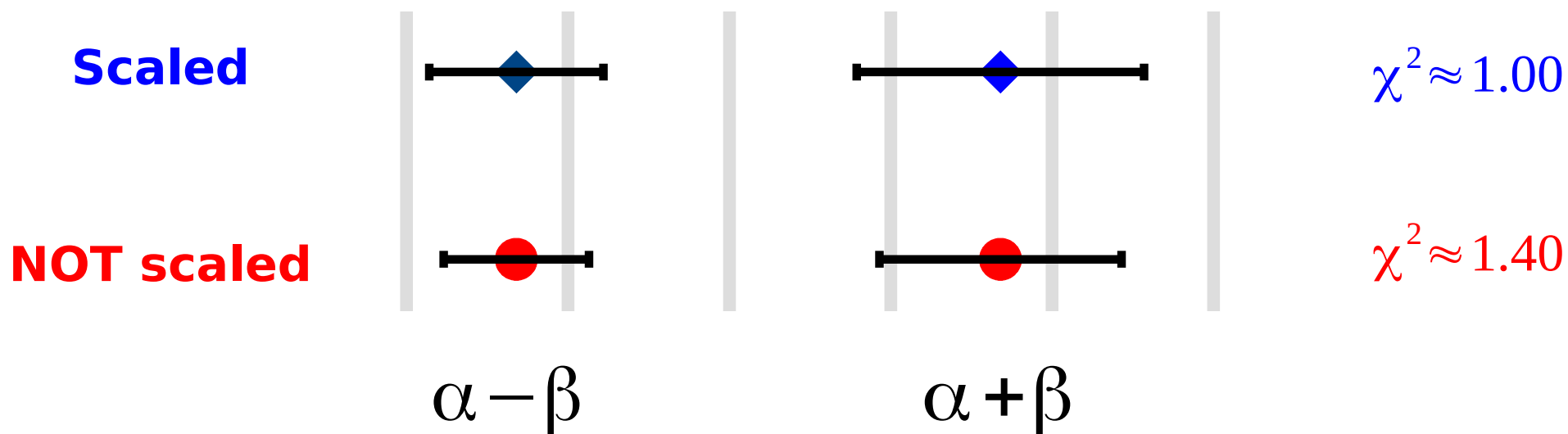
χ^2 curvature close to its minimum



Outliers identification: rescaling of the errors

Outliers → rescaling of all the statistic uncertainties by a factor

$$\sqrt{\chi^2}$$



Effect: enlarging of errors on fitted parameters ($\sim 20\%$)