

Understanding $d^*(2380)$ in a chiral quark model

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Collaborators:

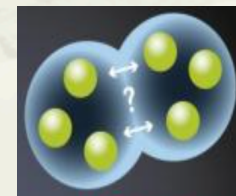
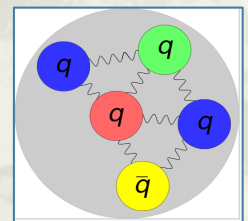
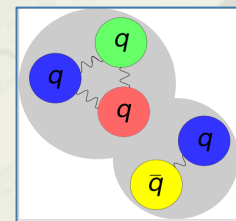
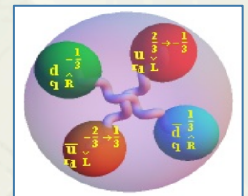
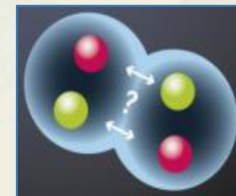
Y.B. Dong, P.N. Shen, Z.Y. Zhang (IHEP, CAS)

Outline

- Motivation
- Experimental observation
- Status of theoretical investigations
- Results from chiral SU(3) quark model calculations
 - Mass & structure
 - Decay widths
 - Charge distribution
- Summary

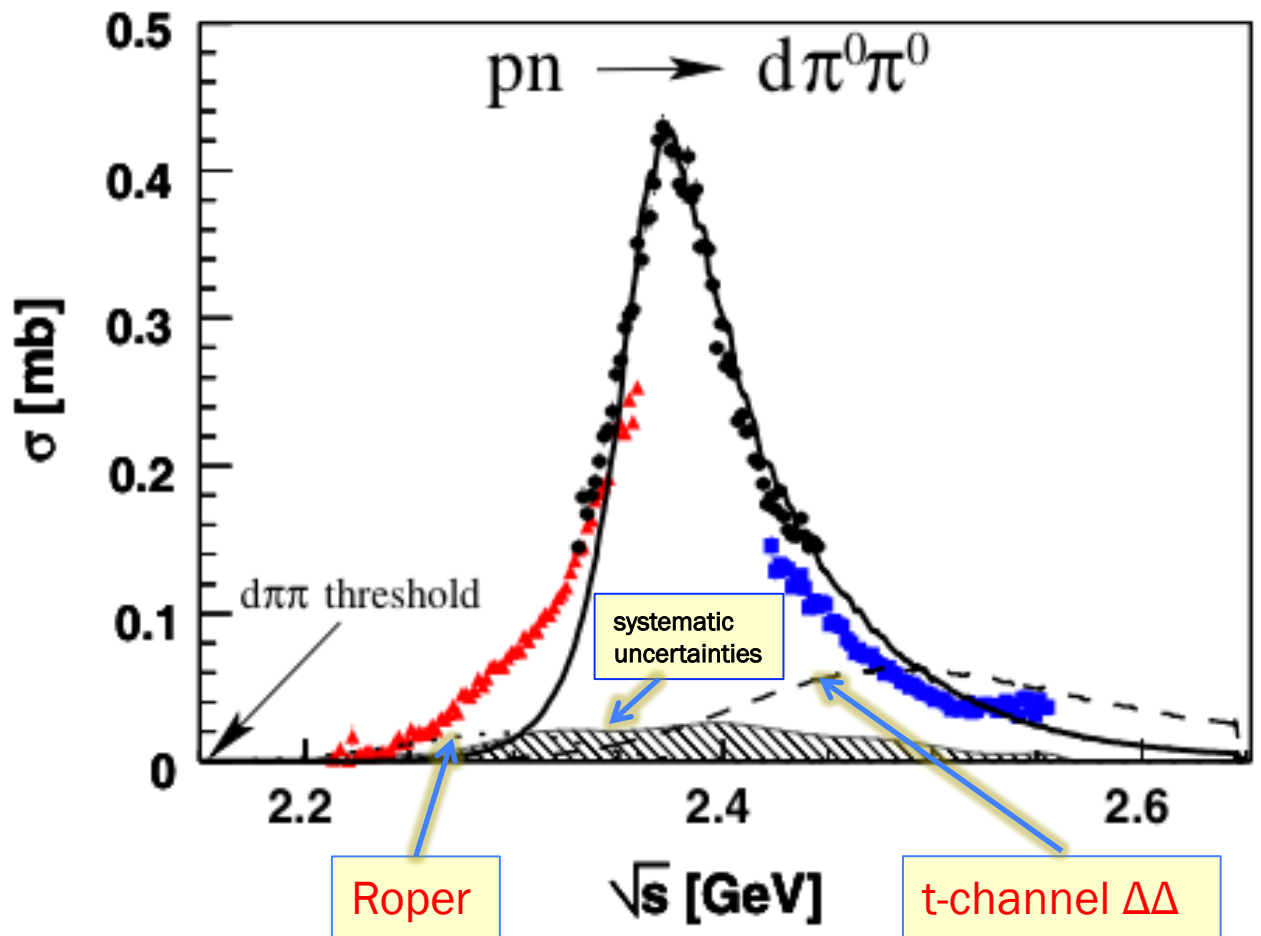
Motivation

- Quarks are confined into hadrons instead of being “bare” & “alone”
- Most hadron states are categorized into mesons ($q\bar{q}$) and baryons (qqq)
- Exotics not excluded by QCD
- 4- or 5-quark configurations claimed in heavy flavor physics
- Possible 6-quark systems: deuteron is the only confirmed one
- Any other BB molecules? Hexaquark states?

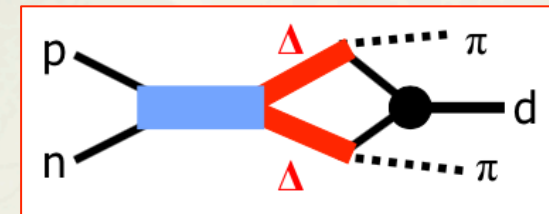


Experiments @ COSY

WASA-at-COSY, PRL106(2011)242302



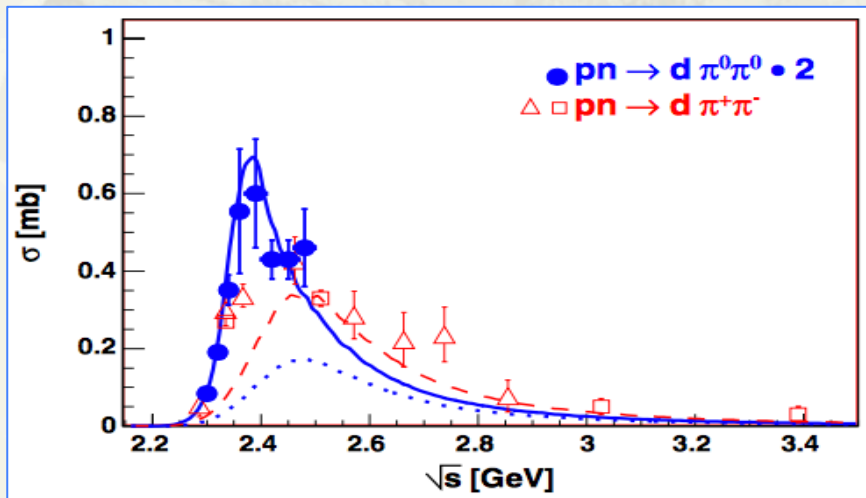
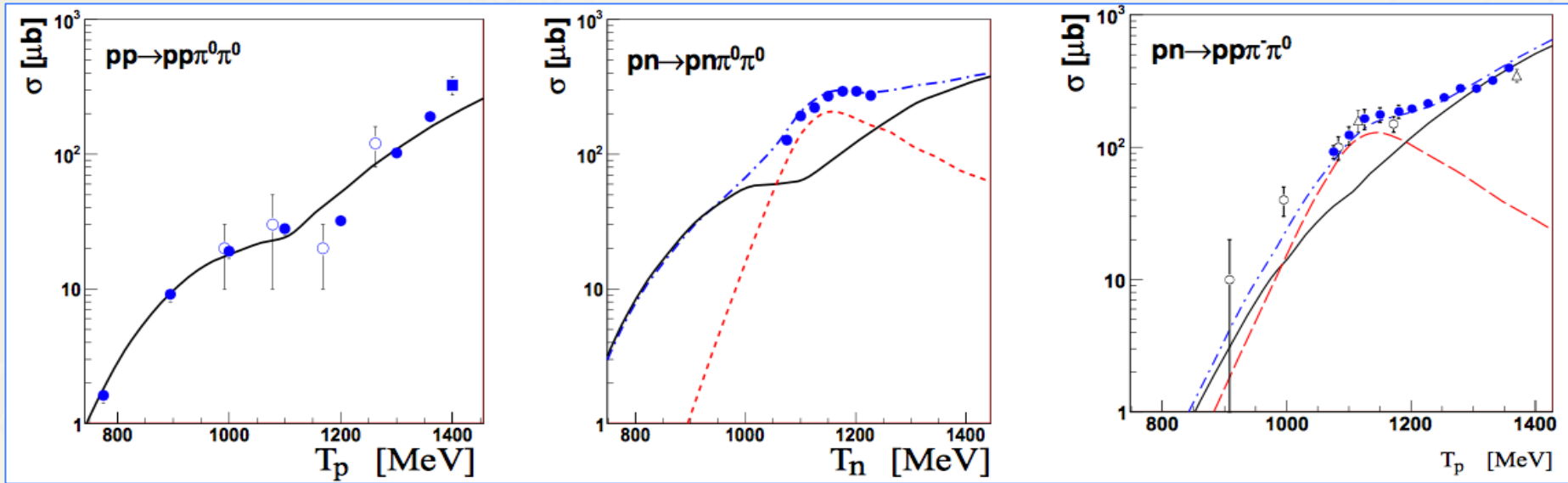
- ✧ Exclusive
- ✧ Kinematically complete



$I(J^p) = 0(3^+)$
 $M \approx 2380 \text{ MeV}$
 $\Gamma \approx 70 \text{ MeV}$

$d^*(2380)$

Signals in other reactions @ COSY



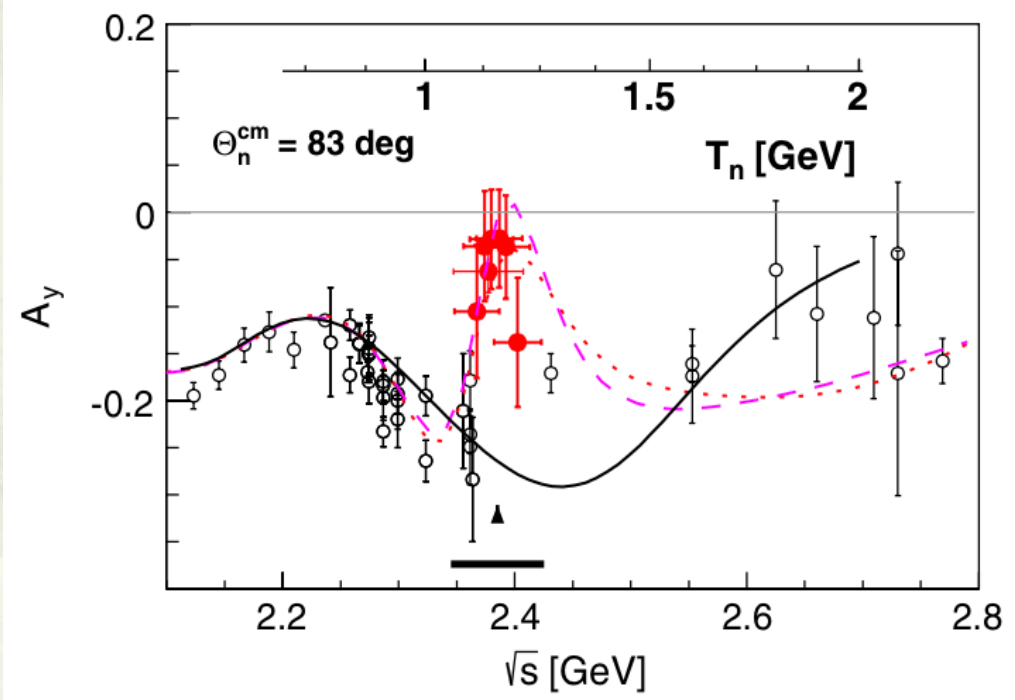
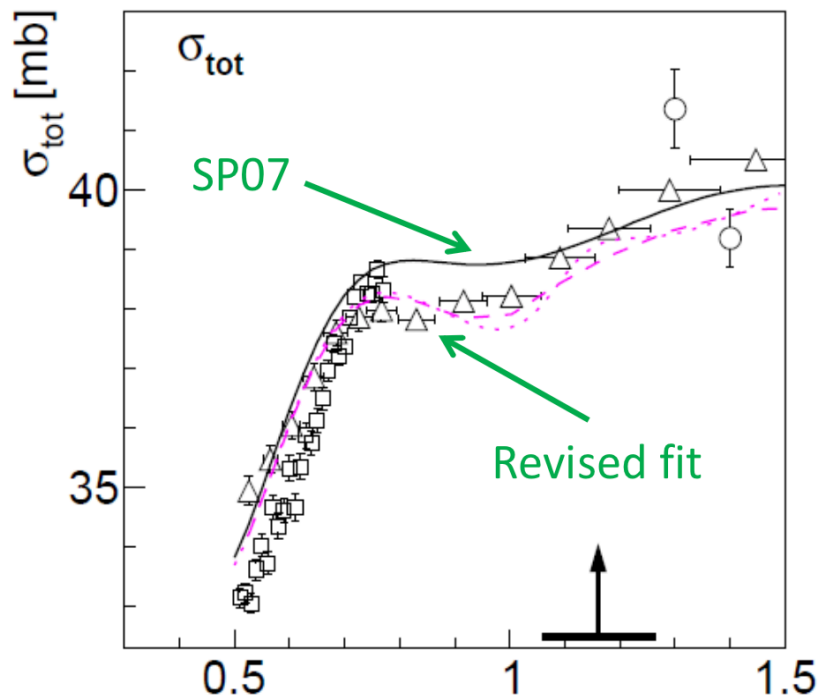
Measured also in fusion reactions to helium isotopes:



Evidence from $\vec{n}p$ scattering

$\vec{d}p \rightarrow np + p_{\text{spectator}}$

$$M = (2380 \pm 10) - i(40 \pm 5)$$



WASA-at-COSY & SAID DAC, PRL112(2014)202301

Unusual narrow width of d^*

2464 $\Delta\Delta$

2380 d^*

2309 $\Delta N\pi$

2154 $NN\pi\pi\pi$

1878 NN

$$M_{d^*} \approx 2380 \text{ MeV}$$

$$\approx 2M_{\Delta} - 84 \text{ MeV}$$

$$> M_{\Delta N\pi}$$

$$> M_{NN\pi\pi\pi}$$

$$> M_{NN}$$

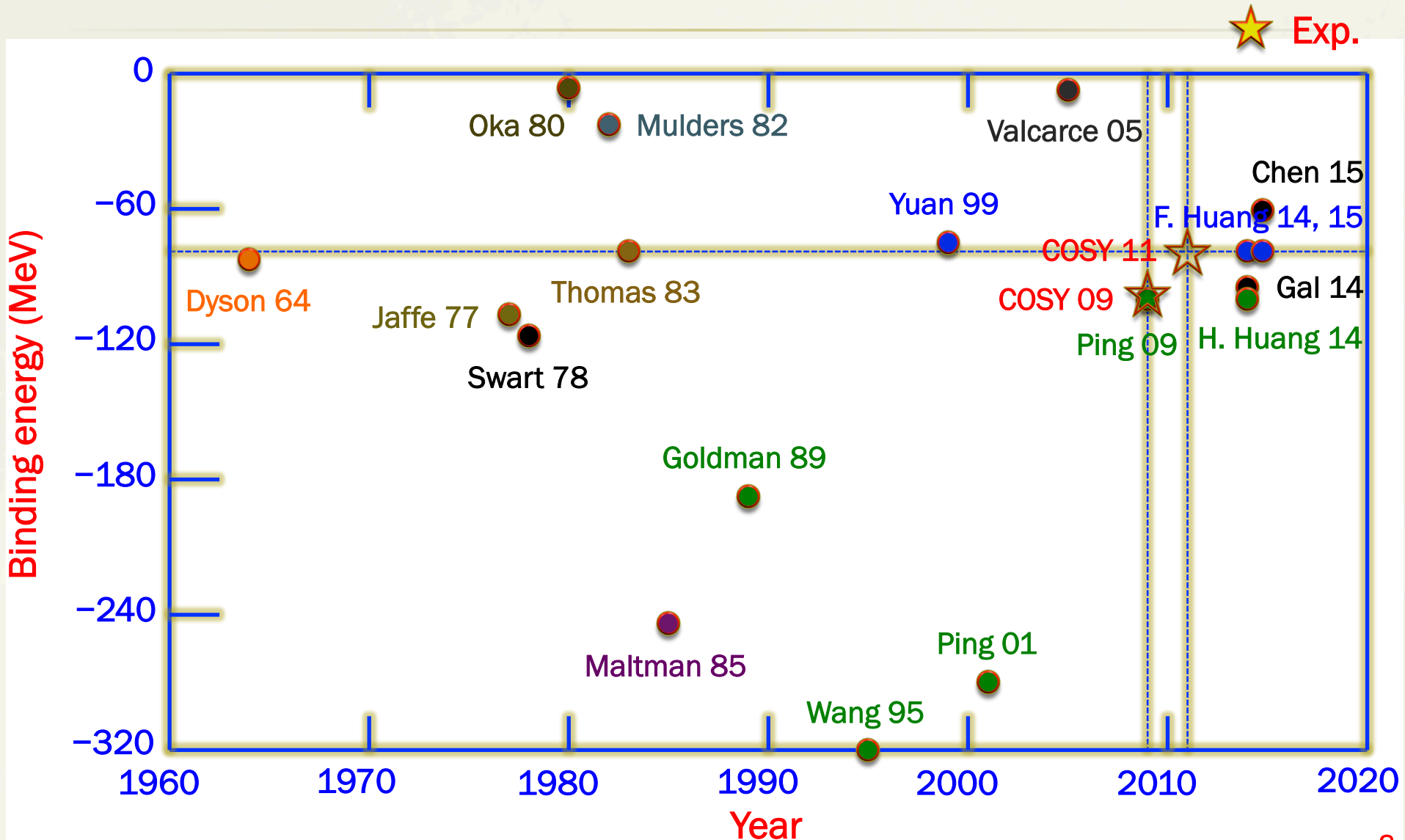
$$2\Gamma_{\Delta} \approx 230 \text{ MeV}$$

$$\Gamma_{d^*} \approx 70 \text{ MeV}$$

$$< 1/3 \times 2\Gamma_{\Delta}$$



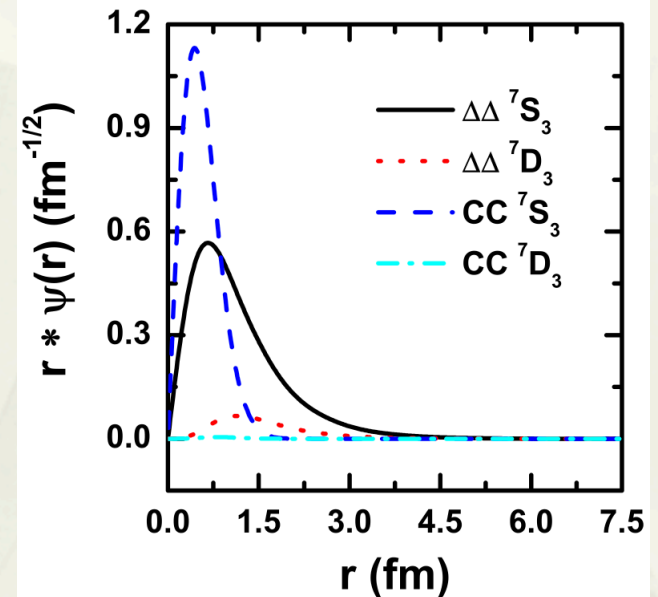
Theoretical $\Delta\Delta$ binding energies



Chiral SU(3) QM study, revisited

Structures & wave functions

- F. Huang, Z.Y. Zhang, P.N. Shen, W.L. Wang, Chin. Phys. C 39 (2015) 071001
- F. Huang, P.N. Shen, Y.B. Dong, Z.Y. Zhang, Sci. China-Phys. Mech. Astron. 59 (2016) 622002



Decay widths & charge distributions

- Y.B. Dong, P.N. Shen, F. Huang, Z.Y. Zhang, Phys. Rev. C 91 (2015) 064002
- Y.B. Dong, F. Huang, P.N. Shen, Z.Y. Zhang, Phys. Rev. C 94 (2015) 014003
- Y.B. Dong, F. Huang, P.N. Shen, Z.Y. Zhang, Phys. Lett. B 769 (2017) 223
- Y.B. Dong, F. Huang, P.N. Shen, Z.Y. Zhang, Phys. Rev. D 96 (2017) 094001

Other explanations in literature

◆ A. Gal & H. Garcilazo, NPA928(2014)73

- Dynamically generated $\Delta N\pi$ 3-body resonance
- Binding energy: 101 MeV
- Width: 66 MeV

$$B_{\text{exp}} \approx 84 \text{ MeV}$$
$$\Gamma_{\text{exp}} \approx 70 \text{ MeV}$$

◆ H.X. Huang, J.L. Ping, & F. Wang, PRC89(2014)034001

- $\Delta\Delta$ bound state
- Binding energy: 71 MeV (ChQM), 107 MeV (QDCSM)
- Width: 150 MeV (ChQM), 110 MeV (QDCSM)

◆ H.X. Chen, E.L. Cui, & W. Chen et al., PRC91(2015)025204

- QCD sum rule analysis
- Mass: 2.4 ± 0.2 GeV

The Chiral SU(3) quark model

SU(2) linear σ model

$$\Sigma = \sigma + i \sum_{a=1}^3 \tau_a \pi_a$$

$$\begin{aligned} \mathcal{L}_I^{\text{ch}} &= -g \left(\bar{\psi}_L \Sigma \psi_R + \bar{\psi}_R \Sigma^\dagger \psi_L \right) \\ &= -g \bar{\psi} \left(\sigma + i \gamma_5 \sum_{a=1}^3 \tau_a \pi_a \right) \psi \end{aligned}$$

Chiral SU(3) quark model

$$\Sigma = \sum_{a=0}^8 \lambda_a \sigma_a + i \sum_{a=0}^8 \lambda_a \pi_a$$

$$\begin{aligned} \mathcal{L}_I^{\text{ch}} &= -g \left(\bar{\psi}_L \Sigma \psi_R + \bar{\psi}_R \Sigma^\dagger \psi_L \right) \\ &= -g \bar{\psi} \left(\sum_{a=0}^8 \lambda_a \sigma_a + i \gamma_5 \sum_{a=0}^8 \lambda_a \pi_a \right) \psi \end{aligned}$$

- Chiral symmetry restored by introducing S & PS fields
- CQ obtains constituent mass via spontaneous CSB
- GB gets mass via explicit CSB caused by tiny current quark mass

Hamiltonian of Chiral QM

Total Hamiltonian for 6q systems:

$$H = \sum_{i=1}^6 \left(m_i + \frac{\vec{P}_i^2}{2m_i} \right) - T_{\text{cm}} + \sum_{1=i<j}^6 \left(V_{ij}^{\text{conf}} + V_{ij}^{\text{OGE}} + V_{ij}^{\text{ch}} \right)$$

Ch. SU(3) QM:

$$V_{ij}^{\text{ch}} = \sum_{a=0}^8 V_{ij}^{\sigma_a} + \sum_{a=0}^8 V_{ij}^{\pi_a}$$

Ext. Ch. SU(3) QM:

$$V_{ij}^{\text{ch}} = \sum_{a=0}^8 V_{ij}^{\sigma_a} + \sum_{a=0}^8 V_{ij}^{\pi_a} + \sum_{a=0}^8 V_{ij}^{\rho_a}$$

Note: OGE almost completely reduced by including VMEs.

Determination of parameters

- **Input:** $m_u = m_d = 313 \text{ MeV}$,
 $b_u = 0.5 \text{ fm (SU(3))}$ & $0.45 \text{ fm (ex. SU(3))}$

- **Coupling between quark & chiral fields:**

$$\frac{g_{\text{ch}}^2}{4\pi} = \left(\frac{3}{5}\right)^2 \frac{g_{NN\pi}^2}{4\pi} \frac{m_u^2}{m_N^2}, \quad \frac{g_{NN\pi}^2}{4\pi} = 13.67$$

- **Mass of mesons:** experimental values except for m_σ

- **Coupling constant for OGE:** $g_u \propto m_\Delta - m_N$

- **Confinement strength & zero point energy:**

$$\frac{\partial m_N}{\partial b_u} = 0, \quad m_N = 939 \text{ MeV}$$

RGM study of $\Delta\Delta$ -CC

RGM wave functions for $\Delta\Delta$ -CC system:

$$\begin{aligned}\Psi_{6q} = & \mathcal{A} \left[\hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \eta_{\Delta\Delta}(\vec{r}) \right]_{S=3, I=0, C=(00)} \\ & + \mathcal{A} \left[\hat{\phi}_C^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_C^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \eta_{CC}(\vec{r}) \right]_{S=3, I=0, C=(00)}\end{aligned}$$

$$\Delta: \quad (0S)^3 [3]_{\text{orb}}, \quad S = \frac{3}{2}, \quad I = \frac{3}{2}, \quad C = (00)$$

$$C: \quad (0S)^3 [3]_{\text{orb}}, \quad S = \frac{3}{2}, \quad I = \frac{1}{2}, \quad C = (11)$$

RGM equation for a bound state problem:

$$\langle \delta\psi_{6q} | H - E | \psi_{6q} \rangle = 0$$

Calculated d^* mass

Without CC: BE \approx 29 – 62 MeV

	$\Delta\Delta (L = 0, 2)$		
	SU(3)	Ext. SU(3) (f/g=0)	Ext. SU(3) (f/g=2/3)
B (MeV)	28.96	62.28	47.90
RMS (fm)	0.96	0.80	0.84

With CC: BE \approx 47 – 84 MeV

	$\Delta\Delta - \text{CC} (L = 0, 2)$		
	SU(3)	Ext. SU(3) (f/g=0)	Ext. SU(3) (f/g=2/3)
B (MeV)	47.27	83.95	70.25
RMS (fm)	0.88	0.76	0.78
$(\Delta\Delta)_{L=0}$ (%)	33.11	31.22	32.51
$(\Delta\Delta)_{L=2}$ (%)	0.62	0.45	0.51
$(\text{CC})_{L=0}$ (%)	66.25	68.33	66.98
$(\text{CC})_{L=2}$ (%)	0.02	0.00	0.00

- d^* : a deeply bound & compact $\Delta\Delta$ -CC state

- Coupling to CC plays a significant role

- Predicted binding energy close to experimental value

$$M_{d^*} \approx 2M_{\Delta} - 84 \text{ MeV}$$

Distinctive features of $\Delta\Delta$: why

Quark-exchange effect:

$$\psi_{6q} = \mathcal{A} \left[\hat{\phi}^{\text{int}}(\bar{\xi}_1, \bar{\xi}_2) \hat{\phi}^{\text{int}}(\bar{\xi}_4, \bar{\xi}_5) \eta(\vec{r}) \right]$$

- For 6 identical quarks: $\mathcal{A} = 1 - 9P_{36}$
- Quark-exchange effect: $\langle \mathcal{A}^{sfc} \rangle \in [0, 2]$
- $(\Delta\Delta)_{S=3, I=0}$: $\langle \Delta\Delta | \mathcal{A}^{sfc} | \Delta\Delta \rangle_{S=3, I=0} = 2$

Strongly “attractive”!

Short-range interaction:

OGE: **attractive** + VMEs: **attractive**

Oka & Yazaki, PLB90(1980)41:

For non-strange BB systems, $(\Delta\Delta)_{S=3, I=0}$ is the only one in which OGE provides attraction at short-range.

Deuteron: $\langle NN | \mathcal{A}^{sfc} | NN \rangle_{S=1, I=0} = 10/9 \sim 1$

OGE: **repulsive** + VMEs: **repulsive**

Channel wave functions

Channel wave functions (Relative wave functions in physical basis):

$$\begin{aligned}
 \chi_{\Delta\Delta}(\vec{r}) &= \left\langle \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \middle| \psi_{6q} \right\rangle \\
 &= \left\langle \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \middle| (1 - 9P_{36}) \left[\hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \eta_{\Delta\Delta}(\vec{r}) \right] \right\rangle \\
 &\quad - 9 \left\langle \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \middle| P_{36} \left[\hat{\phi}_C^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_C^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \eta_{CC}(\vec{r}) \right] \right\rangle
 \end{aligned}$$

$$\begin{aligned}
 \chi_{CC}(\vec{r}) &= \left\langle \hat{\phi}_C^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_C^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \middle| \psi_{6q} \right\rangle \\
 &= \left\langle \hat{\phi}_C^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_C^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \middle| (1 - 9P_{36}) \left[\hat{\phi}_C^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_C^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \eta_{CC}(\vec{r}) \right] \right\rangle \\
 &\quad - 9 \left\langle \hat{\phi}_C^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_C^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \middle| P_{36} \left[\hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \eta_{\Delta\Delta}(\vec{r}) \right] \right\rangle
 \end{aligned}$$

Wave function of d^*

Reorganize the wave function of d^* :

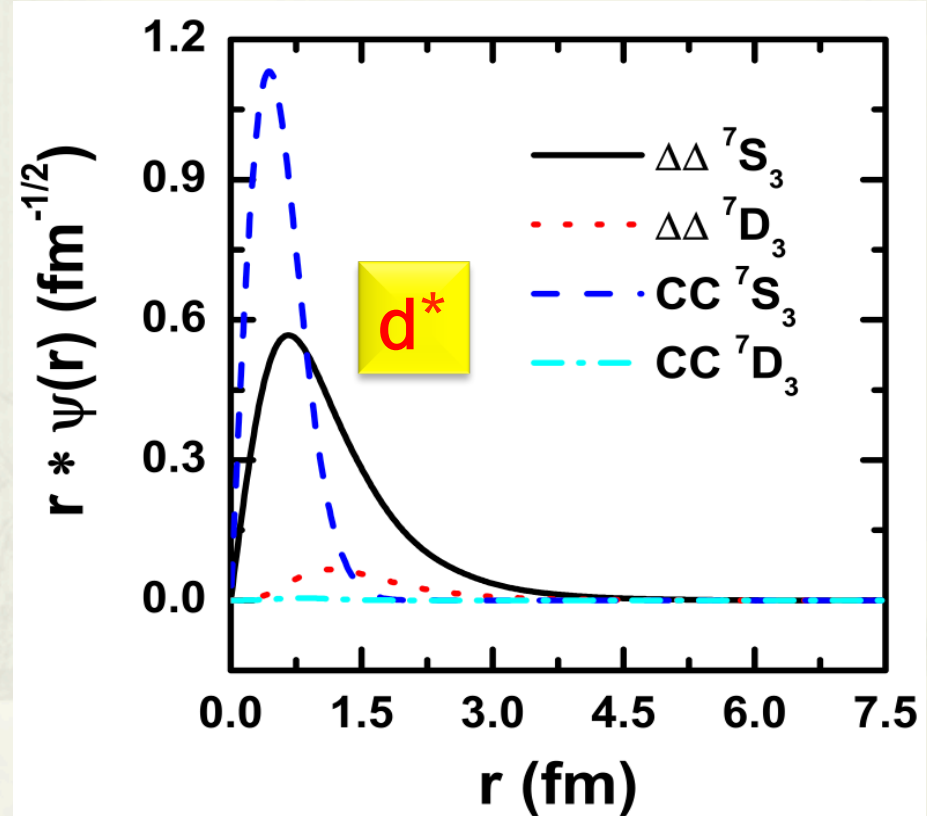
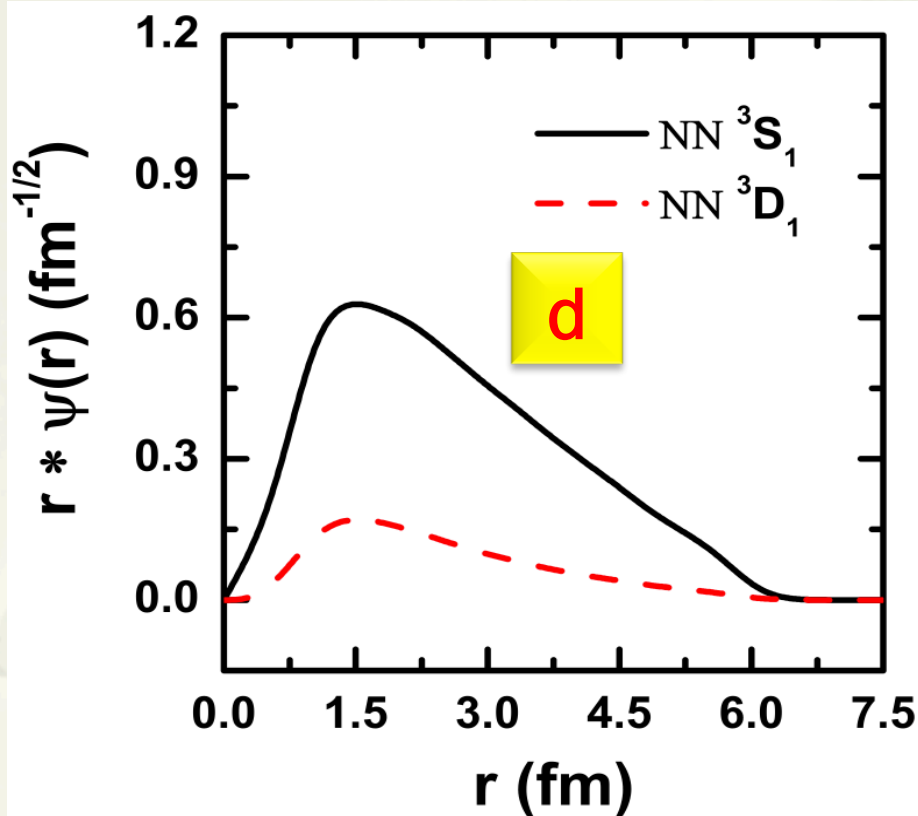
$$\begin{aligned}\psi_{d^*} &= |\Delta\Delta\rangle \chi_{\Delta\Delta}(\vec{r}) + |CC\rangle \chi_{CC}(\vec{r}) \\ &= \sum_{L=0,2} \left[|\Delta\Delta\rangle \frac{\chi_{\Delta\Delta}^L(r)}{r} + |CC\rangle \frac{\chi_{CC}^L(r)}{r} \right] Y_{L0}(\hat{r})\end{aligned}$$

$\Delta\Delta$ & CC parts are now orthogonal to each other:

$$\langle CC | \Delta\Delta \rangle = 0$$

$\chi_{\Delta\Delta}$, χ_{CC} are used to discuss the spatial distribution of d^* and its individual components of $\Delta\Delta$ & CC

Relative wave function



Unlike deuteron, d^* is rather narrowly distributed!

CC component

- d^* has a CC fraction of about 2/3

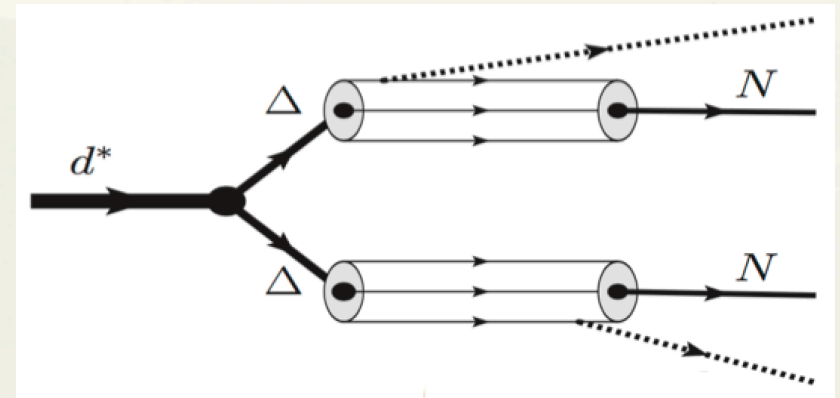
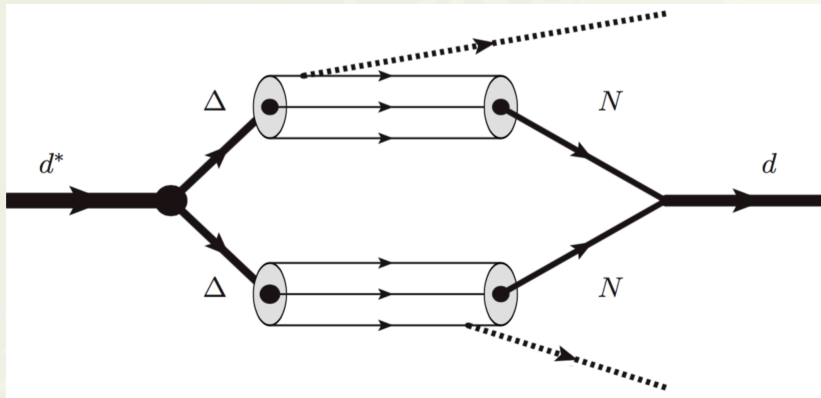
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B (MeV)	47.27	83.95	70.25
RMS (fm)	0.88	0.76	0.78
$(\Delta\Delta)_{L=0}$ (%)	33.11	31.22	32.51
$(\Delta\Delta)_{L=2}$ (%)	0.62	0.45	0.51
$(\text{CC})_{L=0}$ (%)	<u>66.25</u>	<u>68.33</u>	<u>66.98</u>
$(\text{CC})_{L=2}$ (%)	0.02	0.00	0.00

- A pure hexaquark state of $\Delta\Delta$ system has 4/5 CC fraction

$$[6]_{\text{orb}}[33]_{IS=03} = \sqrt{\frac{1}{5}} |\Delta\Delta\rangle_{IS=03} + \sqrt{\frac{4}{5}} |CC\rangle_{IS=03}$$

- d^* is a hexaquark-dominated exotic state!

2π-decay



Given the $qq\pi$ vertex & wave functions of N , Δ , d^* & d , the amplitudes & further decay widths can be calculated explicitly.

$$\mathcal{H} = g_{qq\pi} \vec{\sigma} \cdot \vec{k}_\pi \tau \cdot \phi \times \frac{1}{(2\pi)^{3/2} \sqrt{2\omega_\pi}}$$

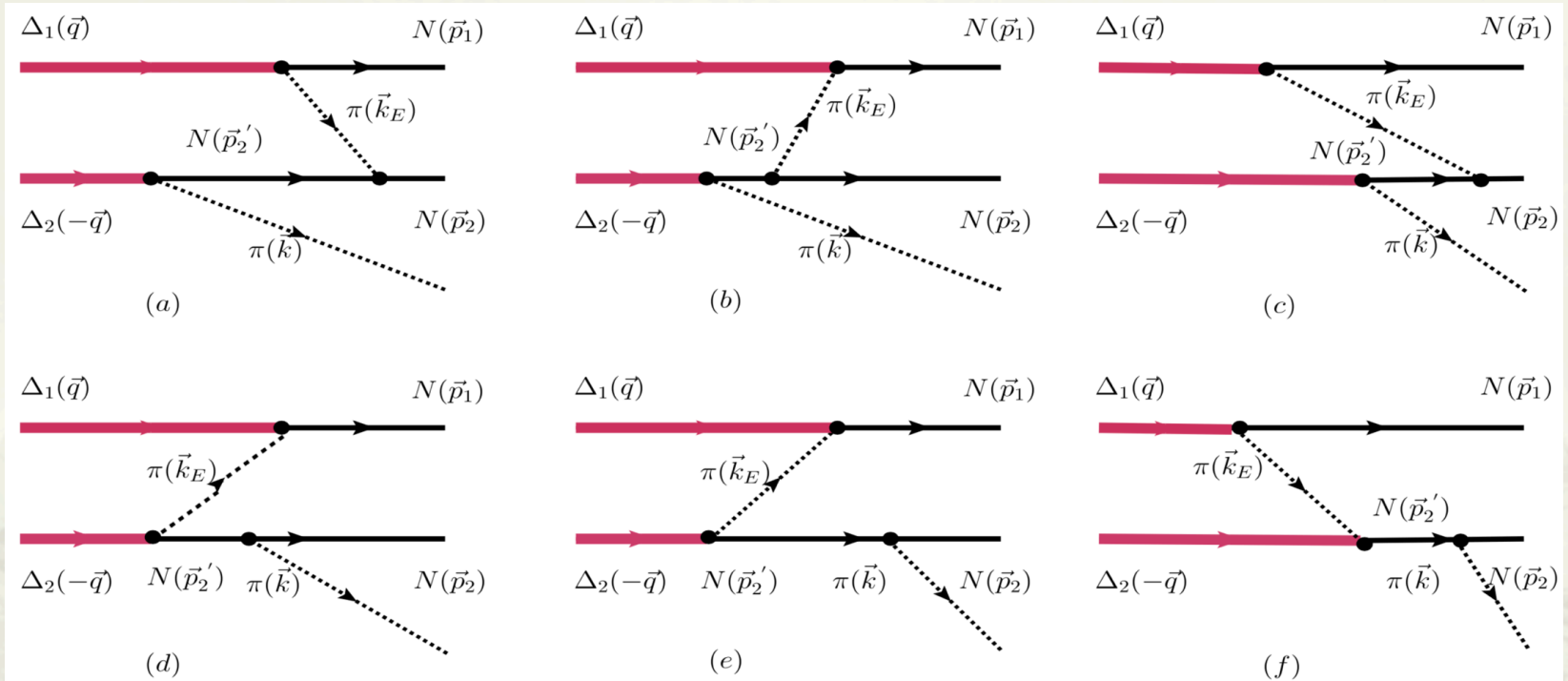
$$\Gamma_{d^* \rightarrow d\pi^0\pi^0} = \frac{1}{2!} \int d^3k_1 d^3k_2 d^3p_d (2\pi) \delta^3(\vec{k}_1 + \vec{k}_2 + \vec{p}_d) \delta(\omega_{k_1} + \omega_{k_2} + E_{p_d} - M_{d^*}) |\overline{\mathcal{M}}_{if}^{\pi^0\pi^0}|^2$$

$$\Gamma_{d^* \rightarrow pn\pi^0\pi^0} = \frac{1}{2!2!} \int d^3k_1 d^3k_2 d^3p_1 (2\pi) \delta(\Delta E) |\overline{\mathcal{M}}(k_1, k_2; p_1)|^2$$

Widths for 2π -decay

	Theor. (MeV)	Expt. (MeV)
$d^* \rightarrow d\pi^+\pi^-$	16.8	16.7
$d^* \rightarrow d\pi^0\pi^0$	9.2	10.2
$d^* \rightarrow pn\pi^+\pi^-$	20.6	21.8
$d^* \rightarrow pn\pi^0\pi^0$	9.6	8.7
$d^* \rightarrow pp\pi^0\pi^-$	3.5	4.4
$d^* \rightarrow nn\pi^0\pi^+$	3.5	4.4
$d^* \rightarrow pn$	8.7	8.7
Total	71.9	74.9

Single π decay



$$\begin{aligned} \mathcal{M}_{d^* \rightarrow NN\pi}^{(a)} &= \int d^3q \frac{\Psi_{d^*}(q)}{2\omega_{k_E} \sqrt{2\omega_k} (2\pi)^6} \delta^3(p_{N_2'} + p_{N_1'} + k - p_{\Delta_1} - p_{\Delta_2}) \\ &\times \tilde{\mathcal{M}}_{\pi(k_E)N(p_2') \rightarrow N(p_2)} \mathcal{D}_{af} \tilde{\mathcal{M}}_{\Delta_1 \rightarrow \pi(k_E)N(p_1)} \mathcal{D}_{ai} \tilde{\mathcal{M}}_{\Delta_2 \rightarrow \pi(k)N(p_2')} \end{aligned}$$

$$\mathcal{D}_{af} = \frac{1}{M_{d^*} - \omega(\vec{k}) - \omega(\vec{k}_E) - E_N(\vec{p}_1) - E_N(\vec{p}_2')}$$

$$\mathcal{D}_{ai} = \frac{1}{M_{d^*} - \omega(\vec{k}) - E_{\Delta_1}(\vec{q}) - E_N(\vec{p}_2')}$$

Results for single π decay

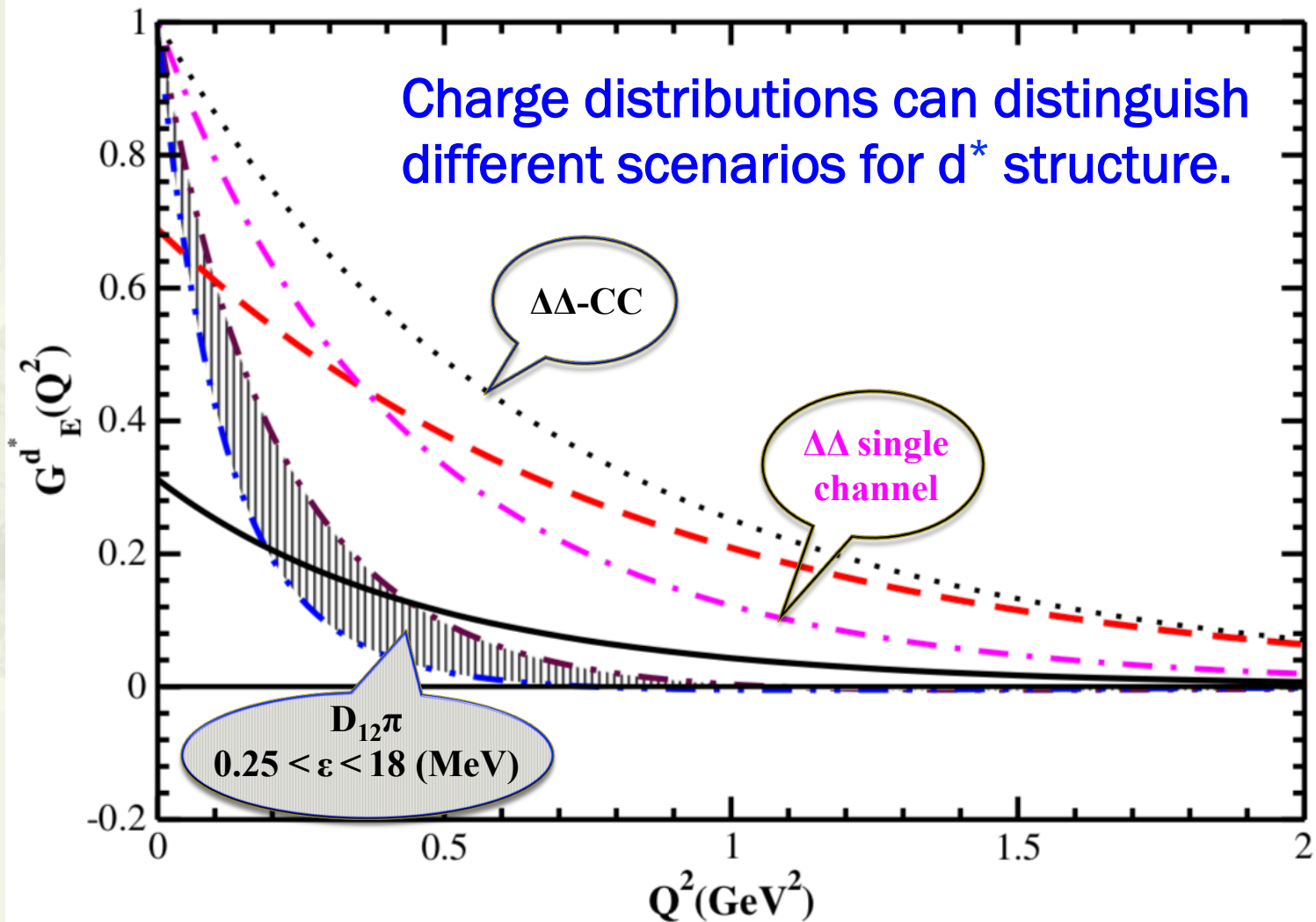
$$\Gamma_{d^* \rightarrow NN\pi} \approx 0.67 \text{ MeV} \quad \frac{\Gamma_{d^* \rightarrow NN\pi}}{\Gamma} \approx 0.9\%$$

The WASA-at-COSY Collaboration / Physics Letters B 774 (2017) 599–607

Exclusive measurements of the quasi-free $pn \rightarrow pp\pi^-$ and $pp \rightarrow pp\pi^0$ reactions have been performed by means of pd collisions at $T_p = 1.2$ GeV using the WASA detector setup at COSY. Total and differential cross sections have been obtained covering the energy region $T_p = 0.95\text{--}1.3$ GeV ($\sqrt{s} = 2.3\text{--}2.46$ GeV), which includes the regions of $\Delta(1232)$, $N^*(1440)$ and $d^*(2380)$ resonance excitations. From these measurements the isoscalar single-pion production has been extracted, for which data existed so far only below $T_p = 1$ GeV. We observe a substantial increase of this cross section around 1 GeV, which can be related to the Roper resonance $N^*(1440)$, the strength of which shows up isolated from the Δ resonance in the isoscalar $(N\pi)_{I=0}$ invariant-mass spectrum. No evidence for a decay of the dibaryon resonance $d^*(2380)$ into the isoscalar $(NN\pi)_{I=0}$ channel is found. An upper limit of 180 μb (90% C.L.) corresponding to a branching ratio of 9% has been deduced.

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d^* charge distribution results



Summary

- $d^*(2380)$ has been reported by WASA-at-COSY with an **unusual narrow width** ($\Gamma \approx 70$ MeV)
- $\Delta\Delta$ -CC with $I(J^P)=0(3^+)$ is dynamically investigated in our chiral SU(3) quark model and its extended version
- d^* has a CC fraction of about $2/3 \rightarrow$ it is a **hexaquark-dominated exotic state**
- Our calculated **binding energy & decay widths consistent with the data**
- **Charge distributions can be used to test different scenarios for the structure of $d^*(2380)$**
- More experiments & theoretical works needed



**THANK YOU FOR
YOUR PATIENCE!**

Our prediction in 1999

X. Q. Yuan, Z. Y. Zhang, Y. W. Yu, and P. N. Shen, Phys. Rev. C 60, 045203 (1999).

TABLE II. Binding energy B and rms \bar{R} of the deltaron $B = -(E_{\text{deltaron}} - 2M_{\Delta})$, $\bar{R} = \sqrt{\langle r^2 \rangle}$.

		$\Delta\Delta(L=0)$	$\Delta\Delta\begin{pmatrix} L=0 \\ +2 \end{pmatrix}$	$\frac{\Delta\Delta}{CC}(L=0)$	$\frac{\Delta\Delta}{CC}\begin{pmatrix} L=0 \\ +2 \end{pmatrix}$
OGE	B (MeV)	29.8	29.9	41.0	42.0
	\bar{R} (fm)	0.92	0.92	0.87	0.87
OGE + π, σ	B (MeV)	50.2	62.6	68.6	79.7
	\bar{R} (fm)	0.87	0.86	0.84	0.83
OGE + SU(3)	B (MeV)	18.4	22.5	31.7	37.3
	\bar{R} (fm)	1.01	1.00	0.92	0.92

- Binding energy: 40 ~ 80 MeV
- CC: 10 ~ 20 MeV increase in binding energy

Argument of Bashkanov et al.

M. Bashkanov, S. J. Brodsky, and H. Clement, Phys. Lett. B 727, 438 (2013).

- $Br(d^* \rightarrow \Delta\Delta)/Br(d^* \rightarrow pn) = 9:1$, but a deltaron with binding energy 90 MeV would have width ~ 160 MeV
- d^* must be of an unconventional origin, possibly indicating a genuine six-quark nature
- Two possible six-quark structures for $I(J^P)=0(3^+)$:

$$\left| \psi_{(0s)^6}^{[6]_{\text{orb}}[33]_{\text{IS}}} \right\rangle = \sqrt{\frac{1}{5}} |\Delta\Delta\rangle + \sqrt{\frac{4}{5}} |CC\rangle$$

$$\left| \psi_{(0s)^4(0p)^2}^{[42]_{\text{orb}}[33]_{\text{IS}}} \right\rangle = \sqrt{\frac{4}{5}} |\Delta\Delta\rangle - \sqrt{\frac{1}{5}} |CC\rangle$$

- It is nature to assign d^* to the former one — six-quark predominantly “hidden color state”

Resonating group method (RGM)

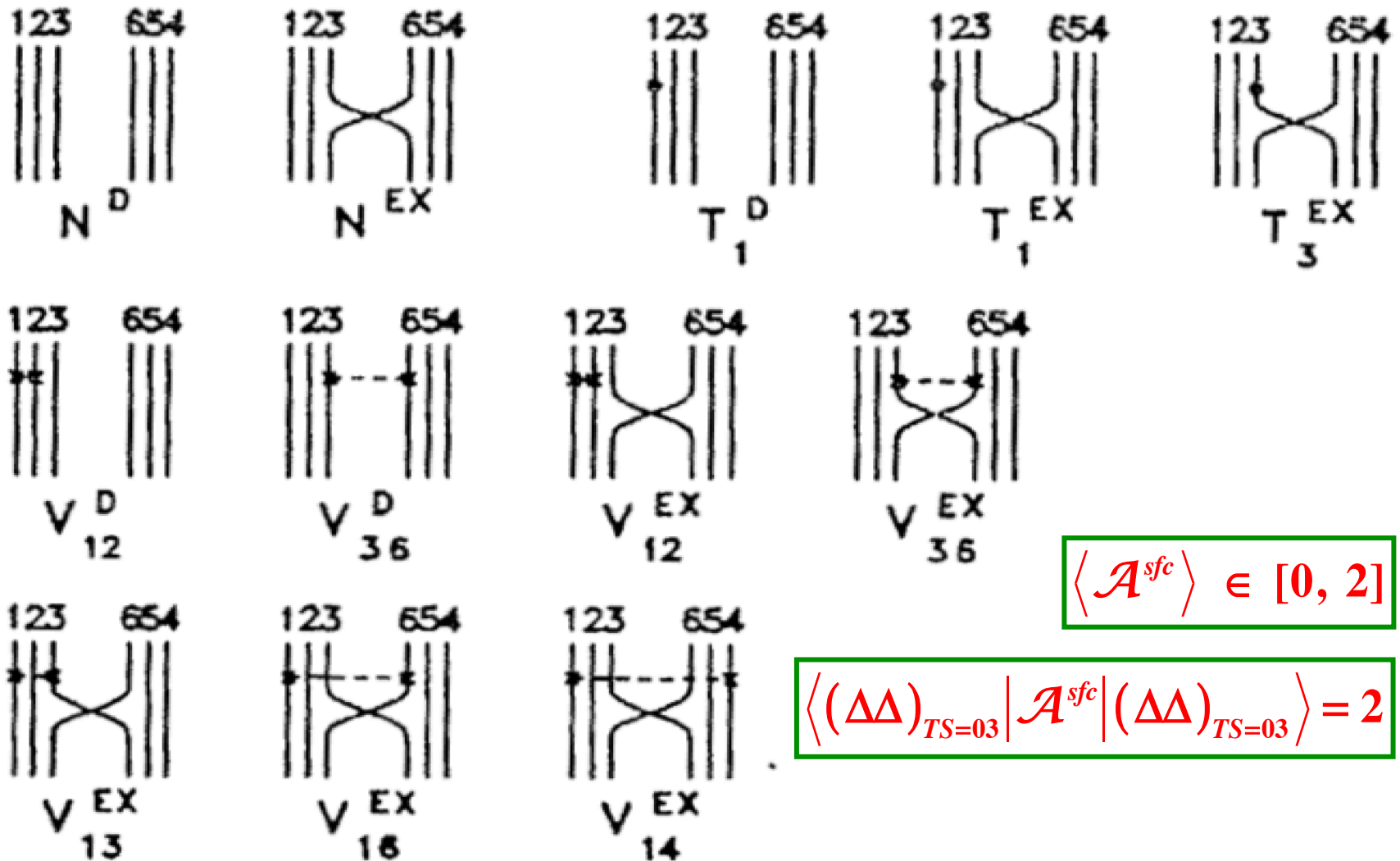
- RGM: well-established method for studying interactions between two composite particles; center of mass motion treated correctly
- Six-quark wave function in C.M. frame:

$$\Psi_{6q} = \mathcal{A} \left[\hat{\Phi}_A^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\Phi}_B^{\text{int}}(\vec{\xi}_3, \vec{\xi}_4) \chi(\vec{R}_{AB}) \right]_{STY}$$
$$\mathcal{A} \equiv \sum_{i \in A, j \in B} (1 - P_{ij}^{OSFC})$$

- Cluster wave functions in coordinate space: Gaussian
- Relative wave function determined by dynamics of the 6-quark system:

$$\langle \delta \Psi_{6q} | H - E | \Psi_{6q} \rangle = 0$$

Six-quark diagrams in RGM



Parameter values

All parameters fixed already in the study of NN scattering.
 No additional parameters introduced for $\Delta\Delta$ system.

TABLE I. Model parameters. The meson masses and the cutoff masses: $m_{\sigma'} = 980$ MeV, $m_{\epsilon} = 980$ MeV, $m_{\pi} = 138$ MeV, $m_{\eta} = 549$ MeV, $m_{\eta'} = 957$ MeV, $m_{\rho} = 770$ MeV, $m_{\omega} = 782$ MeV, and $\Lambda = 1100$ MeV.

	Ch. SU(3)	Ext. Ch. SU(3)	
		f/g=0	f/g=2/3
b_u (fm)	0.5	0.45	0.45
m_u (MeV)	313	313	313
g_u^2	0.766	0.056	0.132
g_{ch}	2.621	2.621	2.621
g_{chv}		2.351	1.973
m_{σ} (MeV)	595	535	547
a_{uu}^c (MeV/fm ²)	46.6	44.5	39.1
a_{uu}^{c0} (MeV)	-42.4	-72.3	-62.9

Masses of baryons & deuteron

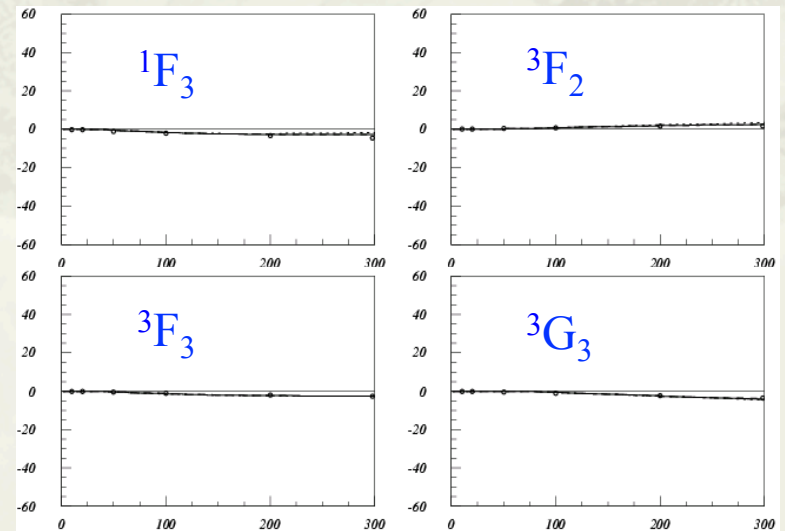
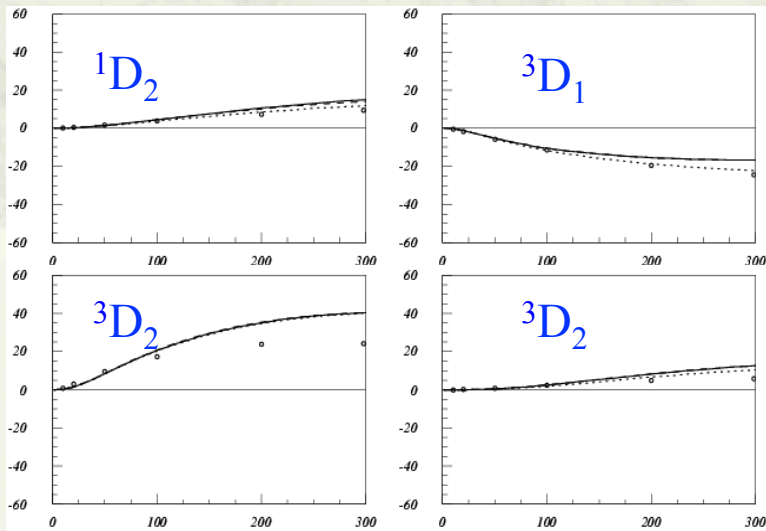
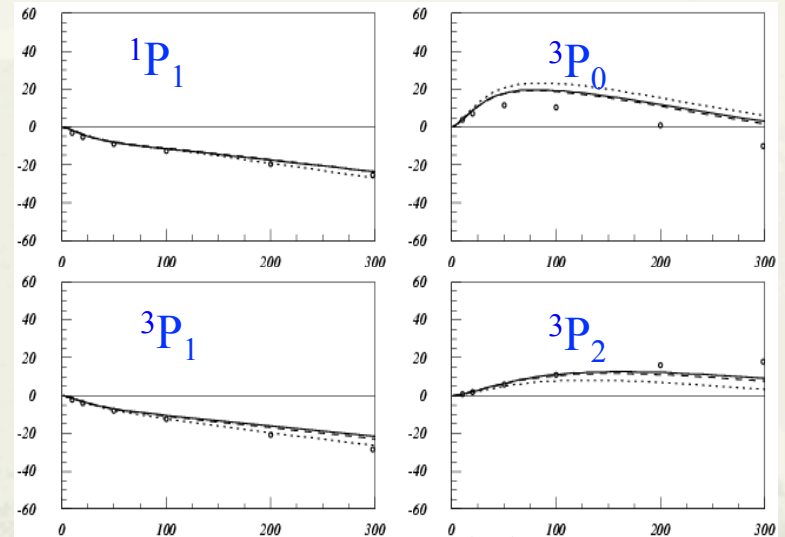
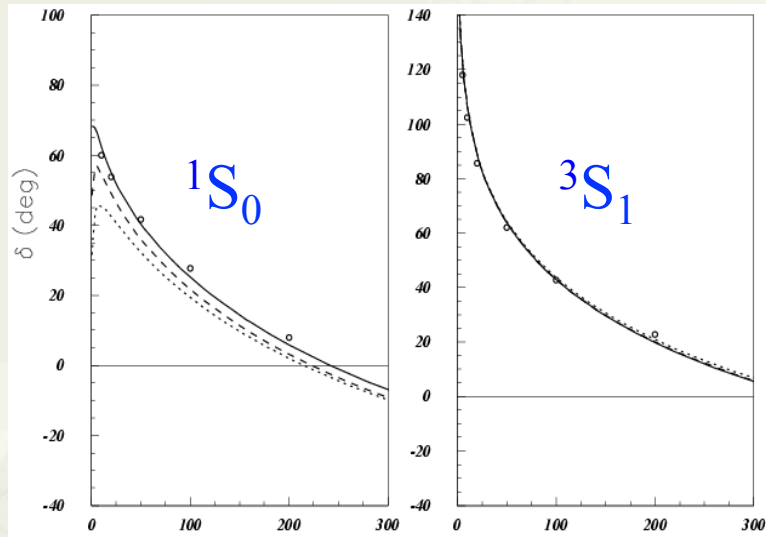
The masses of baryons

	N	Σ	Ξ	Λ	Δ	Σ^*	Ξ^*	Ω
Theor.	939	1194	1335	1116	1232	1370	1511	1656
Expt.	939	1194	1319	1116	1232	1385	1530	1672

The binding energy of deuteron

	Chiral SU(3) quark model	Extended chiral SU(3) quark model	
		$f_{chv}/g_{chv} = 0$	$f_{chv}/g_{chv} = 2/3$
B_{deu} (MeV)	2.13	2.19	2.14

NN phase shifts



RGM wave functions

RGM wave functions:

$$\begin{aligned} \psi_{6q} = & (1 - 9P_{36}) \left[\hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_{\Delta}^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \eta_{\Delta\Delta}(\vec{r}) \right]_{S=3, I=0, C=(00)} \\ & + (1 - 9P_{36}) \left[\hat{\phi}_C^{\text{int}}(\vec{\xi}_1, \vec{\xi}_2) \hat{\phi}_C^{\text{int}}(\vec{\xi}_4, \vec{\xi}_5) \eta_{CC}(\vec{r}) \right]_{S=3, I=0, C=(00)} \end{aligned}$$

Terms on r.h.s. not orthogonal to each other:

$$\langle \Delta\Delta | P_{36}^{sfc} | \Delta\Delta \rangle_{S=3, I=0} = -\frac{1}{9}$$

$$\langle CC | P_{36}^{sfc} | \Delta\Delta \rangle_{S=3, I=0} = -\frac{4}{9}$$

$$\langle CC | P_{36}^{sfc} | CC \rangle_{S=3, I=0} = -\frac{7}{9}$$

Not suitable for clarification of $\Delta\Delta$, CC components in d^*

d* charge distribution calculation

$$\begin{aligned} & \langle N(p') | J_N^\mu | N(p) \rangle \\ &= \frac{1}{1+\eta} \bar{u}_N(p', s') \left[(1+\eta) G_M \gamma^\mu - \frac{G_M - G_E}{2M_N} P^\mu \right] u_N(p, s). \end{aligned}$$

$$G_E^{d^*}(Q^2) = \frac{1}{7} \sum_{m_{d^*}=-3}^3 \langle p', m_{d^*} | J^0 | p, m_{d^*} \rangle$$

$$J^0 = \sum_{i=1}^6 e_i \bar{q}_i \gamma^0 q_i = \sum_{i=1}^6 j_i^0.$$

